On a modified Cahn-Hilliard-Brinkman model with chemotaxis and nonlinear sensitivity

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Abstract

We consider an evolutionary PDE system coupling the Cahn-Hilliard equation with singular potential, mass source and transport effects, to a Brinkman-type relation for the macroscopic velocity field and to a further equation describing the evolution of the concentration of a chemical substance affecting the phase separation process. The main application we have in mind refers to tumor growth models: in particular, the equation for the chemical prescribes that such a substance tends to migrate towards the regions where the tumor cells are more dense and consume it more actively. The cross-diffusion effects characterizing the system are similar to those occurring in the Keller-Segel model for chemotaxis. There is, however, a profound difference between the two settings: actually, the Cahn-Hilliard system prescribes a fourth-order dynamics with respect to space variables, whereas most models for chemotaxis are of the second order in space. This fact has a number of specific consequences regarding regularity properties of solutions and conditions ensuring existence. Our main results are devoted to proving existence of weak solutions in the case when the chemotactic sensitivity function depends nonlinearly on the chemical species concentration, and, more precisely, has a slow growth at infinity so to avoid finite-time blowup. We also analyze the asymptotic problem obtained by letting the viscosity go to zero so to get a Darcy flow regime in the limit.

Key words: Cahn-Hilliard, chemotaxis, mass source, singular potential, nonlinear sensitivity. AMS (MOS) subject classification: 35D30, 35K35, 35K86, 35Q92, 92C17, 92C50.

1 Introduction

4

Letting $\Omega \subset \mathbb{R}^d$, $d \in \{2,3\}$, be a smooth bounded domain of boundary Γ , we consider the following PDE system in $\Omega \times (0,T)$, T > 0 being an arbitrary, but otherwise fixed, final time of the evolution process:

$$\varphi_t + \boldsymbol{u} \cdot \nabla \varphi - \Delta \mu = h(\sigma, \varphi) - \ell \varphi,$$
(1.1)

$$\iota = -\Delta\varphi + f(\varphi) - \chi\sigma,\tag{1.2}$$

$$\sigma_t + \boldsymbol{u} \cdot \nabla \sigma - \Delta \sigma + \chi \operatorname{div}(\alpha(\sigma) \nabla \varphi) = b(\sigma, \varphi), \tag{1.3}$$

$$-\varepsilon \operatorname{div}(D\boldsymbol{u}) + \boldsymbol{u} = \nabla \pi + \mu \nabla \varphi - \chi \varphi \nabla \sigma, \qquad (1.4)$$

$$\operatorname{div} \boldsymbol{u} = 0. \tag{1.5}$$

Relations (1.1)-(1.2) constitute a suitable form of the Cahn-Hilliard-Oono system [25, 26] for the order parameter φ , where the right-hand side of (1.1) plays the role of a mass source decomposed as the sum of a (dominating) linear part $-\ell\varphi$, where $\ell > 0$ is a given constant, and a nonlinear perturbation $h(\sigma, \varphi)$. Convection effects are considered, with the macroscopic flow velocity \boldsymbol{u} satisfying, in an incompressible regime, the Brinkman-type relation (1.4), π denoting the pressure. In (1.4), the expression Du stands for the symmetrized gradient of u, with the last two terms on the right-hand side playing the role of Korteweg forces whose specific expression guarantees the validity of the energy balance law. We assume, in principle, $\varepsilon > 0$, but (in the spirit, e.g., of [22]), we will also study the behavior as $\varepsilon \searrow 0$ corresponding to a Darcy regime in the limit.

The nonlinear function f is assumed as the derivative (or, more precisely, the subdifferential) of a logarithmic potential F of Flory-Huggins type, whose expression reads

$$F(r) = (1+r)\log(1+r) + (1-r)\log(1-r) - \frac{\lambda}{2}r^2, \quad r \in [-1,1], \quad \lambda \ge 0.$$
(1.6)

In particular, as is customary for phase-separation models, the order parameter is normalized so that the pure configurations are represented by $\varphi = \pm 1$; correspondingly, the minima of F are attained in proximity of these states, and they are deeper for larger λ .

The main application we have in mind for system (1.1)-(1.5) refers to tumor growth processes; in this setting φ may represent the local proportion of active cancer cells; hence the right-hand side of (1.1) drives the evolution of the total tumor mass. Indeed, assuming suitable boundary conditions (and, in particular, no-flux boundary conditions for the *chemical potential* μ , which acts as an auxiliary variable), the mass evolution law is obtained by integrating (1.1) over Ω . This process is influenced by the additional variable σ , which represents the concentration of a chemical substance, like a nutrient or a drug. As in related models, σ is assumed to satisfy the second order evolutionary equation (1.3), whose specific formulation can be considered as the main novelty of the present contribution, so deserving some words of explanation.

First of all, we point out that tumor growth models based on the Cahn-Hilliard "diffuseinterface" description are becoming increasingly popular among the scientific community: for mathematical results and a more extensive physical background we may quote, with no claim of exhaustivity, the recent papers [8, 12, 13, 14, 15, 19, 37], the monograph [7], and the references therein. In most of the quoted works, actually, a different expression for the nutrient equation is assumed, which may be written, in the simplest case, as

$$\sigma_t - \Delta \sigma + \chi \Delta \varphi = b(\sigma, \varphi). \tag{1.7}$$

Here, for the sake of clarity, we have taken constant mobility coefficients and omitted the transport effects described by the macroscopic velocity. Both in (1.7) and in (1.3), the coefficient $\chi > 0$ can be interpreted as a transport parameter which drives the cells in the direction of the regions where σ takes larger values, basically providing a chemotactic response of the cells with respect to the nutrient, an effect which is observed in real world situations. The main difference between (1.7) and (1.3) lies in the expression of the cross-diffusion term, which in (1.7) depends linearly on φ . This, in particular, implies that σ may not satisfy a minimum principle, even in the case when the right-hand side source term $b(\sigma, \varphi)$ is properly designed (i.e. it satisfies suitable sign and growth conditions).

Since the nonnegativity of σ is an expected property as σ plays the role of a concentration, it was proposed in [27] to replace (1.7) with

$$\sigma_t - \Delta \sigma + \chi \operatorname{div}(\sigma \nabla \varphi) = b(\sigma, \varphi). \tag{1.8}$$

Actually, such an expression preserves the nonnegativity of σ and is reminiscent of the Keller-Segel model for chemotaxis, so providing a new type of coupling between the Cahn-Hilliard equation and a Keller-Segel-like relation (we observe that a different connection between the two models was analyzed in [9]). One of the main notable features of the model introduced in [27] stands in the fact that, differently from the "genuine" Keller-Segel case, the coercivity of the energy functional is tied to the uniform boundedness of φ , hence on the choice of a singular configuration potential like (1.6). It is clear, however, that the quadratic behavior of the cross-diffusion term in (1.8) may still give rise to blow-up effects; in order to prevent the latter, in [27] b is assumed to have a logistic growth with respect to σ . This choice, which is rather common in the Keller-Segel setting, is aimed at penalizing the large values of the concentration; mathematically speaking, it helps for the sake of obtaining global in time estimates.

In this paper, considering system (1.1)-(1.5), we will see that, suitably designing the expression of the function α in (1.3), we may simultaneously keep the minimum principle for σ and avoid finite-time blow up under very general and natural conditions (and, in particular, with no need for a logistic

growth of b). To be precise, we will assume α to degenerate at 0 as fast as σ , and to behave at infinity like a suitable power σ^a , where $a \in [0, 1)$ is assumed to be small enough in order to avoid blow up. This choice, usually noted as "degenerate sensitivity" in the Keller-Segel literature, is very popular in the chemotaxis community: one can refer to the seminal paper [20] (see also [35] and the references therein for further results and a survey of the more recent literature). Not surprisingly, as happens in similar models characterized by a degenerate sensitivity, our main results, which are devoted to proving existence and regularity of weak solutions for system (1.1)-(1.5), will require suitable restrictions on the exponent a, also depending on the space dimension.

In order to explain the "spirit" of our results and to describe the novelties and the approach of the present work, it is first worth outlining some specific features of system (1.1)-(1.3) compared to other models characterized by similar cross-diffusion terms. In particular, we may observe the following facts:

- The (fourth order) Cahn-Hilliard coupling provides lower time regularity but better space regularity for φ ; this effect results in a somehow different behavior of the cross-diffusion term compared to the Keller-Segel model;
- The choice of the singular potential (1.6) guarantees "for free" the uniform boundedness of φ . This fact not only, as already observed, provides coercivity of the energy, but it also guarantees additional regularity of the coupling terms; in particular, the uniform boundedness of φ is a fundamental information as we seek for additional regularity estimates;
- The Cahn-Hilliard system, which is a subset of (1.1)-(1.5), somehow drives the optimal strategy for getting a-priori estimates. Not surprisingly, the argument used for proving existence is primarily based on the energy balance, which holds as a direct consequence of the variational structure of the system. Then, our assumption $\alpha(\sigma) \sim \sigma^a$ for large σ corresponds to the regularity condition $\sigma \in L^p(\Omega)$, where p = 2 - a; namely, a slower growth of α at infinity corresponds to a greater summability of solutions. Correspondingly, for large a (or, equivalently, small p), the outcome of the energy estimate in terms of a-priori regularity seems not sufficient to pass to the limit in the approximation, with the main difficulties arising in connection with the product terms and, more specifically, the transport and Korteweg terms depending on σ ;
- Looking for additional a-priori regularity, we derive further estimates of the so-called "entropy" type, according to the terminology currently in use for fourth order evolutionary systems. Basically, this consists in testing (1.2) by $-\Delta\varphi$ and (1.3) by σ^{q-1} for suitable $q \in [p, \infty)$. Actually, the information resulting from this type of procedure provides separate regularity bounds for the diffusion terms in (1.3) (and not just for their sum). Moreover, it also yields a L^2 -estimate for $\Delta\varphi$ as well as some additional summability of σ ;
- Our main result will be devoted to the three-dimensional case: we will show that, assuming that $p = 2 a \in (12/11, 2]$, if the initial datum σ_0 lies in $L^q(\Omega)$ for some $q \in [p, 2]$, then σ keeps staying in $L^q(\Omega)$ uniformly on (0, T). In dimension d = 2, with a simpler proof we will show that similar properties hold for any $p \in (1, 2]$. However, the two-dimensional case will be analyzed more carefully in a forthcoming paper, where we will also discuss the limit case p = 1. We finally observe that the (supposedly) critical exponent p = 12/11 for d = 3 seems less restrictive compared to the genuine Keller-Segel setting (cf., e.g., [36]);
- Considering the presence of (mass and nutrient) source terms, we do not expect eventual uniform boundedness of solutions. Indeed, even at the level of the energy estimate, our procedure relies in an essential way on Grönwall's lemma, implying that solutions may grow exponentially in suitable norms. The problem of eventual boundedness of solutions (in the spirit of [20] and several other papers dealing with the Keller-Segel model) as well as other questions related with the long-time behavior (e.g., dissipativity, existence of attractors) may be addressed in some future work referring to the case with no external sources, and possibly neglecting the effects of the macroscopic velocity;
- We finally remark that, in view of the highly nonlinear character of the model, uniqueness in the class of weak solutions is not expected to hold. In particular, the combination of a nonlinear mass

source term in (1.1) with the singular potential (1.6), as observed in [11, 18], may prevent using a contractive argument; in addition to that, the transport effects, especially in the Darcy limit regime, may be difficult to manage. Hence, also the question of uniqueness may be considered in a future work by assuming a simplified setting.

