# Observability on lattice points for heat equations and applications $\stackrel{\text{tr}}{\Rightarrow}$

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# Abstract

Observability inequalities on lattice points are established for non-negative solutions of the heat equation with potentials in the whole space. As applications, some controllability results of heat equations are derived by the above-mentioned observability

Keywords: Observability inequality, heat equation, lattice points

$$\begin{cases} \partial_t u = \Delta u, \quad (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d; \\ u(0, x) = u_0(x) \in L^2(\mathbb{R}^d). \end{cases}$$
(1.1)

$$\int_{\mathbb{R}^d} |u(t,x)|^2 \,\mathrm{d}x \le C(d,t,E) \int_0^t \int_E |u(s,x)|^2 \,\mathrm{d}x \,\mathrm{d}s$$

$$\left| E \bigcap (x + LQ) \right| \ge \gamma L^d$$
 for each  $x \in \mathbb{R}^d$ .

Observability inequalities on lattice points are established for 1 whole space. As applications, some controllability results of 1 inequalities. *Keywords:* Observability inequality, heat equation, lattice point 2010 MSC: 35K05, 93B05, 93B07 **1. Introduction** This is a continuous research of [6, 7] on observability in-equalities for the heat equation in  $\mathbb{R}^d$   $(d \ge 1)$   $\begin{cases} \partial_t u = \Delta u, \quad (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d; \\ u(0, x) = u_0(x) \in L^2(\mathbb{R}^d). \end{cases}$  (1.1) Recall that a measurable set  $E \subset \mathbb{R}^d$  is called an *observable set* if for every t > 0, there exists a constant C(d, t, E) > 0 so that when u solves (1.1),  $\int_{\mathbb{R}^d} |u(t, x)|^2 dx \le C(d, t, E) \int_0^t \int_E |u(s, x)|^2 dx ds.$ It was shown in [6] (see also [1]) that, E is an observable set if and only if E is  $\gamma$ -thick at scale L for some positive  $\gamma, L$ , namely,  $|E \bigcap (x + LQ)| \ge \gamma L^d$  for each  $x \in \mathbb{R}^d$ . Here Q is a unit cube in  $\mathbb{R}^d$ . Clearly, for every N > 0,  $E_N := \mathbb{Z}^d/N = \{n/N : n \in \mathbb{Z}^d\}$  is of zero measure (in the sense of d-dimensional Lebesgue measure), and thus it is not an observable set. sense of *d*-dimensional Lebesgue measure), and thus it is not an observable set.

It was also shown in [7] that, for every  $\varepsilon \in (0, 1)$  and t > 0, there exists a large enough  $N = N(t, \varepsilon) > 0$  so that we can, up to an  $\varepsilon$  error, recover the solution of (1.1) at the time t by observing the solution on the set  $E_N$  at the same time. More precisely, it follows from Theorem 1.2 (i) of [7] that, for every  $(\varepsilon, t) \in (0, 1) \times \mathbb{R}^+$ , there exists a constant C = C(d) > 1 so that,

Email addresses: mwangcug@outlook.com (Ming Wang), zhangcansx@163.com (Can Zhang), thanleon@163.com (Liang Zhang) if  $N \ge \sqrt{\frac{1}{t} \ln \frac{C}{\varepsilon}}$ , then each solution to (1.1) satisfies:

$$\int_{\mathbb{R}^d} |u(t,x)|^2 \,\mathrm{d}x \le 2N^{-d} \sum_{n \in \mathbb{Z}^d} \left| u(t,\frac{n}{N}) \right|^2 + \varepsilon \int_{\mathbb{R}^d} |u_0(x)|^2 \,\mathrm{d}x.$$
(1.2)

Then, the following two natural open questions are remained to study:

- (1) In general, can we remove the  $\varepsilon$ -term on the right hand side of (1.2)?
- (2) If not, for what kind of initial data, the  $\varepsilon$ -term in (1.2) can be removed?

For the first question, since  $E_N$  is not an observable set, it is natural to expect that  $\varepsilon$ -term can not be removed. Actually, we shall construct an explicit example to illustrate it. For the second question, we obtain some sufficient conditions, though it is too hard to give a complete characteristic for such kind of initial data. In all, our answers to these two questions are summarized in the following theorem.

- **Theorem 1.1.** (i) The  $\varepsilon$ -term in (1.2) can not be removed in general.
- (ii) Assume that  $u_0 \ge 0$  (or  $\le 0$ ). Then we have the following estimate for all solutions of (1.1)

$$\int_{\mathbb{R}^d} u^2(t,x) \, \mathrm{d}x \le 36^d e^{\frac{2d}{t}} \sum_{n \in \mathbb{N}^d} u^2(t,n), \quad t > 0.$$
(1.3)

Two remarks are given in order. First, the inequality (1.3) also holds (with a different upper bound constant) if the integral points  $\mathbb{N}^d$  is replaced by  $E_N$ , defined as before. Second, the proof of (1.3) is essentially based on a careful analysis of the heat kernel  $K(t, x, y) = (4\pi t)^{-d/2} e^{-\frac{|x-y|^2}{4t}}$ . In particular, we only

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need a Gaussian type upper bound and a lower bound of the kernel.

As it is well known that for a large class of potentials V(x),  $\Delta + V(x)$  generates an analytic semigroup  $e^{t(\Delta+V)}$  in  $L^2(\mathbb{R}^d)$ , and that the kernel of the semigroup  $e^{t(\Delta+V)}$  satisfies a two-side Gaussian type estimate. Thus, it is natural to extend the estimate in (ii) of Theorem 1.1 to heat equations with potentials.

To this end, we consider the heat equation with a potential

$$\begin{cases} \partial_t u = (\Delta + V(x))u, \quad (t, x) \in \mathbb{R}^+ \times \mathbb{R}^d; \\ u(0, x) = u_0(x) \in L^2(\mathbb{R}^d). \end{cases}$$
(1.4)

Here  $V : \mathbb{R}^d \to \mathbb{R}$  depends only on the spatial variable. To state our result, we need the uniformly local Lebesgue integral spaces  $L^p_{U,loc}(\mathbb{R}^d)$ ,  $p \ge 1$ , which are Banach spaces endowed with norms

$$||f||_{L^p_{U,loc}(\mathbb{R}^d)} := \sup_{x \in \mathbb{R}^d} \left( \int_{|x-y| \le 1} |f(y)|^p \, \mathrm{d}y \right)^{\frac{1}{p}}.$$

Clearly, the usual Lebesgue space  $L^p(\mathbb{R}^d)$  is continuous embedding into  $L^p_{Uloc}(\mathbb{R}^d)$ .