The plan of the paper is as follows: in the next section, we detail our assumptions on coefficients and data and state our mathematical results. In the subsequent Section 3, we prove the basic "physical" a priori estimates resulting from the mass and energy conservation principles and representing the fundamental step in the proof of existence. In Section 4 we obtain further estimates of the so-called "entropy" type which are crucial in order to deal with the cross-diffusion term in (1.3) and to provide some further a-priori information on weak solutions. In Section 5 we show weak sequential stability of families of weak solutions, namely we prove that any sequence of weak solutions complying with the a-priori estimates uniformly with respect to approximation or regularization parameters admits at least one limit point which is still a weak solution. Finally, in Section 6 we propose a specific regularization scheme and show its compatibility both with the estimates and with the argument used to pass to the limit.

2 Assumptions and main results

We start with introducing a set of notation which will be useful in order to rigorously formulate our mathematical results. Letting Ω be a smooth bounded domain of \mathbb{R}^d , $d \in \{2,3\}$, with boundary Γ , we set $H := L^2(\Omega)$ and $V := H^1(\Omega)$. We will generally write H in place of H^d (with similar notation for other spaces), whenever vector-valued functions are considered. We denote by (\cdot, \cdot) the standard scalar product of H and by $\|\cdot\|$ the associated Hilbert norm. Moreover, we equip V with the usual norm $\|\cdot\|_V^2 = \|\cdot\|^2 + \|\nabla\cdot\|^2$. Identifying H with its dual space H' by means of the scalar product introduced above, we obtain the chain of continuous and dense embeddings $V \subset H \subset V'$. We will indicate by $\langle \cdot, \cdot \rangle$ the duality pairing between V' and V, or, more generally, between X' and X, where X is a generic Banach space continuously and densely embedded into H.

Next, we assume F be given by (1.6), we indicate as f the derivative (or, more precisely, the subdifferential) of F, and we denote as $\beta(r) := f(r) + \lambda r$ the "monotone part" of f, namely

$$\beta(r) = \log(1+r) - \log(1-r), \quad r \in (-1,1).$$
(2.1)

We also set

$$\boldsymbol{V} := \left\{ \boldsymbol{a} \in V^d, \quad \text{div}\, \boldsymbol{a} \equiv 0 \text{ in } \Omega, \quad \boldsymbol{a} \cdot \boldsymbol{n} = 0 \text{ on } \Gamma \right\},\tag{2.2}$$

with \boldsymbol{n} denoting the outer normal unit vector to Γ . We also let $H^2_{\boldsymbol{n}}(\Omega)$ be the (closed) subspace of $H^2(\Omega)$ containing the functions with zero normal derivative on Γ . For $\boldsymbol{v} \in \boldsymbol{V}$, we denote as $D\boldsymbol{v} = \frac{\nabla \boldsymbol{v} + (\nabla \boldsymbol{v})^t}{2}$ the symmetrized gradient of \boldsymbol{v} . As in [17] (see also [10] for a "multi-phase" version of this condition), we assume that

$$h \in C^1(\mathbb{R} \times \mathbb{R}), \qquad |h(\sigma, \varphi)| \le H \quad \text{for every } (\sigma, \varphi) \in [0, \infty) \times \mathbb{R},$$

$$(2.3)$$

for a suitable constant H > 0 satisfying the compatibility condition

$$\frac{H}{\ell} < 1. \tag{2.4}$$

Note that the above implies in particular that, uniformly with respect to σ , $h(\sigma, \varphi) - \ell \varphi$ is strictly negative (respectively, strictly positive) in a left neighbourhood of $\varphi = 1$ (respectively, in a right neighbourhood of $\varphi = -1$). In addition, we ask that there exist $C_{h,1}, C_{h,2} > 0$ such that the partial derivatives of h satisfy

$$|h_{\varphi}(\sigma,\varphi)| \le C_{h,1}, \text{ and } |h_{\sigma}(\sigma,\varphi)| \le C_{h,2} \frac{1}{1+|\sigma|}, \text{ for every } (\sigma,\varphi) \in [0,\infty) \times \mathbb{R}.$$
 (2.5)

The second condition, which has a mainly technical character, corresponds to asking that, if the concentration σ is very large, its effects on the total mass tend to become essentially independent

of variations of σ . Actually, this assumption, which seems reasonable from the physical viewpoint, is probably not optimal mathematically: the proof could be likely adaptable to the case when h_{σ} behaves at infinity as σ^{-k} for some suitable $k \in (0, 1)$ depending on p; however, since this does not appear to be an essential point, we decided to assume a simpler condition so to reduce technicalities. We also observe that, in (2.3), (2.5) (and elsewhere below) we considered the possibility that $|\varphi| > 1$: actually, while "in the limit" we will have $|\varphi| \leq 1$ almost everywhere due to the occurrence of the singular potential, this may not be the case in an approximation as F is smoothed out (see Section 6 below for details). On the other hand, the assumptions are only stated for $\sigma \in [0, \infty)$ as it is clear that also in the approximation the minimum principle for σ is satisfied.

As noted in the introduction, we assume that $\alpha(\sigma) \sim \sigma$ near zero and $\alpha(\sigma) \sim \sigma^a$, $a \in [0, 1)$ near infinity (here and below, if a = 0, of course, we interpret $\sigma^a \equiv 1$ for $\sigma \geq 0$). We also set p = 2-a, so that $p \in (1, 2]$, in order to emphasize the L^p -summability provided by the above ansatz. Then, using p, our assumption on α can be stated as

$$\alpha(s) = \frac{s}{1+s^{p-1}}, \quad \text{for } s \ge 0.$$
 (2.6)

Consequently, we have

$$\frac{1}{\alpha(s)} = \frac{1+s^{p-1}}{s} = \frac{1}{s} + s^{p-2}, \quad \text{for } s > 0.$$
(2.7)

We can then define

$$\gamma(s) := \int_{1}^{s} \frac{\mathrm{d}r}{\alpha(r)} = \ln s + \frac{1}{p-1} s^{p-1} - \frac{1}{p-1}, \quad \text{for } s > 0.$$
(2.8)

With the above notation at hand, equation (1.3) can be more conveniently restated as

$$\sigma_t + \boldsymbol{u} \cdot \nabla \sigma - \operatorname{div} \left(\alpha(\sigma) \nabla(\gamma(\sigma) - \chi \varphi) \right) = b(\sigma, \varphi).$$
(2.9)

Next, we assume the nutrient source term to satisfy

$$b \in \operatorname{Lip}_{\operatorname{loc}}([0,\infty) \times \mathbb{R};\mathbb{R}), \quad -b_0 \sigma \le b(\sigma,\varphi) \le b_\infty(1+\sigma), \quad \text{for every } (\sigma,\varphi) \in [0,\infty) \times \mathbb{R}, \quad (2.10)$$

and for some constants $b_0, b_{\infty} > 0$. Note that the above condition is designed so to ensure the applicability of a minimum principle argument for σ uniformly with respect to φ .

System (1.1)-(1.5) is complemented with the Cauchy conditions

$$\varphi|_{t=0} = \varphi_0, \qquad \sigma|_{t=0} = \sigma_0, \quad \text{in } \Omega, \tag{2.11}$$

with the initial data satisfying at least the following properties:

$$\varphi_0 \in V, \quad F(\varphi_0) \in L^1(\Omega), \quad m_0 := (\varphi_0)_\Omega \in (-1, 1),$$
(2.12)

$$\sigma_0 \in L^p(\Omega), \quad \sigma_0 > 0 \text{ a.e. in } \Omega, \quad \ln \sigma_0 \in L^1(\Omega).$$
 (2.13)

The above conditions correspond to the finiteness of the initial energy and to the fact that, as is customary in the setting of the Cahn-Hilliard system with logarithmic potential (1.6), we cannot admit initial configurations where φ is almost everywhere equal to 1 (or to -1). Here and below, for a generic function v defined of Ω we have denoted by v_{Ω} the spatial mean of v, namely

$$v_{\Omega} := \frac{1}{|\Omega|} \int_{\Omega} v, \qquad (2.14)$$

 $|\Omega|$ denoting the *d*-dimensional Lebesgue measure of Ω .

Moreover we shall assume no-flux (i.e., homogeneous Neumann) boundary conditions for μ , φ and σ , and complete slip boundary conditions for the velocity, namely

$$\partial_{\boldsymbol{n}}\mu = \partial_{\boldsymbol{n}}\varphi = \partial_{\boldsymbol{n}}\sigma = 0, \quad \text{on } \Gamma \times (0, T),$$

$$(2.15)$$

$$\boldsymbol{u} \cdot \boldsymbol{n} = 0, \quad -\varepsilon((D\boldsymbol{u})\boldsymbol{n})_{\tau} = 0, \quad \text{on } \Gamma \times (0,T).$$
 (2.16)

Here and below, for a (sufficiently smooth) vector field a defined on Γ , we have denoted by $a_{\tau} := a - (a \cdot n)n$ the tangential component of a.

With the above material at hand, we can now present our main results, which are devoted to proving existence of global in time weak solutions to the Cahn-Hilliard-Brinkman-nutrient model as well as to studying the asymptotic behavior of solutions as $\varepsilon \searrow 0$. As observed in the introduction, we will consider separately the cases d = 2 and d = 3. In the two-dimensional case we will prove the following simpler (but likely not optimal) statement:

Theorem 2.1. Let the assumptions (1.6), (2.3)-(2.5), (2.6), (2.10), (2.12)-(2.13) hold and let d = 2. Let $\varepsilon > 0$, and assume that p > 1. Moreover, let also

$$\sigma_0 \in H. \tag{2.17}$$

Then, there exists at least one quadruple $(\varphi, \mu, \sigma, u)$ with the regularity properties

$$\varphi \in H^1(0,T;V') \cap L^{\infty}(0,T;V) \cap L^4(0,T;H^2(\Omega)) \cap L^2(0,T;W^{2,P}(\Omega)),$$
(2.18)

$$F(\varphi) \in L^{\infty}(0,T;L^{1}(\Omega)), \qquad f(\varphi) \in L^{2}(0,T;L^{P}(\Omega)),$$
(2.19)

$$\mu \in L^2(0,T;V),$$
 (2.20)

$$\boldsymbol{u} \in L^2(0,T;\boldsymbol{V}), \tag{2.21}$$

$$\sigma \in H^1(0,T; W^{1,R}(\Omega)') \cap L^{\infty}(0,T;H) \cap L^2(0,T;V),$$
(2.22)

$$\sigma > 0$$
 a.e. in $\Omega \times (0,T)$, $\ln \sigma \in L^2(0,T;V)$, (2.23)

$$\boldsymbol{H}(\sigma,\varphi) := \alpha^{1/2}(\sigma)\nabla(\gamma(\sigma) - \chi\varphi) \in L^2(0,T;H),$$
(2.24)

where the exponents P in (2.18)-(2.19) and R in (2.22) may be arbitrarily taken in $[2, \infty)$ and in $(2, \infty]$, respectively. The quadruple $(\varphi, \mu, \sigma, u)$ satisfies, a.e. in (0, T), the following weak version of system (1.1)-(1.5):

$$\langle \varphi_t, \xi \rangle - \int_{\Omega} \varphi \boldsymbol{u} \cdot \nabla \xi + \int_{\Omega} \nabla \mu \cdot \nabla \xi = \int_{\Omega} \left(h(\sigma, \varphi) - \ell \varphi \right) \xi \quad \text{for every } \xi \in V,$$
(2.25)

$$\mu = -\Delta \varphi + f(\varphi) - \chi \sigma, \quad \text{a.e. in } \Omega, \tag{2.26}$$

$$\langle \sigma_t, \eta \rangle - \int_{\Omega} \sigma \boldsymbol{u} \cdot \nabla \eta + \int_{\Omega} \alpha^{1/2}(\sigma) \boldsymbol{H}(\sigma, \varphi) \cdot \nabla \eta = \int_{\Omega} b(\sigma, \varphi) \eta, \quad \text{for all } \eta \in W^{1,R}(\Omega),$$
(2.27)

$$\varepsilon \int_{\Omega} (D\boldsymbol{u}) : (D\boldsymbol{\xi}) + \int_{\Omega} \boldsymbol{u} \cdot \boldsymbol{\xi} = -\int_{\Omega} \varphi \boldsymbol{\xi} \cdot \nabla \mu + \chi \int_{\Omega} \sigma \boldsymbol{\xi} \cdot \nabla \varphi \quad \text{for every } \boldsymbol{\xi} \in \boldsymbol{V}.$$
(2.28)