**Theorem 1.2.** Let V be a real-valued function belonging to  $L_{U,loc}^{p}(\mathbb{R}^{d})$  with  $p > \max\{1, \frac{d}{2}\}$ . Assume that  $u_{0} \ge 0$  (or  $\le 0$ ). Then there exists a constant C = C(d, V) > 0 so that the following estimate hold for all solutions of (1.4)

$$\int_{\mathbb{R}^d} u^2(t,x) \,\mathrm{d}x \le e^{C(1+t+\frac{1}{t})} \sum_{n \in \mathbb{N}^d} u^2(t,n), \quad t > 0.$$

The paper is organized as follows. In Section 2, we prove Theorem 1.1 and Theorem 1.2. In Section 3, we give some applications of the observability inequality in (ii) of Theorem 1.1 in Control Theory.

### 2. Proofs of main results

In the sequel, for every  $x \in \mathbb{R}^d$  and r > 0, we use  $Q_r(x)$  to denote the closed cube in  $\mathbb{R}^d$  centered at x with side length r; We denote by  $A^c$  the complement set of A.

**Lemma 2.1.** For any a > 0 and  $y \in Q_4^c(0)$ , we have

$$\sup_{\mathbf{x}\in Q_2(0)} e^{-a|\mathbf{x}-\mathbf{y}|^2} \le 2^{(d-1)} e^{\frac{(d-1)a}{2}} \sum_{n\in Q_2(0)\cap \mathbb{N}^d} e^{-a|n-\mathbf{y}|^2}.$$
 (2.1)

*Proof.* In the case that d = 1, the inequality (2.1) holds obviously. We next assume that  $d \ge 2$ . Arbitrarily give  $y \in Q_4^c(0)$ . Since  $e^{-a|x-y|^2}$  is a continuous function of x in  $Q_2(0)$ , the maximum of  $e^{-a|x-y|^2}$  can be obtained at some point  $x^* = (x_1^*, x_2^*, \dots, x_d^*)$ . Note that

$$\max_{x \in Q_2(0)} e^{-a|x-y|^2} = \max_{x \in Q_2(0)} \prod_{i=1}^d e^{-a|x_i-y_i|^2}$$
$$= \prod_{i=1}^d \max_{-1 \le x_i \le 1} e^{-a|x_i-y_i|^2} = \prod_{i=1}^d e^{-a|x_i^*-y_i|^2},$$

where  $x_i^*$  takes the form

$$x_i^* = \begin{cases} 1, & y_i \ge 1; \\ y_i, & -1 < y_i < 1; \\ -1, & y_i \le -1. \end{cases}$$
(2.2)

We write  $y = (y_1, y_2, \dots, y_d)$  and divide the set  $\{y_i, i = 1, 2, \dots, d\}$  into two groups:  $|y_i| \ge 1$  and  $|y_i| < 1$ . Since  $y \in Q_4^c(0)$ , there exists at least one of  $\{y_i, i = 1, 2, \dots, d\}$  such that  $|y_i| \ge 2$ . Thus there are at most (d - 1) elements of  $\{y_i, i = 1, 2, \dots, d\}$  satisfying  $|y_i| < 1$ .

Without loss of generality we can assume that for some  $j \le d-1$ 

$$\begin{cases} |y_i| \ge 1, \quad j+1 \le i \le d \\ -1 < y_i < 1, \quad i \le j. \end{cases}$$
(2.3)

Then it follows from (2.2) and (2.3) that

$$x_i^* = \begin{cases} 1 \text{ or } -1, & j+1 \le i \le d, \\ y_i, & i \le j. \end{cases}$$
(2.4)

Then we have

$$\sup_{x \in Q_2(0)} e^{-a|x-y|^2} = e^{-a|x^*-y|^2} = \prod_{j+1 \le i \le d} e^{-a|x_i^*-y_i|^2}.$$
 (2.5)

On the other hand, using (2.4), we have

$$\sum_{\substack{n=(n_1,n_2,\cdots,n_d)\\n_i\in\{-1,0,1\},i=1,2,\cdots,d}} e^{-a|n-y|^2} \sum_{\substack{j+1\leq i\leq d}} e^{-a|x_i^*-y_i|^2} \sum_{\substack{n_i\in\{-1,0,1\},\ i\leq j\\i=1,2,\cdots,j}} \prod_{i\leq j} e^{-a|n_i-y_i|^2}.$$
 (2.6)

This is because every term on the right hand side of (2.6) appears on the left, and every term on the left is non-negative. Combining (2.5) and (2.6), we find that

$$\sum_{n \in Q_2(0) \cap \mathbb{N}^d} e^{-a|n-y|^2} = \sum_{\substack{n = (n_1, n_2, \cdots, n_d) \\ n_i \in \{-1, 0, 1\}, i=1, 2, \cdots, d}} e^{-a|n-y|^2} \ge \Theta \sup_{x \in Q_2(0)} e^{-a|x-y|^2},$$

where

$$\Theta = \sum_{\substack{n_i \in \{-1,0,1\}, \\ i=1,2,\cdots, j \le d-1}} \prod_{i \le j \le d-1} e^{-a|n_i - y_i|^2}$$

Thus (2.1) holds if one can show that  $2^{(d-1)}e^{\frac{(d-1)a}{2}}\Theta \ge 1$ . In fact, if we write

$$\sum_{\substack{n_i \in \{-1,0,1\}, i \le j \\ i=1,2,\cdots,j}} \prod_{i \le j} e^{-a|n_i - y_i|^2} = \prod_{i \le j} A_i,$$
(2.7)

with

$$A_i = \sum_{k \in \{-1,0,1\}} e^{-a|k-y_i|^2}, \quad -1 < y_i < 1,$$

then for every  $1 \le i \le j$ , by the definition of  $A_i$  we have

$$A_{i} \geq \begin{cases} \sum_{k \in \{0,1\}} e^{-a|k-y_{i}|^{2}}, & 0 \leq y_{i} < 1; \\ \sum_{k \in \{-1,0\}} e^{-a|k-y_{i}|^{2}}, & -1 < y_{i} < 0. \end{cases}$$
(2.8)

Thanks to (2.8), for  $0 \le y_i < 1$  we have

$$\begin{split} A_i &\geq e^{-a|1-y_i|^2} + e^{-ay_i^2} \geq \frac{1}{2}e^{-\frac{a}{2}(|1-y_i|^2+y_i^2)} \\ &= \frac{1}{2}e^{a(y_i-y_i^2-\frac{1}{2})} \geq \frac{1}{2}e^{-\frac{a}{2}}. \end{split}$$

Similarly, for  $-1 < y_i \le 0$  we also have

$$A_i \ge e^{-a|-1-y_i|^2} + e^{-ay_i^2} \ge \frac{1}{2}e^{-\frac{a}{2}}.$$

Thus, we always have

$$A_i \ge \frac{1}{2}e^{-\frac{a}{2}}, \quad i \le j.$$
 (2.9)

It follows from (2.7) and (2.9) that

$$\sum_{\substack{n_i \in \{-1,0,1\}, \\ i=1,2,\cdots, j \leq d-1}} \prod_{i \leq j \leq d-1} e^{-a|n_i - y_i|^2} \geq 2^{-j} e^{-\frac{aj}{2}} \geq 2^{-(d-1)} e^{-\frac{(d-1)a}{2}}.$$

Thus  $2^{(d-1)}e^{\frac{(d-1)a}{2}}\Theta \ge 1$ . This completes the proof.

**Lemma 2.2.** For any a > 0 and  $y \in Q_4(0)$ , we have

$$\sup_{x \in Q_2(0)} e^{-a|x-y|^2} \le e^{4ad} \sum_{n \in Q_2(0) \cap \mathbb{N}^d} e^{-a|n-y|^2}.$$

*Proof.* Arbitrarily fix  $y \in Q_4(0)$ . Since  $\sup_{x \in Q_2(0)} e^{-a|x-y|^2} \le 1$ , it suffices to show that

$$1 \le e^{4ad} \sum_{n \in Q_2(0) \cap \mathbb{N}^d} e^{-a|n-y|^2}.$$
 (2.10)

Since  $|y| \le 2\sqrt{d}$ , we have  $\sum_{n \in Q_2(0) \cap \mathbb{N}^d} e^{-a|n-y|^2} \ge e^{-a|y|^2} \ge e^{-4ad}$ . This gives (2.10) and finishes the proof.

**Lemma 2.3.** For any a > 0 and  $x, y \in \mathbb{R}^d$ , we have

$$\sup_{z \in Q_2(x)} e^{-a|z-y|^2} \le 2^{d-1} e^{4ad} \sum_{n \in Q_2(x) \cap (\mathbb{N}^d + x)} e^{-a|n-y|^2}.$$

*Proof.* As a direct corollary of Lemma 2.1 and Lemma 2.2, it holds that

$$\sup_{z \in Q_2(0)} e^{-a|z-y|^2} \le 2^{d-1} e^{4ad} \sum_{n \in Q_2(0) \, \bigcap \, \mathbb{N}^d} e^{-a|n-y|^2}, \quad y \in \mathbb{R}^d.$$

By changing the variable  $y \mapsto x + y$ , the desired conclusion follows at once.

**Theorem 2.1.** Assume that  $u_0 \ge 0$ . Then for all solutions of (1.1) and all t > 0 and  $k \in \mathbb{N}^d$ 

$$u(t,x) \le 2^{d-1} e^{\frac{d}{t}} \sum_{n \in Q_2(k) \cap \mathbb{N}^d} u(t,n), \quad x \in Q_2(k).$$
(2.11)

*Proof.* Using the heat kernel, the solution of the heat equation can be written as

$$u(t,x) = \frac{1}{(4\pi t)^{n/2}} \int_{\mathbb{R}^d} e^{-\frac{|x-y|^2}{4t}} u_0(y) \,\mathrm{d}y, \quad x \in \mathbb{R}^d.$$
(2.12)

Applying Lemma 2.3 with  $a = \frac{1}{4t}$ , we obtain

$$e^{-\frac{|x-y|^2}{4t}} \le 2^{d-1} e^{\frac{d}{t}} \sum_{n \in Q_2(k) \cap \mathbb{N}^d} e^{-\frac{|n-y|^2}{4t}}, \quad x \in Q_2(k).$$
(2.13)

Since  $u_0 \ge 0$ , the inequality (2.11) follows from (2.12) and (2.13).

**Remark 2.1.** Notice that Theorem 2.1 does not follow from the parabolic Harnack inequality. The classical Harnack inequality [2] says that, for every  $t' > t > 0, k \in \mathbb{N}^d$ , every non-negative solution of (1.1) satisfies that

$$\max_{x \in Q_2(k)} u(t, x) \le C(d, t, t') \inf_{x \in Q_2(k)} u(t', x).$$
(2.14)

The condition t' > t is essential here. The time t' cannot be equal to t in (2.14). To see this, without loss of generality, it suffices to construct a non-negative solution such that the following fails:

$$u(t, x_0) \le 2^{d-1} e^{\frac{d}{t}} u(t, 0), \tag{2.15}$$

where  $x_0 = (1, 0, \dots, 0) \in \mathbb{R}^d$ .

To this end, for every M > 0, set  $u_{0M}(x) = (4\pi)^{-\frac{d-1}{2}}\chi_{\{M \le x_1 \le M+1\}}e^{-\frac{|x'|^2}{4}}$ , where  $\chi_{\{M \le x_1 \le M+1\}}$  is the characteristic function of the interval [M, M+1],  $x = (x_1, x') \in \mathbb{R}^d$ . Clearly,  $u_{0M}$  is uniformly bounded in  $L^2(\mathbb{R}^d)$ . Using the heat kernel, we find the solution of the heat equation (1.1) with initial datum  $u_{0M}$  is given by

$$u_M(t,x) = (4\pi(t+1))^{-\frac{d-1}{2}} e^{-\frac{|x'|^2}{4(t+1)}} (4\pi t)^{-1/2} \int_M^{M+1} e^{-\frac{|x_1-y_1|^2}{4t}} \, \mathrm{d}y_1,$$

for all  $t > 0, x \in \mathbb{R}^d$ . By some computations, we have for t > 0

$$u_{M}(t, x_{0}) = (4\pi(t+1))^{-\frac{d-1}{2}} (4\pi t)^{-\frac{1}{2}} \int_{M}^{M+1} e^{-\frac{|1-y_{1}|^{2}}{4t}} dy_{1}$$
$$\geq \frac{1}{2} (4\pi(t+1))^{-\frac{d-1}{2}} (4\pi t)^{-\frac{1}{2}} e^{-\frac{|\frac{1}{2}-M|^{2}}{4t}}, \qquad (2.16)$$

$$u_M(t,0) = (4\pi(t+1))^{-\frac{d-1}{2}} (4\pi t)^{-\frac{1}{2}} \int_M^{M+1} e^{-\frac{|y_1|^2}{4t}} \, \mathrm{d}y_1$$
  
$$\leq (4\pi(t+1))^{-\frac{d-1}{2}} (4\pi t)^{-\frac{1}{2}} e^{-\frac{M^2}{4t}}.$$
(2.17)

Combining (2.16) and (2.17) gives that

$$u_M(t, x_0) \ge \frac{1}{2} u_M(t, 0) e^{\frac{M - \frac{1}{4}}{4t}}, \quad t > 0.$$
 (2.18)

When M is large enough, (2.18) obviously contradicts with (2.15).