Moreover, there hold the initial conditions (2.11) as well as the boundary condition

$$\partial_{\boldsymbol{n}}\varphi = 0 \quad \text{a.e. on } \Gamma \times (0,T).$$
 (2.29)

Finally, let, for $\varepsilon \in (0,1)$, $(\varphi_{\varepsilon}, \mu_{\varepsilon}, \sigma_{\varepsilon}, u_{\varepsilon})$ be a family of solutions satisfying the system in the sense specified above. Then, there exist a (nonrelabelled) subsequence of $\varepsilon \searrow 0$ and a limit quadruple $(\varphi, \mu, \sigma, u)$ such that

$$\varphi_{\varepsilon} \to \varphi \quad \text{weakly star in } H^1(0,T;V') \cap L^{\infty}(0,T;V) \cap L^4(0,T;H^2(\Omega)), \tag{2.30}$$

$$\mu_{\varepsilon} \to \mu \quad \text{weakly in } L^2(0,T;V),$$
(2.31)

$$\sigma_{\varepsilon} \to \sigma \quad \text{weakly star in } W^{1,4/3}(0,T;W^{1,4}(\Omega)') \cap L^{\infty}(0,T;H) \cap L^{2}(0,T;V), \tag{2.32}$$

$$\boldsymbol{u}_{\varepsilon} \to \boldsymbol{u} \quad \text{weakly in } L^2(0,T;H).$$
 (2.33)

The quadruple $(\varphi, \mu, \sigma, \boldsymbol{u})$ satisfies (2.25)-(2.28) with $\varepsilon = 0$ (where, now, R = 4, i.e., η has to lie in $W^{1,4}(\Omega)$, in (2.27)), the additional boundary condition (2.29), and the initial conditions (2.11).

Remark 2.2. In the Darcy limit $\varepsilon = 0$, the velocity \boldsymbol{u} just acts as an auxiliary variable and the system could be restated directly in terms of the pressure π (or, more precisely, of its gradient as π is determined up to an additive constant). Actually, π turns out to satisfy the limit elliptic system

$$-\Delta \pi = \operatorname{div}(\mu \nabla \varphi - \chi \varphi \nabla \sigma), \qquad (2.34)$$

complemented with the natural (no-flux) boundary conditions resulting from a comparison with (2.15). Once π is given by (2.34), (1.4) with $\varepsilon = 0$ can then be seen as a *definition* of \boldsymbol{u} in terms of the other variables.

We now move to our main statement referring to the three-dimensional setting. In this case, we will first prove that, if p is strictly larger than the (supposedly) critical exponent 12/11, then the two summands in the cross-diffusion term $H(\sigma, \varphi)$ can be controlled separately; this information, which is achieved by means of the so-called "entropy" method for fourth order PDE's, is crucial for the sake of proving existence of weak solutions. Furthermore, using a more refined version of the entropy estimate, we can also show that, if $\sigma_0 \in L^q(\Omega)$ for $q \in [p, 2]$, then a natural parabolic estimate for $\sigma^{q/2}$ holds (see (2.42) below). Finally, if q > 6/5, we can take the limit $\varepsilon \to 0$ and prove convergence to the Darcy flow. This further restriction is motivated by the need for having sufficient regularity to control, uniformly as $\varepsilon \searrow 0$, the transport and Korteweg terms depending on σ . These results are summarized in the following statement:

Theorem 2.3. Let the assumptions (1.6), (2.3)-(2.5), (2.6), (2.10), (2.12)-(2.13) hold, and let d = 3. Let, moreover, $\varepsilon > 0$ and let

$$p \in (12/11, 2]. \tag{2.35}$$

In addition to that, let us additionally assume

$$\sigma_0 \in L^q(\Omega) \quad \text{where } q \in [p, 2]. \tag{2.36}$$

Then, there exists at least one quadruple $(\varphi, \mu, \sigma, u)$ with the regularity properties

$$\varphi \in H^1(0,T;V') \cap L^{\infty}(0,T;V) \cap L^{P_0}(0,T;H^2(\Omega)) \cap L^2(0,T,W^{2,\frac{3q}{3-q}}(\Omega)),$$
(2.37)

$$F(\varphi) \in L^{\infty}(0,T;L^{1}(\Omega)), \qquad f(\varphi) \in L^{2}(0,T;L^{\frac{3q}{3-q}}(\Omega)),$$
(2.38)

$$\mu \in L^2(0,T;V), \tag{2.39}$$

$$\boldsymbol{u} \in L^2(0,T;\boldsymbol{V}),\tag{2.40}$$

$$\sigma \in W^{1,Z}(0,T; W^{1,R}(\Omega)') \cap L^2(0,T; W^{1,S}(\Omega)) \quad \text{for some } Z > 1,$$
(2.41)

$$\sigma^{q/2} \in L^{\infty}(0,T;H) \cap L^{2}(0,T;V), \tag{2.42}$$

$$\sigma > 0 \quad \text{a.e. in } \Omega \times (0,T), \qquad \ln \sigma \in L^2(0,T;V), \tag{2.43}$$

$$\boldsymbol{H}(\sigma,\varphi) := \alpha^{1/2}(\sigma)\nabla(\gamma(\sigma) - \chi\varphi) \in L^2(0,T;H),$$
(2.44)

where the exponents P_0, R, S in (2.37), (2.41) satisfy (note that S > 1 due to (2.35))

$$P_0 = \min\left\{\frac{18q - 6p}{12 - 5p}, 4\right\}, \qquad S = \min\left\{\frac{6p}{12 - 5p}, p\right\}, \qquad R = \max\left\{4, p'\right\}, \tag{2.45}$$

p' being the conjugate exponent to p. The quadruple $(\varphi, \mu, \sigma, \mathbf{u})$ satisfies, a.e. in (0, T), relations (2.25)-(2.28), where, now, the exponent R in (2.27) is specified by (2.45). Moreover, there hold the boundary condition (2.29) and, in the sense of traces, the initial conditions (2.11). Next, let us additionally assume

$$q > 6/5 \tag{2.46}$$

and let, for $\varepsilon \in (0,1)$, $(\varphi_{\varepsilon}, \mu_{\varepsilon}, \sigma_{\varepsilon}, u_{\varepsilon})$ be a family of solutions satisfying the system in the sense specified above. Then, there exist a (nonrelabelled) subsequence of $\varepsilon \searrow 0$ and a limit quadruple $(\varphi, \mu, \sigma, u)$ such that

$$\varphi_{\varepsilon} \to \varphi \quad \text{weakly in } H^1(0,T;V') \cap L^{P_0}(0,T;H^2(\Omega)) \cap L^2(0,T,W^{2,\frac{3q}{3-q}}(\Omega)), \tag{2.47}$$

$$\varphi_{\varepsilon} \to \varphi$$
 weakly star in $L^{\infty}(0, T; V)$. (2.48)

$$\mu_{\varepsilon} \to \mu \quad \text{weakly in } L^2(0,T;V),$$
(2.49)

$$\sigma_{\varepsilon} \to \sigma \quad \text{weakly in } W^{1,Z}(0,T;W^{1,Z'}(\Omega)') \cap L^2(0,T;W^{1,S}(\Omega)), \quad \text{for some } Z > 1, \tag{2.50}$$

$$\boldsymbol{u}_{\varepsilon} \to \boldsymbol{u} \quad \text{weakly in } L^2(0,T;H).$$
 (2.51)

The quadruple $(\varphi, \mu, \sigma, \mathbf{u})$ satisfies (2.25)-(2.28) with $\varepsilon = 0$ (where, now, in (2.27), η has to lie in $W^{1,Z'}(\Omega)$, Z' being the "large" conjugate exponent to Z), the additional boundary condition (2.29), and the initial conditions (2.11).

We conclude this section with some remarks aimed at clarifying some aspects of the above statements.

Remark 2.4. It is worth observing that the boundary conditions for μ and u are embedded, respectively, in the weak formulation (2.25) and in the regularity (2.21) (or (2.40)) combined with relation (2.28). We also remark that relation (2.27) might be more standardly rewritten as

$$\langle \sigma_t, \eta \rangle - \int_{\Omega} \sigma \boldsymbol{u} \cdot \nabla \eta + \int_{\Omega} \alpha(\sigma) \nabla(\gamma(\sigma) - \chi \varphi) \cdot \nabla \eta = \int_{\Omega} b(\sigma, \varphi) \eta.$$
(2.52)

We preferred to keep formulation (2.27) since it better emphasizes the additional regularity (2.24) satisfied by the mixed term \boldsymbol{H} . In particular we point out that, while \boldsymbol{H} , as a sum, lies in L^2 , as it is decomposed into the summands $\alpha^{1/2}(\sigma)\nabla\gamma(\sigma)$ and $-\chi\alpha^{1/2}(\sigma)\nabla\varphi$, these summands may fulfill lower regularity properties.

Remark 2.5. Referring to the statement of Theorem 2.3, some comments about assumptions (2.35) and (2.36) are in order. First of all, the condition p > 12/11 corresponds to the most general situation in which we are able to prove, via the entropy method, decoupled regularity for the two summands in the cross-diffusion term $H(\sigma, \varphi)$. In particular, one may well consider the case q = p, which corresponds to proving existence for initial data having exactly the "energy" regularity. On the other hand, looking at (2.36), one may assume that σ_0 enjoys some additional summability; in that case, we can prove that this level of summability is maintained in time (cf. (2.42)).

Remark 2.6. We decided not to treat the case when $\sigma_0 \in L^q(\Omega)$ for q strictly larger than 2. In that situation, however, it is expected that the regularity of solutions could be further improved by bootstrapping. Moreover, assuming some condition on $\nabla \sigma_0$ one may also prove some summability of σ_t and $\Delta \sigma$ (which are now controlled only in Sobolev spaces of negative order with respect to space variables). In turn, this may help improving the information on φ by means of the so-called "second energy estimate" that can be performed for the Cahn-Hilliard equation with singular potential. The question of additional regularity of solutions may be treated in a future work, possibly neglecting, also in this case, the effects of convection.

3 Mass balance and energy estimate

In this part we derive a number of a priori estimates satisfied by any hypothetical solution of our system and holding under the natural "physical" conditions on the initial data, namely finiteness of the physical energy, compatibility of the initial mass with the configuration potential, and positivity of the initial concentration. For simplicity, we will work directly on the "original" system (1.1)-(1.5): in Section 6 we will see how the procedure may be adapted to a specific approximation scheme. We start with deriving the evolution of mass.

Balance of mass. Integrating (1.1) over Ω and using the incompressibility constraint and the boundary conditions, we immediately deduce

$$\frac{\mathrm{d}}{\mathrm{d}t}\varphi_{\Omega} + \ell\varphi_{\Omega} = \frac{1}{|\Omega|} \int_{\Omega} h(\sigma, \varphi).$$
(3.1)

Applying (2.4) and using the last (2.12) with a simple argument based on the comparison principle for ODE's, we deduce that there exists a sufficiently small $\delta \in (0, 1)$, depending only on m_0 , ℓ and H(and in particular independent of T) and such that

$$-1 + \delta \le m(t) = (\varphi(t))_{\Omega} \le 1 - \delta \quad \text{for every } t \in [0, T].$$

$$(3.2)$$

Positivity of the nutrient concentration. As in similar models, this property may be obtained by means of a standard argument based on the Stampacchia truncation method combined with a regularization of equation (1.3), which is needed in order to justify the use of the truncated σ as a test function. In particular, one may refer to the regularization scheme that will be detailed in Section 6 and possibly operate a further smoothing of u and φ . Note also that, by (2.10), the right-hand side of (1.3) is suitably designed so to allow for such an argument. It is also worth observing that the spirit of the Stampacchia truncation argument is to prove

$$\varphi_0(\cdot) \ge 0$$
 a.e. in $\Omega \implies \varphi(\cdot, \cdot) \ge 0$ a.e. in $\Omega \times (0, T)$. (3.3)

However, we point out that (3.3) is not sufficient for our purposes; actually, in the sequel, based on (2.13), we will improve the above argument obtaining a stronger information.

Energy estimate. We now prove the most important a-priori bound, which is a direct consequence of the variational structure of the system. To derive it, we test (1.1) by μ , (1.2) by φ_t , and take the difference so to obtain

$$\frac{\mathrm{d}}{\mathrm{d}t} \Big(\frac{1}{2} \|\nabla\varphi\|^2 + \int_{\Omega} F(\varphi) \Big) - \chi \int_{\Omega} \sigma\varphi_t + \|\nabla\mu\|^2 = -\int_{\Omega} \mu \boldsymbol{u} \cdot \nabla\varphi + \int_{\Omega} \big(h(\sigma,\varphi) - \ell\varphi \big) \mu.$$
(3.4)

Moreover, we test (2.9) by $\gamma(\sigma) - \chi \varphi$. This gives

$$\frac{\mathrm{d}}{\mathrm{d}t} \int_{\Omega} \widehat{\gamma}(\sigma) - \chi \int_{\Omega} \sigma_t \varphi + \int_{\Omega} \alpha(\sigma) \big| \nabla(\gamma(\sigma) - \chi \varphi) \big|^2 = \int_{\Omega} b(\sigma, \varphi) (\gamma(\sigma) - \chi \varphi) + \chi \int_{\Omega} \varphi \boldsymbol{u} \cdot \nabla \sigma, \quad (3.5)$$

where the incompressibility constraint (1.5) and the first (2.16) have also been used. Note also that $\gamma(\sigma)$ contains $\ln \sigma$ a summand; as will be clear in the sequel of the argument, $\ln \sigma$ can be shown to have some a-priori regularity, which justifies its use as a test function. Here, we have also set

$$\widehat{\gamma}(s) := \int_{1}^{s} \gamma(r) \, \mathrm{d}r = \frac{1}{p(p-1)} s^{p} + s \ln s - \frac{p}{p-1} s + \frac{p+1}{p}.$$
(3.6)

Testing now (1.4) by \boldsymbol{u} , and using once more (1.5) and (2.16), we deduce

$$\varepsilon \|D\boldsymbol{u}\|^2 + \|\boldsymbol{u}\|^2 = \int_{\Omega} \mu \boldsymbol{u} \cdot \nabla \varphi - \chi \int_{\Omega} \varphi \boldsymbol{u} \cdot \nabla \sigma.$$
(3.7)

Summing (3.4), (3.5) and (3.7), we obtain

$$\frac{\mathrm{d}}{\mathrm{d}t} \left(\frac{1}{2} \|\nabla\varphi\|^2 + \int_{\Omega} F(\varphi) + \int_{\Omega} \left(\widehat{\gamma}(\sigma) - \chi\sigma\varphi \right) \right) + \|\nabla\mu\|^2 + \int_{\Omega} \alpha(\sigma) |\nabla(\gamma(\sigma) - \chi\varphi)|^2 + \varepsilon \|D\boldsymbol{u}\|^2 + \|\boldsymbol{u}\|^2 = \int_{\Omega} b(\sigma,\varphi)(\gamma(\sigma) - \chi\varphi) + \int_{\Omega} \left(h(\sigma,\varphi) - \ell\varphi \right) \mu.$$
(3.8)

Then, defining now the energy functional

$$\mathcal{E}(\varphi,\sigma) := \frac{1}{2} \|\nabla\varphi\|^2 + \int_{\Omega} F(\varphi) + \int_{\Omega} \left(\widehat{\gamma}(\sigma) - \chi\sigma\varphi\right),\tag{3.9}$$

and recalling (3.6), it is immediate to check the coercivity property

$$\mathcal{E}(\varphi,\sigma) \ge \kappa \left(\|\varphi\|_V^2 + \|F(\varphi)\|_{L^1(\Omega)} \right) + \kappa_p \|\sigma\|_{L^p(\Omega)}^p - c_p, \tag{3.10}$$

for suitable constants $\kappa, \kappa_p > 0$ and $c_p \ge 0$. Analogously, one may also prove a control from below, namely for some C > 0 there holds

$$\mathcal{E}(\varphi, \sigma) \le C \big(\|\varphi\|_V^2 + \|F(\varphi)\|_{L^1(\Omega)} + \|\sigma\|_{L^p(\Omega)}^p + 1 \big).$$
(3.11)

Of course, properties (3.10) and (3.11) are to be intended to hold for (φ, σ) belonging to the natural domain of the energy functional (namely, for the pairs (φ, σ) such that $\varphi \in V$, $F(\varphi) \in L^1(\Omega)$, $\sigma \in L^p(\Omega)$ with $\sigma \geq 0$ a.e. in Ω).

Remark 3.1. In order to deduce (3.10) and (3.11), we have used that

$$\chi \left| \int_{\Omega} \sigma \varphi \right| \le c \|\sigma\|_{L^{p}(\Omega)} \|\varphi\|_{L^{p'}(\Omega)} \le c \|\sigma\|_{L^{p}(\Omega)},$$
(3.12)

where p' is the conjugate exponent to p and the last inequality follows from the uniform boundedness of φ , which is a consequence of the choice of the potential (1.6). It is worth however observing that, in order for the above to work, it would be sufficient to assume that F satisfies the weaker coercivity property

$$F(\varphi) \ge \kappa \|\varphi\|_{L^{q_0}(\Omega)}^{q_0} - c, \tag{3.13}$$

for some $q_0 > p'$, where $c, \kappa > 0$. Actually, in Section 6 below, we will exhibit an approximation of F such that (3.13) holds with $c, \kappa > 0$ independent of the approximation parameter (see (6.14) in the statement of Lemma 6.1 for details, see also [27] for additional considerations).

In order to control the terms on the right-hand side of (3.8), we first observe that, by (2.10) and (2.8),

$$\int_{\Omega} b(\sigma, \varphi)(\gamma(\sigma) - \chi \varphi) = \int_{\{\sigma < 1\}} b(\sigma, \varphi)(\gamma(\sigma) - \chi \varphi) + \int_{\{\sigma \ge 1\}} b(\sigma, \varphi)(\gamma(\sigma) - \chi \varphi) \\
\leq c \int_{\{\sigma < 1\}} (1 + |\sigma \ln \sigma|) + c \int_{\{\sigma \ge 1\}} (1 + \sigma^p) \\
\leq c (1 + ||\sigma||_{L^p(\Omega)}^p).$$
(3.14)

Here, we used once more the uniform boundedness of φ (but there hold the considerations made in Remark 3.1). Next, regarding the term depending on the chemical potential, we replace the expression (1.2) of μ therein, so to obtain

$$\int_{\Omega} \left(h(\sigma,\varphi) - \ell\varphi \right) \mu = \int_{\Omega} \left(h(\sigma,\varphi) - \ell\varphi \right) \left(-\Delta\varphi + f(\varphi) - \chi\sigma \right)$$
$$= -\int_{\Omega} \left(h(\sigma,\varphi) - \ell\varphi \right) \Delta\varphi + \int_{\Omega} \left(h(\sigma,\varphi) - \ell\varphi \right) f(\varphi) - \int_{\Omega} \chi \left(h(\sigma,\varphi) - \ell\varphi \right) \sigma$$
$$=: -I_1 + I_2 - I_3. \tag{3.15}$$

Let us manage the quantities on the right-hand side. First, let us recall that, as noted above, $h(\sigma, \varphi) - \ell \varphi$ has the opposite sign to φ (hence to $f(\varphi)$) when $|\varphi| \sim 1$ and uniformly in σ . This implies that $I_2 \leq c$. Next, by (2.3) with the boundedness of φ (recall, however, Remark 3.1), we deduce

$$|I_3| \le c \|\sigma\|_{L^1(\Omega)} \le c \left(1 + \|\sigma\|_{L^p(\Omega)}^p\right).$$
(3.16)

In order to control I_1 we observe that, recalling assumption (2.5) and applying Young's inequality,

$$-I_{1} = \int_{\Omega} h_{\sigma}(\sigma,\varphi) \nabla \sigma \cdot \nabla \varphi + \int_{\Omega} (h_{\varphi}(\sigma,\varphi) - \ell) |\nabla \varphi|^{2}$$
$$\leq C_{h,2} \int_{\Omega} \frac{1}{1+\sigma} |\nabla \sigma| |\nabla \varphi| + c \|\nabla \varphi\|^{2} \leq \frac{1}{4} \|\nabla \ln \sigma\|^{2} + c \|\nabla \varphi\|^{2}.$$
(3.17)

Then, collecting the above considerations, (3.8) yields

$$\frac{\mathrm{d}}{\mathrm{d}t} \left(\frac{1}{2} \|\nabla\varphi\|^2 + \int_{\Omega} F(\varphi) + \int_{\Omega} \left(\widehat{\gamma}(\sigma) - \chi\sigma\varphi \right) \right) + \|\nabla\mu\|^2 + \int_{\Omega} \alpha(\sigma) |\nabla(\gamma(\sigma) - \chi\varphi)|^2 + \varepsilon \|D\boldsymbol{u}\|^2 + \|\boldsymbol{u}\|^2 \le c \left(1 + \|\sigma\|_{L^p(\Omega)}^p \right) + c \|\nabla\varphi\|^2 + \frac{1}{4} \|\nabla\ln\sigma\|^2.$$
(3.18)

The simplest way to control the last term consists in improving a bit the positivity estimate on σ : to this aim, we test (2.9) by $-1/\sigma$. Also this procedure could be justified by truncation. We then deduce

$$\frac{\mathrm{d}}{\mathrm{d}t} \int_{\Omega} -\ln\sigma + \int_{\Omega} \frac{\gamma'(\sigma)\alpha(\sigma)}{\sigma^2} |\nabla\sigma|^2 = -\int_{\Omega} \frac{b(\sigma,\varphi)}{\sigma} + \chi \int_{\Omega} \frac{\alpha(\sigma)}{\sigma^2} \nabla\varphi \cdot \nabla\sigma.$$
(3.19)

Now, owing to (2.10) and considering that $\sigma \geq 0$ by the Stampacchia truncation argument, we deduce that $\sigma^{-1}b(\sigma, \varphi) \leq b_0$ almost everywhere. Next, by (2.8) one has (as expected)

$$\int_{\Omega} \frac{\gamma'(\sigma)\alpha(\sigma)}{\sigma^2} |\nabla\sigma|^2 = \int_{\Omega} \frac{1+\sigma^{p-1}}{\sigma} \frac{\sigma}{1+\sigma^{p-1}} \frac{1}{\sigma^2} |\nabla\sigma|^2 = \int_{\Omega} \frac{1}{\sigma^2} |\nabla\sigma|^2 = \|\nabla\ln\sigma\|^2.$$
(3.20)