**Proof of (ii) of Theorem 1.1.** Without loss of generality, we can assume that  $u_0 \ge 0$ . Arbitrarily fix  $k \in \mathbb{N}^d$ . The number of points of the set  $Q_2(k) \cap \mathbb{N}^d$  is  $3^d$ . Using the elementary inequality  $(\sum_{1 \le i \le m} a_i)^2 \le m \sum_{1 \le i \le m} a_i^2$  with  $m = 3^d$ , we deduce from (2.11) that

$$u^{2}(t,x) \leq 12^{d} e^{\frac{2d}{t}} \sum_{n \in Q_{2}(k) \cap \mathbb{N}^{d}} u^{2}(t,n), \quad x \in Q_{2}(k).$$
(2.19)

Integrating (2.19) on  $x \in Q_2(k)$ , noting that the volume of  $Q_2(k)$  is  $2^d$ , we get

$$\int_{Q_2(k)} u^2(t,x) \, \mathrm{d}x \le 24^d e^{\frac{2d}{t}} \sum_{n \in Q_2(k) \cap \mathbb{N}^d} u^2(t,n).$$
(2.20)

Finally, summarizing (2.20) for  $k \in \mathbb{N}^d$  we deduce that

$$\int_{\mathbb{R}^d} u^2(t, x) \, \mathrm{d}x = 2^{-d} \sum_{k \in \mathbb{N}^d} \int_{Q_2(k)} u^2(t, x) \, \mathrm{d}x$$
  
$$\leq 12^d e^{\frac{2d}{t}} \sum_{k \in \mathbb{N}^d} \sum_{n \in Q_2(k) \cap \mathbb{N}^d} u^2(t, n) \leq 36^d e^{\frac{2d}{t}} \sum_{n \in \mathbb{N}^d} u^2(t, n).$$

This proves the theorem.

**Proof of (i) of Theorem 1.1.** For every T > 0 and N > 0, it suffices to show that there exists an initial datum  $u_0 \in L^2(\mathbb{R}^d)$  such that the following inequality fails

$$\int_{\mathbb{R}^d} |u(T, x)|^2 \, \mathrm{d}x \le N^{-d} \sum_{n \in \mathbb{Z}^d} |u(T, \frac{n}{N})|^2.$$
(2.21)

We first consider the case that d = 1. Define an initial datum  $u_0$  via Fourier transform<sup>1</sup>

$$\widehat{u_0}(\xi) = \frac{\pi}{i} e^{T\xi^2} \Big( \widehat{f}(\xi - N\pi) - \widehat{f}(\xi + N\pi) \Big),$$

where  $\widehat{f}(\xi) = e^{-(T+1)\xi^2}, \xi \in \mathbb{R}$ . Clearly,

$$e^{T\xi^2}\widehat{f}(\xi - N\pi) = e^{-\xi^2 + 2(T+1)N\pi\xi - (T+1)(N\pi)^2}$$

belongs to  $L^2(\mathbb{R})$ . So does  $e^{T\xi^2}\widehat{f}(\xi + N\pi)$ . Thus,  $||u_0||_{L^2(\mathbb{R})} = ||\widehat{u_0}||_{L^2(\mathbb{R})} < \infty$ . Since  $u(T, x) = (e^{T\Delta}u_0)(x)$ , we find

$$\widehat{u(T,\cdot)}(\xi) = e^{-T\xi^2} \widehat{u_0}(\xi) = \frac{\pi}{i} \Big( \widehat{f}(\xi - N\pi) - \widehat{f}(\xi + N\pi) \Big).$$

Taking the inverse Fourier transform we obtain

$$u(T, x) = \frac{1}{2i} (e^{iN\pi x} - e^{-iN\pi x}) f(x) = \sin(N\pi x) f(x).$$

Since *f* is a bounded smooth function, we find  $u(T, \frac{n}{N}) = 0$  for  $n \in \mathbb{Z}$ . However, it is clear that  $||u(T, \cdot)||_{L^2(\mathbb{R})} > 0$ . This leads to a contradiction with (2.21) in one dimension.

In higher dimensions, set

$$\widehat{u_0}(\xi) = \frac{(2\pi)^d}{2i} e^{T|\xi|^2} \prod_{i=1}^d \left(\widehat{f}(\xi_i - N\pi) - \widehat{f}(\xi_i + N\pi)\right)$$

with  $\widehat{f}(\xi_i) = e^{-(T+1)\xi_i^2}, \xi_i \in \mathbb{R}$ . Similar to the analysis above, we find

$$u(T, \frac{n}{N}) = 0, n \in \mathbb{Z}^d$$
, but  $||u(T, \cdot)||_{L^2(\mathbb{R}^d)} > 0$ .

This completes the proof.

**Definition 2.1.** We say that a function  $V(\cdot)$  in  $L^1_{loc}(\mathbb{R}^d)$  satisfies two-side Gaussian type heat kernel estimates, if the operator  $\Delta + V$  generates an analytic semigroup  $e^{t(\Delta+V(x))}$  in  $L^2(\mathbb{R}^d)$ , and if there exist positive constants  $c_i$  (i = 1, 2, 3, 4), with  $c_2 \le c_4$ , so that for all  $t > 0, x, y \in \mathbb{R}^d$ 

$$t^{-\frac{d}{2}}e^{-c_1(t+1)}e^{-\frac{|x-y|^2}{c_2t}} \le K(t,x,y) \le t^{-\frac{d}{2}}e^{c_3(t+1)}e^{-\frac{|x-y|^2}{c_4t}},$$

where K(t, x, y) is the kernel of the semigroup  $e^{t(\Delta+V)}$ , namely,

$$(e^{t(\Delta+V)}f)(x) = \int_{\mathbb{R}^d} K(t,x,y)f(y)\,\mathrm{d}y.$$

**Theorem 2.2.** Let V be a real-valued function satisfying twoside Gaussian type heat kernel estimates. Assume that  $u_0 \ge 0$ (or  $\le 0$ ). Then there exists a positive constant C = C(d, V) such that the following estimate hold for all solutions of (1.4)

$$\int_{\mathbb{R}^d} u^2(t, x) \, \mathrm{d}x \le e^{C(1+t+\frac{1}{t})} \sum_{n \in \mathbb{N}^d} u^2(t, n), \quad t > 0$$