On the other hand,

$$\chi \int_{\Omega} \frac{\alpha(\sigma)}{\sigma^2} \nabla \varphi \cdot \nabla \sigma = \chi \int_{\Omega} \frac{1}{\sigma + \sigma^p} \nabla \varphi \cdot \nabla \sigma.$$
(3.21)

Hence,

$$\chi \int_{\Omega} \frac{\alpha(\sigma)}{\sigma^2} \nabla \varphi \cdot \nabla \sigma \le \frac{1}{2} \int_{\Omega} \frac{1}{\sigma^2} |\nabla \sigma|^2 + c \|\nabla \varphi\|^2 = \frac{1}{2} \|\nabla \ln \sigma\|^2 + c \|\nabla \varphi\|^2.$$
(3.22)

Replacing the above calculations into (3.19), we deduce

$$\frac{\mathrm{d}}{\mathrm{d}t} \int_{\Omega} -\ln\sigma + \frac{1}{2} \|\nabla\ln\sigma\|^2 \le c \left(1 + \|\nabla\varphi\|^2\right). \tag{3.23}$$

Summing (3.23) to (3.18) we then obtain

$$\frac{\mathrm{d}}{\mathrm{d}t} \left(\frac{1}{2} \| \nabla \varphi \|^{2} + \int_{\Omega} F(\varphi) + \int_{\Omega} \left(\widehat{\gamma}(\sigma) - \chi \sigma \varphi - \ln \sigma \right) \right) + \| \nabla \mu \|^{2} + \int_{\Omega} \alpha(\sigma) \left| \nabla (\gamma(\sigma) - \chi \varphi) \right|^{2} \\
+ \varepsilon \| D \boldsymbol{u} \|^{2} + \| \boldsymbol{u} \|^{2} + \frac{1}{4} \| \nabla \ln \sigma \|^{2} \le c \left(1 + \| \sigma \|_{L^{p}(\Omega)}^{p} + \| \nabla \varphi \|^{2} \right).$$
(3.24)

Then, recalling (3.6) and (3.10)-(3.11), an application of Grönwall's lemma permits us to deduce the following properties:

$$\|\varphi\|_{L^{\infty}(0,T;V)} \le c,$$
 (3.25)

$$|F(\varphi)||_{L^{\infty}(0,T;L^{1}(\Omega))} \leq c, \qquad (3.26)$$

$$\nabla \mu \|_{L^2(0,T;H)} \le c,$$
(3.27)

$$\|\boldsymbol{u}\|_{L^{2}(0,T;H)} + \varepsilon^{1/2} \|D\boldsymbol{u}\|_{L^{2}(0,T;H)} \le c, \qquad (3.28)$$

$$\|\sigma\|_{L^{\infty}(0,T;L^{p}(\Omega))} \le c, \tag{3.29}$$

$$\|\alpha + (0)\nabla(\gamma(0) - \chi\varphi)\|_{L^{2}(0,T;H)} \le C,$$
(3.50)

$$\|\ln\sigma\|_{L^{\infty}(0,T;L^{1}(\Omega))} + \|\ln\sigma\|_{L^{2}(0,T;V)} \le c.$$
(3.31)

Note that the above estimates, as well as the ones that follow, are to be intended as a-priori bounds holding uniformly with respect to any hypothetical regularization parameter. Moreover, as is implicitly written in (3.31) we have in particular that $\sigma > 0$ a.e. in Q. Analogously, from (3.26) there formally follows that

$$\|\varphi\|_{L^{\infty}(0,T;L^{\infty}(\Omega))} \le c. \tag{3.32}$$

However, it is worth remarking that (3.32) holds in the limit, and with c = 1, thanks to the specific expression (1.6) of the logarithmic potential F, which constrains φ to take values into [-1, 1]. As F is replaced by a regularizing family F_n , (3.32) in principle may be lost, or replaced by a weaker information (cf. Remark 3.1). Nevertheless, in Section 6, we shall prove that if F_n are suitably designed, then the analogue of (3.32) keeps holding in the approximation, with the constant c (possibly larger than 1), being independent of the approximation parameter n. This fact, which has not been used in deducing the energy estimate (cf. Remark 3.1), will be crucial in the sequel.

That said, applying Korn's inequality, from (3.28) there follows that

$$\varepsilon^{1/2} \| \boldsymbol{u} \|_{L^2(0,T;V)} \le c.$$
 (3.33)

Next, we provide an estimate of the logarithmic term $f(\varphi)$. To this aim, we test (1.2) by $\varphi - \varphi_{\Omega}$ so to obtain

$$\|\nabla\varphi\|^2 + \int_{\Omega} f(\varphi)(\varphi - \varphi_{\Omega}) = \chi \int_{\Omega} \sigma(\varphi - \varphi_{\Omega}) + \int_{\Omega} (\mu - \mu_{\Omega})(\varphi - \varphi_{\Omega}).$$
(3.34)

Then, accounting for (3.2), a well-known argument from [23] (which is robust with respect to regularizations of the potential) permits us to see that

$$\int_{\Omega} f(\varphi)(\varphi - \varphi_{\Omega}) \ge \kappa \| f(\varphi) \|_{L^{1}(\Omega)} - c, \qquad (3.35)$$

where the constants $\kappa > 0$ and $c \ge 0$ depend only on the expression of f and on the parameter δ in (3.2). Then, using (3.35) and (3.32) in (3.34), and applying the Poincaré-Wirtinger inequality, it is not difficult to arrive at

$$\kappa \| f(\varphi) \|_{L^1(\Omega)} \le c \big(1 + \| \sigma \|_{L^1(\Omega)} + \| \nabla \mu \| \big), \tag{3.36}$$

whence, squaring, integrating in time, and recalling (3.27) and (3.29), we deduce

$$\|f(\varphi)\|_{L^2(0,T;L^1(\Omega))} \le c. \tag{3.37}$$

Then, integrating (1.2) over Ω and using the above relation together with (3.29), we obtain

$$\|\mu_{\Omega}\|_{L^2(0,T)} \le c. \tag{3.38}$$

Recalling (3.27), we then have

$$\|\mu\|_{L^2(0,T;V)} \le c. \tag{3.39}$$

We finally observe that, using (3.29) and (3.39), we may reinterpret (1.2) as a time-dependent family of elliptic problems with maximal monotone nonlinearities. This fact allows us the use of a nowadays standard regularity argument, which we just sketch for the reader's convenience. Namely, one can test (1.2) by $|\beta(\varphi)|^{p-1} \operatorname{sign}(\beta(\varphi))$, which is monotone with respect to φ (recall that β denotes the "monotone part" of f). Then, using conditions (3.29) and (3.39) (with Sobolev's embeddings), it is not difficult to arrive at

$$\|\varphi\|_{L^2(0,T;W^{2,p}(\Omega))} + \|f(\varphi)\|_{L^2(0,T;L^p(\Omega))} \le c.$$
(3.40)

More precisely, the second bound is a direct consequence of the estimate, while the first one follows by a further comparison of terms and applying L^p -elliptic regularity results of Agmon-Douglis-Nirenberg type. Note that this procedure will be repeated in the sequence with different (and better) exponents.

Remark 3.2. One may wonder if, under stronger (compared to (2.13)) positivity conditions on the initial datum, relation (3.31) could be improved, getting for instance a strictly positive uniform lower bound for σ . Actually, a natural way to prove such a property would be to test (1.3) by (a truncation of) $-\sigma^{-\eta}$ for larger values of the exponent η , and then possibly applying some iteration procedure. However, even though condition (2.10) seems to allow such a procedure, the regularity (3.25) available at this level for $\nabla \varphi$ seems not sufficient in order to close the estimate.

4 Entropy-type estimates

The energy estimate derived in the previous section holds without any restriction on the exponent p and on the the dimension. However, if p is very close to 1 (meaning that the cross diffusion term behaves at infinity similarly to the genuine Keller-Segel case), then, at least for d = 3, the outcome, in terms of a-priori regularity, seems not sufficient in order to pass to the limit in the approximation. The main issue arises when dealing with the product terms (more precisely, the transport term in (1.3) and the Korteweg term depending on σ in (1.4)), which may not be identified under the sole regularity provided by the energy bound.

The following additional argument, based on the so-called "entropy" method for fourth order PDE's, is crucial in order to improve the available a-priori regularity and, in particular, to prove separate bounds for the summands in H (cf. (2.24)). We shall present two versions of the entropy bound, corresponding to the regularity settings of Theorem 2.1 and of Theorem 2.3. The first version, referring to the case d = 2, is simpler (but likely not optimal): indeed, in this situation for every p > 1 we can "directly" obtain an estimate corresponding to the "natural" parabolic regularity that would hold also in the "linear" case. We decided to start detailing this argument because it gives a first idea of the method, with reduced technicalities compared to the subsequent 3D proof.

4.1 Entropy estimate: simple version for d = 2

The basic strategy to obtain the entropy bound simply consists in testing (1.2) by $-\Delta\varphi$. As done for the previous estimates, we shall directly work on the original system (1.1)-(1.5); in Section 6, we will discuss how also this procedure might be adapted so to fit the framework of a regularization scheme.

That said, standard manipulations give

$$\|\Delta\varphi\|^2 = -\int_{\Omega} f'(\varphi) |\nabla\varphi|^2 + \int_{\Omega} \nabla\varphi \cdot \nabla\mu + \chi \int_{\Omega} \nabla\varphi \cdot \nabla\sigma.$$
(4.1)

Now, using (3.25) and recalling that $f'(\cdot) \geq -\lambda$, we easily obtain

$$-\int_{\Omega} f'(\varphi) |\nabla \varphi|^2 + \int_{\Omega} \nabla \varphi \cdot \nabla \mu + \chi \int_{\Omega} \nabla \varphi \cdot \nabla \sigma \le c \left(1 + \|\nabla \mu\| + \|\nabla \sigma\|\right).$$
(4.2)

Hence, replacing the above into (4.1) and squaring, we deduce

$$\|\Delta\varphi\|^{4} \le c_{1} \left(1 + \|\nabla\mu\|^{2} + \|\nabla\sigma\|^{2}\right), \tag{4.3}$$

where the constant c_1 depends only on the outcome of the previous a-priori estimates (hence it is a computable quantity independent of the approximation parameters).

In order to control the last term on the right-hand side, we test (1.3) by σ . Then, assumption (2.10) and simple integrations by parts give

$$\frac{1}{2}\frac{\mathrm{d}}{\mathrm{d}t}\|\sigma\|^2 + \|\nabla\sigma\|^2 \le c\left(1 + \|\sigma\|^2\right) + \chi \int_{\Omega} \sigma^{2-p} |\nabla\sigma| |\nabla\varphi|.$$

$$(4.4)$$

Now, the last term on the right-hand side can be estimated as follows:

$$\begin{split} \chi \int_{\Omega} \sigma^{2-p} |\nabla \sigma| |\nabla \varphi| &\leq c \|\sigma^{1/2}\|_{L^{\frac{4(1+\epsilon)}{\epsilon}}(\Omega)} \|\sigma^{\frac{3-2p}{2}}\|_{L^{2(1+\epsilon)}(\Omega)} \|\nabla \sigma\| \|\nabla \varphi\|_{L^{\frac{4(1+\epsilon)}{\epsilon}}(\Omega)} \\ &\leq c \|\sigma\|_{L^{\frac{2(1+\epsilon)}{\epsilon}}(\Omega)}^{1/2} \|\sigma^{\frac{3-2p}{2}}\|_{L^{2(1+\epsilon)}(\Omega)} \|\nabla \sigma\| \|\nabla \varphi\|_{L^{\frac{2}{2(1+\epsilon)}}}^{\frac{\epsilon}{2(1+\epsilon)}} \|\varphi\|_{H^{2}(\Omega)}^{\frac{2+\epsilon}{2(1+\epsilon)}} \\ &\leq c_{\epsilon} \|\sigma\|_{V}^{3/2} \|\sigma^{\frac{3-2p}{2}}\|_{L^{2(1+\epsilon)}(\Omega)} \left(1 + \|\Delta \varphi\|_{L^{\frac{2+\epsilon}{2(1+\epsilon)}}}\right) \\ &\leq \frac{1}{2} \|\sigma\|_{V}^{2} + c_{\epsilon} \|\sigma^{\frac{3-2p}{2}}\|_{L^{2(1+\epsilon)}(\Omega)}^{4} \left(1 + \|\Delta \varphi\|_{L^{\frac{2}{\epsilon}}}^{\frac{2(2+\epsilon)}{1+\epsilon}}\right) \\ &\leq \frac{1}{2} \|\sigma\|_{V}^{2} + \delta \|\Delta \varphi\|^{4} + c + c_{\epsilon,\delta} \|\sigma^{\frac{3-2p}{2}}\|_{L^{2(1+\epsilon)}(\Omega)}^{\frac{8(1+\epsilon)}{\epsilon}}. \end{split}$$
(4.5)

Here, we have used two dimensional embeddings with the Gagliardo-Nirenberg and Young inequalities. Moreover, $\epsilon > 0$ and $\delta > 0$ are "small" parameters to be specified below, whereas, for instance, $c_{\epsilon} > 0$ denotes a "large" constant depending on the choice of ϵ . To specify these quantities, we first observe that, as p > 1, we may use (3.29) to control the last term on the right-hand side, provided that $\epsilon > 0$ is taken so small that

$$p \ge \frac{3(1+\epsilon)}{3+2\epsilon}, \quad \text{i.e.} \ (3-2p)(1+\epsilon) \le p.$$

$$(4.6)$$

Next, let us observe that, by elliptic regularity and (3.25),

$$\|\varphi\|_{H^2(\Omega)}^4 \le c_\Omega \left(\|\varphi\|^4 + \|\Delta\varphi\|^4\right) \le c_\Omega \left(1 + \|\Delta\varphi\|^4\right).$$

$$(4.7)$$

Then, replacing (4.5) into (4.4) and using (4.6), we obtain

$$\frac{1}{2}\frac{\mathrm{d}}{\mathrm{d}t}\|\sigma\|^2 + \frac{1}{2}\|\nabla\sigma\|^2 \le c(1+\|\sigma\|^2) + c_\Omega\delta\|\Delta\varphi\|^4 + c_\delta\|\sigma\|_{L^p(\Omega)}^P,\tag{4.8}$$

where $P = 4(1 + \epsilon)(3 - 2p)/\epsilon$ is a "large", but otherwise fixed, exponent depending on ϵ . Next, recalling (4.3), we take $\delta > 0$ so small that $2c_1c_\Omega\delta \leq 1/4$. Then, we multiply (4.3) by $2c_\Omega\delta$ and subsequently sum the resulting relation to (4.8) so to obtain

$$\frac{1}{2}\frac{\mathrm{d}}{\mathrm{d}t}\|\sigma\|^{2} + c_{\Omega}\delta\|\Delta\varphi\|^{4} + \frac{1}{4}\|\nabla\sigma\|^{2} \le c(1+\|\sigma\|^{2}+\|\nabla\mu\|^{2}) + c_{\delta}\|\sigma\|_{L^{p}(\Omega)}^{P}.$$
(4.9)

Then, integrating in time and using Grönwall's lemma with (3.39) and (3.29), we readily deduce the additional estimates

$$\|\varphi\|_{L^4(0,T;H^2(\Omega))} \le c,\tag{4.10}$$

$$\|\sigma\|_{L^{\infty}(0,T;H)} + \|\sigma\|_{L^{2}(0,T;V)} \le c, \tag{4.11}$$

where we also used assumption (2.17) on the initial datum. With the above regularity at hand, we can also improve the control of the logarithmic term in (1.2). In particular, using that, for d = 2, $V \subset L^P(\Omega)$ for every $P \in [1, \infty)$, estimates (3.39) and (4.11) and the same elliptic regularity argument used at the end of Section 3 permit us to deduce

$$\|f(\varphi)\|_{L^2(0,T;L^P(\Omega))} + \|\varphi\|_{L^2(0,T;W^{2,P}(\Omega))} \le c_P \text{ for all } P \in [1,\infty).$$

$$(4.12)$$

4.2 Entropy estimate: improved version for d = 3

In this part we adapt the above argument so to deal with the assumptions of Theorem 2.3 regarding the three-dimensional case. In this setting, we will obtain a more accurate estimation of the cross-diffusion terms by valorizing the following chain of inequalities:

$$\|\nabla\varphi\|_{L^{4}(\Omega)} \le c \|\varphi\|_{L^{\infty}(\Omega)}^{1/2} \|\varphi\|_{H^{2}(\Omega)}^{1/2} \le c (1 + \|\Delta\varphi\|^{1/2}),$$
(4.13)

which holds in any space dimension and follows from the Gagliardo-Nirenberg interpolation inequality [24] and standard elliptic regularity results (recall that φ satisfies a no-flux boundary condition). Relation (4.13) uses in an essential way the uniform boundedness of φ . This surely holds "in the limit" thanks to the structure of the potential (1.6), but it can be assumed to be satisfied (uniformly) in the approximation as shown in Lemma 6.1 below.

That said, we start, as before, by testing (1.2) by $-\Delta\varphi$ so to obtain the analogue of (4.1), namely

$$\|\Delta\varphi\|^2 \le c \big(1 + \|\nabla\mu\|\big) + \chi \int_{\Omega} \nabla\varphi \cdot \nabla\sigma.$$
(4.14)

Moreover, we test (2.9) by σ^{p-1} , with the purpose of obtaining a decoupled estimate for the crossdiffusion term. We deduce

$$\frac{1}{p}\frac{\mathrm{d}}{\mathrm{d}t}\int_{\Omega}\sigma^{p} + \frac{4(p-1)}{p^{2}}\int_{\Omega}|\nabla\sigma^{\frac{p}{2}}|^{2} \leq c + c\int_{\Omega}\sigma^{p} + (p-1)\chi\int_{\Omega}\frac{\sigma^{p-1}}{1+\sigma^{p-1}}\nabla\varphi\cdot\nabla\sigma.$$
(4.15)

In order to control the terms on the right-hand side of (4.14) and (4.15), which are of the same type, we may proceed as follows:

$$\int_{\Omega} \nabla \varphi \cdot \nabla \sigma \leq c \int_{\Omega} |\nabla \varphi| |\nabla \sigma| \leq c \int_{\Omega} \sigma^{\frac{2-p}{2}} |\nabla \sigma^{p/2}| |\nabla \varphi|
\leq c \|\sigma^{\frac{2-p}{2}}\|_{L^{4}(\Omega)} \|\nabla \sigma^{p/2}\| \|\nabla \varphi\|_{L^{4}(\Omega)}
\leq c \|\sigma^{\frac{2-p}{2}}\|_{L^{\frac{2p}{2-p}}(\Omega)}^{\eta} \|\sigma^{\frac{2-p}{2}}\|_{L^{\frac{6p}{2-p}}(\Omega)}^{1-\eta} \|\nabla \sigma^{p/2}\| (1+\|\Delta \varphi\|^{1/2})
\leq c \|\sigma\|_{L^{p}(\Omega)}^{\frac{\eta(2-p)}{2}} \|\sigma^{\frac{p}{2}\frac{2-p}{p}}\|_{L^{\frac{6p}{2-p}}(\Omega)}^{1-\eta} \|\nabla \sigma^{p/2}\| (1+\|\Delta \varphi\|^{1/2})
\leq c \|\sigma\|_{L^{p}(\Omega)}^{\frac{\eta(2-p)}{2}} \|\sigma^{p}\|_{L^{6}(\Omega)}^{(2-p)(1-\eta)} \|\nabla \sigma^{p/2}\| (1+\|\Delta \varphi\|^{1/2}),$$
(4.16)

where the interpolation exponent η is given by the relation

$$\eta \frac{2-p}{2p} + (1-\eta)\frac{2-p}{6p} = \frac{1}{4}, \quad \text{i.e., } \eta = \frac{5p-4}{8-4p}, \quad 1-\eta = \frac{12-9p}{8-4p}.$$
(4.17)

Note that here we have implicitly assumed that $p \le 4/3$, whence $2p/(2-p) \le 4$, which justifies the use of interpolation. However, it is clear that, for p > 4/3, the procedure can be modified as follows:

$$\int_{\Omega} \nabla \varphi \cdot \nabla \sigma \leq c \| \sigma^{\frac{2-p}{2}} \|_{L^{4}(\Omega)} \| \nabla \sigma^{p/2} \| \| \nabla \varphi \|_{L^{4}(\Omega)}
\leq c (1 + \| \sigma^{1/3} \|_{L^{4}(\Omega)}) \| \nabla \sigma^{p/2} \| \| \nabla \varphi \|_{L^{4}(\Omega)}
\leq c (1 + \| \sigma \|_{L^{p}(\Omega)}^{1/3}) \| \nabla \sigma^{p/2} \| \| \nabla \varphi \|_{L^{4}(\Omega)} \leq c \| \nabla \sigma^{p/2} \| (1 + \| \Delta \varphi \|^{1/2}),$$
(4.18)

the last inequality following from (3.29). It is then not difficult to adapt the remainder of the procedure so to cover also this (simpler) situation.

Thus, going back to the case $p \leq 4/3$, (4.16) can be continuated as follows:

$$\int_{\Omega} \nabla \varphi \cdot \nabla \sigma \leq c \|\sigma\|_{L^{p}(\Omega)}^{\frac{\eta(2-p)}{2}} \|\sigma^{\frac{p}{2}}\|_{L^{6}(\Omega)}^{\frac{(2-p)(1-\eta)}{p}} \|\nabla \sigma^{p/2}\| \left(1 + \|\Delta \varphi\|^{1/2}\right) \\
\leq c \|\sigma\|_{L^{p}(\Omega)}^{\frac{5p-4}{8}} \left(1 + \|\nabla \sigma^{p/2}\|^{\frac{12-5p}{4p}}\right) \left(1 + \|\Delta \varphi\|^{1/2}\right).$$
(4.19)

Replacing the above into (4.14), we deduce

$$\|\Delta\varphi\|^{3/2} \le c + c \|\nabla\mu\|^{3/4} + c \|\sigma\|_{L^p(\Omega)}^{\frac{5p-4}{8}} \left(1 + \|\nabla\sigma^{p/2}\|^{\frac{12-5p}{4p}}\right).$$
(4.20)

Now, let us go back to (4.15). Estimating the last term on the right-hand side as done in (4.19), we have

$$\frac{1}{p}\frac{\mathrm{d}}{\mathrm{d}t}\|\sigma\|_{L^{p}(\Omega)}^{p} + \frac{4(p-1)}{p^{2}}\|\nabla\sigma^{\frac{p}{2}}\|^{2} \le c + c\|\sigma\|_{L^{p}(\Omega)}^{p} + c\|\sigma\|_{L^{p}(\Omega)}^{\frac{5p-4}{8}} (1 + \|\nabla\sigma^{p/2}\|^{\frac{12-5p}{4p}})(1 + \|\Delta\varphi\|^{1/2}).$$
(4.21)

Using (4.20) to control the term with the Laplacian, we then infer

$$\frac{1}{p} \frac{\mathrm{d}}{\mathrm{d}t} \|\sigma\|_{L^{p}(\Omega)}^{p} + \frac{4(p-1)}{p^{2}} \|\nabla\sigma^{\frac{p}{2}}\|^{2} \leq c + C \\
+ C \left(1 + \|\nabla\sigma^{p/2}\|^{\frac{12-5p}{4p}}\right) \left(1 + \|\nabla\mu\|^{1/4} + C \left(1 + \|\nabla\sigma^{p/2}\|^{\frac{12-5p}{12p}}\right)\right).$$
(4.22)

Here and below, C represents a "large", but otherwise computable, constant depending on (some power of) the $L^p(\Omega)$ -norm of σ , which has already been estimated uniformly in time thanks to (3.29).