*Proof.* We only consider the case that  $u_0 \ge 0$ . Since both *V* and  $u_0$  are real-valued, the solution *u* of (1.4) is also real-valued. According to the definition of the kernel K(t, x, y), we have

$$u(t, x) = \int_{\mathbb{R}^d} K(t, x, y) u_0(y) \, \mathrm{d}y.$$

Since *V* satisfies two-side Gaussian type heat kernel estimates, we find that for all t > 0 and  $x \in \mathbb{R}^d$ 

$$0 \le u(t, x) \le \int_{\mathbb{R}^d} t^{-\frac{d}{2}} e^{c_3(t+1)} e^{-\frac{|x-y|^2}{c_4 t}} u_0(y) \, \mathrm{d}y$$
$$= (c_4 \pi)^{\frac{d}{2}} e^{c_3(t+1)} (e^{\frac{c_4}{4} t\Delta} u_0)(x).$$
(2.22)

We apply Theorem 1.1 with t replaced by  $\frac{c_4}{4}t$  to obtain that

$$\int_{\mathbb{R}^d} ((e^{\frac{c_4}{4}t\Delta}u_0)(x))^2 \, \mathrm{d}x \le 36^d e^{\frac{8d}{c_4t}} \sum_{n\in\mathbb{N}^d} ((e^{\frac{c_4}{4}t\Delta}u_0)(n))^2.$$
(2.23)

Combining (2.22) and (2.23) gives that for all t > 0

$$\int_{\mathbb{R}^d} u^2(t,x) \,\mathrm{d}x$$
  

$$\leq 36^d (c_4 \pi)^d e^{2c_3(t+1)} e^{\frac{8d}{c_4 t}} \sum_{n \in \mathbb{N}^d} ((e^{\frac{c_4}{4} t\Delta} u_0)(n))^2. \quad (2.24)$$

Replacing *t* by  $c_2t/c_4$  in (2.24) gives that for all t > 0

$$\int_{\mathbb{R}^{d}} u^{2}(\frac{c_{2}}{c_{4}}t, x) \, \mathrm{d}x$$

$$\leq 36^{d}(c_{4}\pi)^{d} e^{2c_{3}(\frac{c_{2}}{c_{4}}t+1)} e^{\frac{8d}{c_{2}t}} \sum_{n \in \mathbb{N}^{d}} ((e^{\frac{c_{2}}{4}t\Delta}u_{0})(n))^{2}. \quad (2.25)$$

On the other hand, using the lower bound of the kernel, we find that

$$u(t,x) \ge \int_{\mathbb{R}^d} t^{-\frac{d}{2}} e^{-c_1(t+1)} e^{-\frac{|x-y|^2}{c_2 t}} u_0(y) \, \mathrm{d}y$$
  
=  $(c_2 \pi)^{\frac{d}{2}} e^{-c_1(t+1)} (e^{\frac{c_2}{4} t\Delta} u_0)(x).$  (2.26)

<sup>&</sup>lt;sup>1</sup>Here the Fourier transform is defined as  $\mathcal{F}f = \widehat{f}(\xi) = \int_{\mathbb{R}^d} e^{-ix\xi} f(x) \, dx$ , and the inverse Fourier transform is  $\mathcal{F}^{-1}f(x) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{ix\xi} f(\xi) \, d\xi$ .

It follows from (2.26) that for all t > 0 and  $n \in \mathbb{N}^d$ 

$$\left(\left(e^{\frac{c_2}{4}t\Delta}u_0\right)(n)\right)^2 \le (c_2\pi)^{-d}e^{2c_1(t+1)}u^2(t,n).$$
(2.27)

Inserting (2.27) into (2.25) we get

$$\int_{\mathbb{R}^d} u^2 (\frac{c_2}{c_4}t, x) \, \mathrm{d}x \le (\frac{36c_4}{c_2})^d e^{2c_3(\frac{c_2}{c_4}t+1)} e^{2c_1(t+1)} e^{\frac{8d}{c_2t}} \sum_{n \in \mathbb{N}^d} u^2(t, n).$$
(2.28)

Moreover, using the upper bound of *K* again, we find that for all t > 0

$$\begin{aligned} \|e^{t(\Delta+V)}\|_{L^{2}(\mathbb{R}^{d}),L^{2}(\mathbb{R}^{d})} &\leq \|t^{-\frac{d}{2}}e^{c_{3}(t+1)}e^{-\frac{|x|^{2}}{c_{4}t}}\|_{L^{1}(\mathbb{R}^{d})} \\ &= (c_{4}\pi)^{d/2}e^{c_{3}(t+1)}. \end{aligned}$$
(2.29)

Noting that  $c_2/c_4 \le 1$ , it follows from (2.29) (since  $u(t, x) = (e^{t(\Delta+V)}u_0)(x)$ ) that

$$\int_{\mathbb{R}^d} u^2(t,x) \, \mathrm{d}x \le (c_4 \pi)^d e^{2c_3(2 - \frac{c_2}{c_4})} \int_{\mathbb{R}^d} u^2(\frac{c_2}{c_4}t,x) \, \mathrm{d}x.$$
(2.30)

Finally, combining (2.28) and (2.30) gives the desired conclusion.  $\hfill \Box$ 

**Proof of Theorem 1.2.** Since  $V \in L^p_{U,loc}(\mathbb{R}^d)$  with  $p > \max\{1, \frac{d}{2}\}$ , it is easy to check (see also [9, Proposition 2.1]) that V belongs to the Kato class  $K_d$  (see [4, p. 453] for a precise definition). According to [4, Theorem 7.1], for all  $\varepsilon > 0$ , there exist positive constants  $C_1(\varepsilon), C_2(\varepsilon)$  so that the kernel K(t, x, y) of the analytic semigroup  $e^{t(\Delta+V)}$  satisfies

$$t^{-\frac{d}{2}}e^{-C_1(\varepsilon)(1+t)}e^{-\frac{|x-y|^2}{4(1-\varepsilon)t}} \le K(t,x,y) \le t^{-\frac{d}{2}}e^{C_2(\varepsilon)(1+t)}e^{-\frac{|x-y|^2}{4(1+\varepsilon)t}}$$

for all  $t > 0, x, y \in \mathbb{R}^d$ . Thus *V* satisfies a two-side Gaussian type heat kernel estimate. Then Theorem 1.2 follows from Theorem 2.2 directly.