Now, we first observe that, as p > 12/11, it follows

$$\frac{12-5p}{4p} < \frac{7}{4}, \quad \text{whence} \quad \boldsymbol{C} \|\nabla \sigma^{p/2}\|^{\frac{12-5p}{4p}} \|\nabla \mu\|^{1/4} \le \frac{p-1}{p^2} \|\nabla \sigma^{p/2}\|^2 + \boldsymbol{C} \|\nabla \mu\|^2 + c.$$
(4.23)

Moreover, the condition p > 12/11 also guarantees

$$\frac{12-5p}{4p} + \frac{12-5p}{12p} = \frac{12-5p}{3p} < 2, \quad \text{whence} \quad C \|\nabla \sigma^{p/2}\|^{\frac{12-5p}{3p}} \le \frac{p-1}{p^2} \|\nabla \sigma^{p/2}\|^2 + C.$$
(4.24)

On account of these considerations, (4.22) gives

$$\frac{1}{p}\frac{\mathrm{d}}{\mathrm{d}t}\|\sigma\|_{L^{p}(\Omega)}^{p} + \frac{2(p-1)}{p^{2}}\|\nabla\sigma^{\frac{p}{2}}\|^{2} \leq C(1+\|\nabla\mu\|^{2}).$$
(4.25)

Hence, applying Grönwall's lemma, we obtain

$$\|\sigma^{p/2}\|_{L^2(0,T;V)} \le c,\tag{4.26}$$

which implies the desired decoupled regularity bound for the cross-diffusion term in (1.3). Finally, going back to (4.20), we prove an estimate of $\Delta \varphi$. To this aim, we need to distinguish between two case: if $p \leq 3/2$, then taking a suitable power of that relation we deduce

$$\|\Delta\varphi\|^{\frac{12p}{12-5p}} \le c + c \|\nabla\mu\|^{\frac{6p}{12-5p}} + c \|\sigma\|^{\frac{p(5p-4)}{12-5p}}_{L^p(\Omega)} (1 + \|\nabla\sigma^{p/2}\|^2), \tag{4.27}$$

and it turns out that $\frac{6p}{12-5p} \leq 2$. Hence, integrating in time and using (4.26) we infer

$$\|\varphi\|_{L^{\frac{12p}{12-5p}}(0,T;H^2(\Omega))} \le c.$$
(4.28)

On the other hand, if p > 3/2, we take the (8/3)-power of (4.20), which gives

$$\|\Delta\varphi\|^{4} \le c + c\|\nabla\mu\|^{2} + c\|\sigma\|_{L^{p}(\Omega)}^{\frac{5p-4}{3}} \left(1 + \|\nabla\sigma^{p/2}\|^{\frac{24-10p}{3p}}\right)$$
(4.29)

Then, p > 3/2 implies $\frac{24-10p}{3p} < 2$. Hence, integrating once more in time and using (4.26), in place of (4.28) we get

$$\|\varphi\|_{L^4(0,T;H^2(\Omega))} \le c. \tag{4.30}$$

4.3 Additional regularity for d = 3 by iteration of the entropy estimate

We now see how the previous procedure can be adapted in order to achieve additional summability of σ in the case when $\sigma_0 \in L^q(\Omega)$. For technical reasons, we will start assuming $p < q \leq 20/11$. For larger q the procedure will be carried out by means of a two-step iteration. That said, we start with testing (2.9) by σ^{q-1} . Then, using once more (2.10), we deduce

$$\frac{1}{q}\frac{\mathrm{d}}{\mathrm{d}t}\int_{\Omega}\sigma^{q} + \frac{4(q-1)}{q^{2}}\int_{\Omega}|\nabla\sigma^{\frac{q}{2}}|^{2} \leq c + c\int_{\Omega}\sigma^{q} + (q-1)\chi\int_{\Omega}\frac{\sigma^{q-1}}{1+\sigma^{p-1}}\nabla\varphi\cdot\nabla\sigma \leq c + c\int_{\Omega}\sigma^{q} + c\int_{\Omega}\sigma^{q-p}|\nabla\varphi||\nabla\sigma|.$$
(4.31)

In order to estimate the last term on the right-hand side in an optimal way, we go back to relation (4.14) and modify (4.16) as follows:

$$\begin{split} \chi \int_{\Omega} \nabla \varphi \cdot \nabla \sigma &\leq c \int_{\Omega} \sigma^{\frac{2-q}{2}} |\nabla \sigma^{q/2}| ||\nabla \varphi| \\ &\leq c ||\sigma^{\frac{2-q}{2}}||_{L^{4}(\Omega)} ||\nabla \sigma^{q/2}|| ||\nabla \varphi||_{L^{4}(\Omega)} \\ &\leq c ||\sigma^{\frac{2-q}{2}}||_{L^{\frac{2p}{2-q}}(\Omega)}^{\eta} ||\sigma^{\frac{2-q}{2}}||_{L^{\frac{6q}{2-q}}(\Omega)}^{1-\eta} ||\nabla \sigma^{q/2}|| (1 + ||\Delta \varphi||^{1/2}) \\ &\leq c ||\sigma||_{L^{p}(\Omega)}^{\frac{\eta(2-q)}{2}} ||\sigma^{\frac{q}{2}\frac{2-q}{q}}||_{L^{\frac{6q}{2-q}}(\Omega)}^{1-\eta} ||\nabla \sigma^{q/2}|| (1 + ||\Delta \varphi||^{1/2}) \\ &\leq c ||\sigma||_{L^{p}(\Omega)}^{\frac{\eta(2-q)}{2}} ||\sigma^{\frac{q}{2}}||_{L^{6}(\Omega)}^{\frac{(2-q)(1-\eta)}{q}} ||\nabla \sigma^{q/2}|| (1 + ||\Delta \varphi||^{1/2}). \end{split}$$
(4.32)

As before, the previous computation is rigorous provided that

$$\frac{2p}{2-q} < 4 < \frac{6q}{2-q}.$$
(4.33)

Clearly the right inequality is always satisfied as far as 1 < q < 2. Concerning the left one, it corresponds to asking 2 - q > p/2. On the other hand, if $2 - q \le p/2$, as before we may notice that

$$\begin{split} \chi \int_{\Omega} \nabla \varphi \cdot \nabla \sigma &\leq c \int_{\Omega} \sigma^{\frac{2-q}{2}} |\nabla \sigma^{q/2}| |\nabla \varphi| \\ &\leq c \| \sigma^{\frac{2-q}{2}} \|_{L^4(\Omega)} \| \nabla \sigma^{q/2} \| \| \nabla \varphi \|_{L^4(\Omega)} \leq c \big(1 + \| \sigma^{\frac{p}{4}} \|_{L^4(\Omega)} \big) \| \nabla \sigma^{q/2} \| \| \nabla \varphi \|_{L^4(\Omega)} \\ &\leq C \| \nabla \sigma^{q/2} \| \big(1 + \| \Delta \varphi \|^{1/2} \big), \end{split}$$

$$(4.34)$$

and the argument proceeds with straighforward adaptations. That said, going back to the case $\frac{2p}{2-q} < 4$, the interpolation exponent η is then provided by the relation

$$\eta \frac{2-q}{2p} + (1-\eta)\frac{2-q}{6q} = \frac{1}{4}, \quad \text{i.e., } \eta = \frac{(5q-4)p}{2(2-q)(3q-p)}, \quad 1-\eta = \frac{q(12-6q-3p)}{2(2-q)(3q-p)}.$$
(4.35)

Hence, (4.32) can be continuated as follows:

$$\begin{split} \chi &\int_{\Omega} \nabla \varphi \cdot \nabla \sigma \leq c \|\sigma\|_{L^{p}(\Omega)}^{\frac{\eta(2-q)}{2}} \|\sigma^{\frac{q}{2}}\|_{L^{6}(\Omega)}^{\frac{(2-q)(1-\eta)}{q}} \|\nabla \sigma^{q/2}\| \left(1 + \|\Delta \varphi\|^{1/2}\right) \\ &\leq c \|\sigma\|_{L^{p}(\Omega)}^{\frac{(5q-4)p}{4(3q-p)}} \left(1 + \|\nabla \sigma^{q/2}\|^{\frac{12-6q-3p}{6q-2p}}\right) \|\nabla \sigma^{q/2}\| \left(1 + \|\Delta \varphi\|^{1/2}\right) \\ &\leq C \left(1 + \|\nabla \sigma^{q/2}\|^{\frac{12-5p}{6q-2p}}\right) \left(1 + \|\Delta \varphi\|^{1/2}\right), \end{split}$$
(4.36)

where, as before, C > 0 is a (possibly large) constant depending on the outcome of the previous estimates (specifically, on (3.29)).

Hence, in view of the above procedure, (4.20) is replaced by

$$\|\Delta\varphi\|^{3/2} \le c + c \|\nabla\mu\|^{3/4} + C\left(1 + \|\nabla\sigma^{q/2}\|^{\frac{12-5p}{6q-2p}}\right),\tag{4.37}$$

or, in other words,

$$\|\Delta\varphi\|^{\frac{18q-6p}{12-5p}} \le c + c\|\nabla\mu\|^{\frac{9q-3p}{12-5p}} + C(1+\|\nabla\sigma^{q/2}\|^2).$$
(4.38)

Now, let us go back to (4.31). We propose the following estimation strategy:

$$c \int_{\Omega} \sigma^{q-p} |\nabla \sigma| |\nabla \varphi| \leq c \int_{\Omega} \sigma^{\frac{q-2p+2}{2}} |\nabla \sigma^{q/2}| |\nabla \varphi|$$

= $c \int_{\Omega} \sigma^{\frac{q}{2} \frac{q-2p+2}{q}} |\nabla \sigma^{q/2}| |\nabla \varphi| \leq c \|\sigma^{\frac{q}{2} \frac{q-2p+2}{q}}\|_{L^{4}(\Omega)} \|\nabla \sigma^{q/2}\| \|\nabla \varphi\|_{L^{4}(\Omega)}$
 $\leq c \|\sigma^{q/2}\|_{L^{\frac{4(q-2p+2)}{q}}(\Omega)}^{\frac{q-2p+2}{q}} \|\nabla \sigma^{q/2}\| (1 + \|\Delta \varphi\|^{1/2}).$ (4.39)

The first factor on the right-hand side is controlled by interpolation with the following choice:

$$\frac{q}{4(q-2p+2)} = \eta \frac{q}{2p} + (1-\eta)\frac{1}{6},$$
(4.40)

whence

$$\eta = \frac{pq + 4p^2 - 4p}{(q - 2p + 2)(6q - 2p)} = \frac{pq + 4p^2 - 4p}{6q^2 + 4p^2 - 14pq - 4p + 12q},$$
(4.41)

$$1 - \eta = \frac{6q^2 - 15pq + 12q}{(q - 2p + 2)(6q - 2p)} = \frac{6q^2 - 15pq + 12q}{6q^2 + 4p^2 - 14pq - 4p + 12q}.$$
(4.42)

To be precise, in deducing the last inequality in (4.39) we have implicitly used that $\frac{4(q-2p+2)}{q} \ge 1$ (which is always true), while in (4.40) we have used the more restrictive condition

$$\frac{2p}{q} \le \frac{4(q-2p+2)}{q} \le 6. \tag{4.43}$$

Actually, it is easy to check that the right inequality in (4.43) is always satisfied in our exponents range. This is not always true for the left inequality, instead; nevertheless, if the opposite holds, then it follows 2(q - 2p + 2) < p. Hence, going back to (4.39), one has

$$c \int_{\Omega} \sigma^{q-p} |\nabla \sigma| |\nabla \varphi| \le c \| \sigma^{\frac{q-2p+2}{2}} \|_{L^{4}(\Omega)} \|\nabla \sigma^{q/2}\| \|\nabla \varphi\|_{L^{4}(\Omega)}$$

$$= c \left(\int_{\Omega} \sigma^{2(q-2p+2)} \right)^{1/4} \|\nabla \sigma^{q/2}\| \|\nabla \varphi\|_{L^{4}(\Omega)} \le c \left(1 + \|\sigma\|_{L^{p}(\Omega)}^{p/4}\right) \|\nabla \sigma^{q/2}\| \|\nabla \varphi\|_{L^{4}(\Omega)}$$

$$\le C \|\nabla \sigma^{q/2}\| \left(1 + \|\Delta \varphi\|^{1/2}\right), \tag{4.44}$$

and the remainder of the procedure can be adapted up to small variations.