# 3. Applications to controllability

In this section, we will show an application of Theorem 1.1 for an impulsive controllability for the heat equation in  $\mathbb{R}^d$ . We refer the interesting reader to [1, 6] for the null controllability of the heat equation in  $\mathbb{R}^d$  with distributed controls.

Arbitrarily fix T > 0 and  $\tau \in (0, T)$ . Consider the heat equation with impulsive control

$$\begin{cases} \partial_t y(t, x) - \Delta y(t, x) = 0 & \text{ in } (0, T) \times \mathbb{R}^d, \\ y(\tau, x) = y(\tau^-, x) + Bv & \text{ in } \mathbb{R}^d, \\ y(0) = y_0(x) & \text{ in } \mathbb{R}^d. \end{cases}$$
(3.1)

Here, *y* is the state variable,  $y_0 \in L^2(\mathbb{R}^d)$ ,  $y(\tau^-, x)$  denotes the left limit of  $y(\cdot, x)$  (treated as a function from  $\mathbb{R}^+$  to  $\mathbb{R}^d$ ) for each *x* at time  $\tau$ , and  $v \in \ell^2(\mathbb{R}^d)$  is the control. The control operator  $B : \ell^2(\mathbb{R}^d) \to \mathcal{D}'(\mathbb{R}^d)$  is defined by

$$Bv := \sum_{n \in \mathbb{Z}^d} v_n \delta_n$$
 for each  $v = (v_n)_{n \in \mathbb{Z}^d} \in \ell^2(\mathbb{R}^d)$ ,

where  $\delta_n(x) := \delta_n(x - n)$ ,  $x \in \mathbb{R}^d$ , with  $\delta$  being the Dirac measure. In fact, it is not hard to check that *B* is linear and bounded from  $\ell^2(\mathbb{R}^d)$  to  $H^{-s}(\mathbb{R}^d)$  when s > d/2 (see also [7]).

We first quote from [7] the following result concerning the well-posedness of (3.1).

**Proposition 3.1.** If s > d/2,  $y_0 \in L^2(\mathbb{R}^d)$  and  $v \in \ell^2(\mathbb{R}^d)$ , then (3.1) has a unique solution in  $C([0, \tau) \cup (\tau, T]; L^2(\mathbb{R}^d)) \cap C([\tau, T]; H^{-s}(\mathbb{R}^d))$ .

The main result of this section is stated as follows.

**Theorem 3.1.** Let  $0 < \tau < T$ . For each  $y_0 \in L^2(\mathbb{R}^d)$ , there is a control  $v \in \ell^2(\mathbb{R}^d)$ , with

$$||v||_{\ell^2} \le C e^{\frac{C}{T-\tau}} ||y_0||,$$

such that the solution of (3.1) verifies  $y(T, x; v) \ge 0$  for a.e.  $x \in \mathbb{R}^d$ .

Here and in the sequel, we write  $\langle \cdot, \cdot \rangle$  and  $\|\cdot\|$  for the usual inner product and norm in  $L^2(\mathbb{R}^d)$ , and denote by  $\langle \cdot, \cdot \rangle_{\ell^2}$  and  $\|\cdot\|_{\ell^2}$  the usual inner product and norm in  $\ell^2(\mathbb{R}^d)$ , respectively.

**Remark 3.1.** Similarly, for each  $y_0 \in L^2(\mathbb{R}^d)$ , there exists  $v \in \ell^2(\mathbb{R}^d)$  such that the solution of (3.1) satisfies  $y(T, x; v) \leq 0$  for *a.e.*  $x \in \mathbb{R}^d$ .

**Remark 3.2.** Analogous results can be established for the heat equation with potentials by using Theorem 1.2 instead.

We point out that the proof of Theorem 3.1 is motivated and adapted from the arguments in [8, Theorem 2.4] (see also [3]). Before giving the proof of Theorem 3.1, we start with some preliminaries. First of all, we quote from [5, Chapter I, Theorem 1.2] the following classical result in the Calculus of Variations.

**Proposition 3.2.** Let X be a reflexive Banach space and let K be a weakly closed subset. If a weakly lower semi-continuous functional  $F : K \to \mathbb{R}$  satisfies the following coercive condition

$$\lim_{\substack{z \in K \\ \|z\|_X \to +\infty}} F(z) = +\infty,$$

then F attains its minimum in K, i.e., there exists at least one  $\hat{z}$  so that

$$F(\hat{z}) = \min_{z \in K} F(z).$$

Moreover, if F is strictly convex in K then it has a unique minimizer.

**Remark 3.3.** Notice that closed and convex subsets of Banach spaces are important examples of weakly closed sets.

In the sequel, we define

$$L^{2}_{+}(\mathbb{R}^{d}) := \{ y \in L^{2}(\mathbb{R}^{d}) : y(x) \ge 0 \text{ for a.e. } x \in \mathbb{R}^{d} \}.$$

Clearly, it is a closed and convex subset of  $L^2(\mathbb{R}^d)$ .

For each  $T > \tau > 0$  and  $\varepsilon > 0$ , we define a functional  $F_{\varepsilon}^{T,\tau}$ :  $L^2_+(\mathbb{R}^d) \to \mathbb{R}$  by

$$F_{\varepsilon}^{T,\tau}(\varphi_T) := \frac{1}{2} \|B^*\varphi(\tau)\|_{\ell^2}^2 + \langle y_0, \varphi(0) \rangle + \varepsilon \|\varphi_T\|$$

for any  $\varphi_T \in L^2_+(\mathbb{R}^d)$ , where  $y_0$  is the given initial state of (3.1) and  $\varphi$  solves the so-called adjoint equation

$$\begin{cases} \partial_t \varphi(t, x) + \Delta \varphi(t, x) = 0 & \text{ in } (0, T) \times \mathbb{R}^d, \\ \varphi(T, x) = \varphi_T(x) & \text{ in } \mathbb{R}^d, \end{cases}$$
(3.2)

and  $B^* : C_0^{\infty}(\mathbb{R}^d) \to \ell^2(\mathbb{R}^d)$  is the adjoint operator of *B*. It is clear that  $B^*$  is linear and bounded from  $H^s(\mathbb{R}^d)$  to  $\ell^2(\mathbb{R}^d)$  with s > d/2.

**Lemma 3.1.** Given  $T > \tau > 0$  and  $y_0 \in L^2(\mathbb{R}^d)$ , for each  $\varepsilon > 0$ ,  $F_{\varepsilon}^{T,\tau}$  has a unique minimizer, denoted by  $\widehat{\varphi}_T^{\varepsilon}$ , in  $L^2_+(\mathbb{R}^d)$ . Furthermore, for all  $\varphi_T \in L^2_+(\mathbb{R}^d)$ 

$$\langle B^* \widehat{\varphi}^{\varepsilon}(\tau), B^* \varphi(\tau) \rangle_{\ell^2} + \langle y_0, \varphi(0) \rangle + \varepsilon ||\varphi_T|| \ge 0, \tag{3.3}$$

where  $\widehat{\varphi}^{\varepsilon}$  and  $\varphi$  are the solutions of (3.2) with  $\widehat{\varphi}^{\varepsilon}_{T}$  and  $\varphi_{T}$ , respectively.

*Proof.* It is not hard to check that  $F_{\varepsilon}^{T,\tau}$  is strictly convex and weakly lower semi-continuous in  $L_{+}^{2}(\mathbb{R}^{d})$ . We next show that  $F_{\varepsilon}^{T,\tau}$  satisfies the coercive condition, i.e.,

$$\liminf_{\substack{\varphi_T \in L^2_+(\mathbb{R}^d) \\ \|\varphi_T\| \to +\infty}} F^{T,\tau}_{\varepsilon}(\varphi_T) = +\infty$$

To seek a contradiction, we would assume that there was a sequence  $\{\varphi_{T}^{n}\}_{n\geq 1} \subset L^{2}_{+}(\mathbb{R}^{d})$  to be such that

$$\lim_{n \to +\infty} \|\varphi_T^n\| = +\infty$$

and

$$F_{\varepsilon}^{T,\tau}(\varphi_T^n) < +\infty \text{ for all } n \in \mathbb{N}.$$

Let us set

$$\widetilde{\varphi}_T^n := \frac{\varphi_T^n}{\|\varphi_T^n\|}$$
 for all  $n \in \mathbb{N}$ .

Clearly,  $\widetilde{\varphi}_T^n \in L^2_+(\mathbb{R}^d)$ . Then

$$\frac{F_{\varepsilon}^{T,\tau}(\varphi_T^n)}{\|\varphi_T^n\|} = \frac{1}{2} \|\varphi_T^n\| \|B^* \widetilde{\varphi}^n(\tau)\|_{\ell^2}^2 + \langle y_0, \widetilde{\varphi}^n(0) \rangle + \varepsilon,$$

where  $\widetilde{\varphi}^n$  is the solution to

$$\begin{cases} \partial_t \widetilde{\varphi}^n + \Delta \widetilde{\varphi}^n = 0 & \text{ in } (0, T) \times \mathbb{R}^d, \\ \widetilde{\varphi}^n(T) = \widetilde{\varphi}^n_T & \text{ in } \mathbb{R}^d. \end{cases}$$

As  $F_{\varepsilon}^{T,\tau}(\varphi_T^n)$  is uniformly bounded, we have

$$\lim_{n \to +\infty} \|B^* \widetilde{\varphi}^n(\tau)\|_{\ell^2} = 0.$$
(3.4)

Because

$$\|\widetilde{\varphi}_T^n\| = 1$$
 for all  $n \ge 1$ ,

there exists  $\widetilde{\varphi}_T \in L^2_+(\mathbb{R}^d)$  and a subsequence  $\{\widetilde{\varphi}_T^{n_k}\}_{k\geq 1}$  such that

$$\widetilde{\varphi}_T^{n_k} \rightarrow \widetilde{\varphi}_T$$
 weakly in  $L^2(\mathbb{R}^d)$ 

Hence,

$$\widetilde{\varphi}^{n_k}(s) \to \widetilde{\varphi}(s) \text{ in } L^2(\mathbb{R}^d) \text{ for each } s \in [0, T),$$
 (3.5)

where  $\widetilde{\varphi}$  is the solution to

$$\begin{cases} \partial_t \widetilde{\varphi} + \Delta \widetilde{\varphi} = 0 & \text{ in } (0, T) \times \mathbb{R}^d, \\ \widetilde{\varphi}(T) = \widetilde{\varphi}_T & \text{ in } \mathbb{R}^d. \end{cases}$$

By (3.4) and (3.5), it holds that

$$\|B^*\widetilde{\varphi}(\tau)\|_{\ell^2}=0.$$

This, along with (ii) in Theorem 1.1, implies that  $\tilde{\varphi} \equiv 0$  in  $[0, T] \times \mathbb{R}^d$ . Consequently,

$$\liminf_{k\to+\infty}\frac{F_{\varepsilon}^{I,\tau}(\varphi_T^{n_k})}{\|\varphi_T^{n_k}\|}\geq\varepsilon.$$

This leads to a contradiction with the uniform boundedness of  $F_{\varepsilon}^{T,\tau}(\varphi_T^{n_k})$  for all  $k \ge 1$ . Hence, the first part of this lemma follows from Proposition 3.2 immediately.

For the second part of this lemma, we first note that

$$\lim_{\rho \to 0} \frac{F_{\varepsilon}^{T,\tau}(\widehat{\varphi}_{T}^{\varepsilon} + \rho \varphi_{T}) - F_{\varepsilon}^{T,\tau}(\widehat{\varphi}_{T}^{\varepsilon})}{\rho} \ge 0$$
(3.6)

for all  $\rho > 0$  and  $\varphi_T \in L^2_+(\mathbb{R}^d)$ . If  $\widehat{\varphi}_T^{\varepsilon} = 0$ , then (3.3) is obviously valid by (3.6). Otherwise, by the definition of  $F_{\varepsilon}^{T,\tau}$ , one can easily derive that

$$\lim_{\rho \to 0} \frac{F_{\varepsilon}^{T,\tau}(\widehat{\varphi}_{T}^{\varepsilon} + \rho \varphi_{T}) - F_{\varepsilon}^{T,\tau}(\widehat{\varphi}_{T}^{\varepsilon})}{\rho} = \langle B^{*} \widehat{\varphi}^{\varepsilon}(\tau), B^{*} \varphi(\tau) \rangle_{\ell^{2}} + \langle y_{0}, \varphi(0) \rangle + \varepsilon \langle \frac{\widehat{\varphi}_{T}^{\varepsilon}}{||\widehat{\varphi}_{T}^{\varepsilon}||}, \varphi_{T} \rangle \quad (3.7)$$

for any  $\varphi_T \in L^2_+(\mathbb{R}^d)$ . Noting that

$$\langle \frac{\widehat{\varphi}_T^{\varepsilon}}{\|\widehat{\varphi}_T^{\varepsilon}\|}, \varphi_T \rangle \leq \|\varphi_T\|,$$

therefore, (3.3) follows from (3.6) and (3.7). It completes the proof.

**Lemma 3.2.** For each  $\varepsilon > 0$ , let  $\widehat{\varphi}_T^{\varepsilon}$  be the minimizer of  $F_{\varepsilon}^{T,\tau}$  in  $L^2_+(\mathbb{R}^d)$ , and let  $\widehat{\varphi}^{\varepsilon}$  be the solution of (3.2) with  $\widehat{\varphi}_T^{\varepsilon}$ . Then there exists a positive constant *C*, independent of  $\varepsilon$ , so that

$$\|B^*\widehat{\varphi}^{\varepsilon}(\tau)\|_{\ell^2} \le Ce^{\frac{C}{T-\tau}}\|y_0\| \quad for \ all \quad \varepsilon > 0.$$
(3.8)

*Proof.* Since  $F_{\varepsilon}^{T,\tau}(\widehat{\varphi}_{T}^{\varepsilon}) \leq F_{\varepsilon}^{T,\tau}(0) = 0$ , we see that for any  $\varepsilon > 0$ 

$$\|B^*\widehat{\varphi}^{\varepsilon}(\tau)\|_{\ell^2}^2 \le 2\|y_0\|\|\widehat{\varphi}^{\varepsilon}(0)\| \le 2\|y_0\|\|\widehat{\varphi}^{\varepsilon}(\tau)\|.$$
(3.9)

Thanks to (ii) of Theorem 1.1, we have

$$\|\widehat{\varphi}^{\varepsilon}(\tau)\| \le C e^{\frac{c}{T-\tau}} \|B^* \widehat{\varphi}^{\varepsilon}(\tau)\|_{\ell^2} \quad \text{for all } \varepsilon > 0,$$

with a positive constant *C* independent of  $\varepsilon$ . This, together with (3.9), implies (3.8) at once.

**Proof of Theorem 3.1.** Arbitrarily fix  $\varphi_T \in L^2_+(\mathbb{R}^d)$ . Let  $\varphi$  be any solution of (3.2) with  $\varphi(T) = \varphi_T$ . Multiplying the equation (3.1) by  $\varphi$  and then integrating the resulting over  $[0, T] \times \mathbb{R}^d$ , we obtain that for any  $v \in \ell^2(\mathbb{R}^d)$ 

$$\langle y(T; v), \varphi_T \rangle - \langle y_0, \varphi(0) \rangle = \langle v, B^* \varphi(\tau) \rangle_{\ell^2}.$$
 (3.10)

Then, for any  $\varepsilon > 0$ , we let  $\widehat{\varphi}_T^{\varepsilon}$  be the minimizer of  $F_{\varepsilon}^{T,\tau}$  and let  $\widehat{\varphi}^{\varepsilon}$  be the solution of (3.2) with  $\varphi(T) = \widehat{\varphi}_T^{\varepsilon}$ . Now, by setting

$$v_{\varepsilon} := B^* \widehat{\varphi}_{\varepsilon}(\tau), \quad \varepsilon > 0,$$

in the above identity (3.10), we get

$$\langle y(T; v_{\varepsilon}), \varphi_T \rangle - \langle y_0, \varphi(0) \rangle = \langle B^* \widehat{\varphi}_{\varepsilon}(\tau), B^* \varphi(\tau) \rangle_{\ell^2}.$$

This, combined with Lemma 3.1, indicates that

$$\langle y(T; v_{\varepsilon}), \varphi_T \rangle + \varepsilon ||\varphi_T|| \ge 0 \quad \text{for any} \quad \varepsilon > 0.$$
 (3.11)

Finally, by Lemma 3.2, we see that  $v_{\varepsilon}$  is uniformly bounded in  $\ell^2(\mathbb{R}^d)$ , and thus there exists  $\hat{v} \in \ell^2(\mathbb{R}^d)$  satisfying

$$\|\hat{v}\|_{\ell^2} \le C e^{\frac{C}{T-\tau}} \|y_0\|,$$

where the constant C > 0 is independent of  $\varepsilon$ , such that (up to a subsequence)

 $v_{\varepsilon} \rightarrow \hat{v}$  weakly in  $\ell^2(\mathbb{R}^d)$  as  $\varepsilon \rightarrow 0$ .

Hence, by letting  $\varepsilon$  goes to zero in (3.11), we at once obtain that

$$\langle y(T; \hat{v}), \varphi_T \rangle \ge 0.$$

This completes the proof because of the arbitrariness of  $\varphi_T$  in  $L^2_+(\mathbb{R}^d)$ .

### References

## References

- M. Egidi, I. Veselić, Sharp geometric condition for null-controllability of the heat equation on ℝ<sup>d</sup> and consistent estimates on the control cost, Arch. Math., 111 (2018), 85–99.
- [2] L. C. Evans, Partial Differential Equations. Second edition. American Mathematical Society, Provindence, Rhode Island, 2010.
- [3] X. Liu, H. Gao, Controllability of a class of Newtonian filtration equations with control and state constraints, SIAM J. Control Optim., 46 (2007), 2256–2279.
- [4] B. Simon, Schrödinger semigroups, Bull. Amer. Math. Soc., 7 (1982), 447–526.
- [5] M. Struwe, Variational Methods. Applications to Nonlinear Partial Differential Equations and Hamiltonian Systems. Fourth edition. Springer-Verlag, Berlin, 2008.
- [6] G. Wang, M. Wang, C. Zhang, Y. Zhang, Observable set, observability, interpolation inequality and spectral inequality for the heat equation in  $\mathbb{R}^n$ , arXiv:1711.04279, to appear in J. Math. Pure. Appl.
- [7] G. Wang, M. Wang, Y. Zhang, Asymptotic observability identity for the heat equation in R<sup>d</sup>, arXiv:1810.10849v1.
- [8] H. Yu, Controllability of the heat equation with a control acting on a measurable set, Chin. Ann. Math. Ser. B, 33 (2012), 149–160.
- [9] Q. Zheng, X. Yao, Higher-order Kato class potentials for Schrödinger operators, Bull. London Math. Soc., 41 (2009), 293–301.