Hence, going back to the (more delicate) case when (4.43) holds, (4.39) can be continuated as follows:

$$\begin{split} c \int_{\Omega} \sigma^{q-p} |\nabla \sigma| |\nabla \varphi| \\ &\leq c \|\sigma^{q/2}\| \frac{\frac{n(q-2p+2)}{q}}{L^{\frac{2p}{q}}(\Omega)} \|\sigma^{q/2}\|_{L^{6}(\Omega)}^{(1-\eta)\frac{q-2p+2}{q}} \|\nabla \sigma^{q/2}\| \left(1 + \|\Delta \varphi\|^{1/2}\right) \\ &\leq c \|\sigma^{q/2}\| \frac{\frac{pq+4p^{2}-4p}{q(6q-2p)}}{L^{\frac{2p}{q}}(\Omega)} \|\sigma^{q/2}\| \frac{\frac{6q-15p+12}{6q-2p}}{V} \|\nabla \sigma^{q/2}\| \left(1 + \|\Delta \varphi\|^{1/2}\right) \\ &\leq c \|\sigma\| \frac{\frac{pq+4p^{2}-4p}{2(6q-2p)}}{L^{p}(\Omega)} \left(1 + \|\nabla \sigma^{q/2}\| \frac{\frac{12q-17p+12}{6q-2p}}{6q-2p}\right) \left(1 + \|\Delta \varphi\|^{1/2}\right). \end{split}$$
(4.45)

At this point, we replace (4.45) into (4.31) and subsequently estimate the term with the Laplacian by means of (4.37). We obtain

$$\frac{1}{q} \frac{\mathrm{d}}{\mathrm{d}t} \|\sigma\|_{L^{q}(\Omega)}^{q} + \frac{4(q-1)}{q^{2}} \|\nabla\sigma^{\frac{q}{2}}\|^{2} \leq c + c \|\sigma\|_{L^{q}(\Omega)}^{q} + C\left(1 + \|\nabla\sigma^{q/2}\|^{\frac{12q-17p+12}{6q-2p}}\right)\left(1 + \|\Delta\varphi\|^{1/2}\right) \\
\leq c + c \|\sigma\|_{L^{q}(\Omega)}^{q} + C\left(1 + \|\nabla\sigma^{q/2}\|^{\frac{12q-17p+12}{6q-2p}}\right)\left(1 + \|\nabla\mu\|^{1/4} + \|\nabla\sigma^{q/2}\|^{\frac{12-5p}{18q-6p}}\right).$$
(4.46)

Now, in order to control the right-hand side via Grönwall's lemma we need two conditions. The first one, in analogy with (4.23), corresponds to

$$\frac{12q - 17p + 12}{6q - 2p} \le \frac{7}{4}, \quad \text{or, equivalently, } q \le 9p - 8.$$
(4.47)

It is then easy to check that the above inequality holds under the technical condition $p < q \le 20/11$ (actually, if $p \ge 10/9$ then it holds for every $q \in (p, 2]$). This implies

$$C \|\nabla \sigma^{q/2}\|^{\frac{12q-17p+12}{6q-2p}} \|\nabla \mu\|^{1/4} \le \frac{q-1}{q^2} \|\nabla \sigma^{q/2}\|^2 + C \|\nabla \mu\|^2 + c.$$
(4.48)

Moreover, let us compute

$$\frac{12q - 17p + 12}{6q - 2p} + \frac{12 - 5p}{18q - 6p} = \frac{18q - 28p + 24}{9q - 3p}.$$
(4.49)

Then, we notice that

$$\frac{18q - 28p + 24}{9q - 3p} < 2 \quad \text{corresponds exactly to } 22p > 24, \text{ i.e. } p > 12/11.$$
(4.50)

Hence, in our admissible range of exponents, independently of the value assumed by q, we have

$$C\left(1 + \|\nabla\sigma^{q/2}\|^{\frac{12q-17p+12}{6q-2p}}\right)\left(1 + \|\nabla\sigma^{q/2}\|^{\frac{12-5p}{18q-6p}}\right) \le C + \frac{q-1}{q^2}\|\nabla\sigma^{q/2}\|^2.$$
(4.51)

Replacing (4.48) and (4.51) into (4.46), we then arrive at

$$\frac{1}{q}\frac{\mathrm{d}}{\mathrm{d}t}\|\sigma\|_{L^{q}(\Omega)}^{q} + \frac{2(q-1)}{q^{2}}\|\nabla\sigma^{\frac{q}{2}}\|^{2} \le c + c\|\sigma\|_{L^{q}(\Omega)}^{q} + C(1 + \|\nabla\mu\|^{2}).$$
(4.52)

Applying Grönwall's lemma and recalling (2.36), we obtain

$$\|\sigma^{q/2}\|_{L^{\infty}(0,T;H)} + \|\sigma^{q/2}\|_{L^{2}(0,T;V)} \le c.$$
(4.53)

To conclude, we deduce some additional regularity for $\Delta \varphi$. To this aim, we go back to (4.38) and notice that, if $\frac{18g-6p}{12-5p} < 4$ then the exponent of the norm of $\nabla \mu$ on the right-hand side is less than 2. Integrating in time, we then deduce

$$\|\varphi\|_{L^{\frac{18q-6p}{12-5p}}(0,T;H^2(\Omega))} \le c.$$
(4.54)

On the other hand, if $\frac{18q-6p}{12-5p} \ge 4$, as before we need to rescale the inequality before integrating in time. In that case, we obtain

$$\|\varphi\|_{L^4(0,T;H^2(\Omega))} \le c. \tag{4.55}$$

Notice that the above conditions correspond exactly to the choice of the exponent P_0 in the statement of Theorem 2.3 (cf. in particular the first (2.45)).

Let us now remove the technical condition $p < q \leq 20/11$ and, for simplicity, let us just consider the worst case scenario, i.e., when $p \sim 12/11$. In that case, let us assume, as stated in (2.36), that $\sigma_0 \in L^q(\Omega)$ for q > 20/11. In this case, we proceed by bootstrap: namely, we first perform the above argument for $q_0 = 20/11$ (we note this "temporary" exponent as q_0 so to distinguish it from q which now corresponds to the summability of the initial datum). Then, we achieve the regularity properties (4.53) (where $q_0 = 20/11$) and (4.55) (note that, as $q_0 = 20/11$, then $\frac{18q_0-6p}{12-5p} \geq 4$ for any $p \in (12/11, q_0)$). Then, to bootstrap regularity, assuming for simplicity q = 2 (the "intermediate" case when σ_0 lies in $L^q(\Omega)$ for $q \in (20/11, 2)$ can be treated with small modifications) one may achieve the "linear" parabolic regularity simply by testing (1.3) by σ . This gives

$$\frac{1}{2}\frac{\mathrm{d}}{\mathrm{d}t}\|\sigma\|^2 + \|\nabla\sigma\|^2 \le c\left(1 + \|\sigma\|^2\right) + c\int_{\Omega}\sigma^{2-p}\nabla\sigma\cdot\nabla\varphi.$$
(4.56)

The last term, being p > 12/11, can be controlled this way:

$$\int_{\Omega} \sigma^{2-p} \nabla \sigma \cdot \nabla \varphi \leq c \left(1 + \| \sigma^{10/11} \|_{L^{4}(\Omega)} \right) \| \nabla \sigma \| \| \nabla \varphi \|_{L^{4}(\Omega)}
\leq c \left(1 + \| \sigma^{10/11} \|_{L^{6}(\Omega)}^{3/4} \| \sigma^{10/11} \|^{1/4} \right) \| \nabla \sigma \| \left(1 + \| \Delta \varphi \|^{1/2} \right)
\leq c \left(1 + \| \sigma \|_{L^{6}(\Omega)}^{3/4} \right) \| \nabla \sigma \| \left(1 + \| \Delta \varphi \|^{1/2} \right) \leq c + c \| \Delta \varphi \|^{4} + \frac{1}{2} \| \nabla \sigma \|^{2}.$$
(4.57)

Hence, integrating in time and using (2.36), we obtain once more (4.53) (for q = 2). With simple adaptations, one can get the analogue also in the case $q \in (20/11, 2)$. By interpolation, (4.53) implies

$$\|\sigma\|_{L^2(0,T;L^{\frac{3q}{3-q}}(\Omega))} \le c.$$
(4.58)

Hence, interpreting once more (1.2) as a time-dependent family of elliptic problems and using (4.58) together with (3.39), the usual regularity argument gives

$$\|\varphi\|_{L^{2}(0,T;W^{2,\frac{3q}{3-q}}(\Omega))} + \|f(\varphi)\|_{L^{2}(0,T;L^{\frac{3q}{3-q}}(\Omega))} \le c.$$
(4.59)

5 Passage to the limit

In this section we complete the proof of our results. We will proceed by showing the so-called "weak sequential stability" of solution families. Namely, we let $(\varphi_n, \mu_n, \sigma_n, u_n)$, $n \in \mathbb{N}$, be a sequence of solutions complying with the a-priori estimates obtained before uniformly with respect to n. Then, proving weak sequential stability means that at least a (non-relabelled) subsequence has to converge to a quadruple $(\varphi, \mu, \sigma, u)$ solving (1.1)-(1.5) in the sense specified in the statements of the existence theorems.

Of course, in order for the argument to be fully rigorous, we should rather assume that the sequence $(\varphi_n, \mu_n, \sigma_n, \boldsymbol{u}_n)$ solves some approximation of the system (for instance the one sketched in Section 6 below), with the approximating parameters varying with n in a suitable way. However, for the sake of simplicity, we prefer to assume that $(\varphi_n, \mu_n, \sigma_n, \boldsymbol{u}_n)$ complies directly with the "limit", or "original", system, possibly written in the "strong" form (1.1)-(1.5). We actually believe that such a simplified method permits us to focus on the substantial points of the compactness argument. In the next section, we will explain how the procedure could be adapted in order to fit the proposed approximation scheme.

We will mostly focus on the (more difficult) case of Theorem 2.3, referring to the situation where d = 3 and p > 12/11 (at the end we will sketch the differences occurring for d = 2 in the setting of Theorem 2.1). Moreover, in order to achieve the limit $n \nearrow \infty$, we will use only the information resulting from the mass conservation, the minimum principle, the energy estimate, and the most general version of the entropy bound, namely (cf. (4.28) and (4.26))

$$\|\sigma_n^{p/2}\|_{L^2(0,T;V)} + \|\varphi_n\|_{L^{\frac{12p}{12-5p}}(0,T;H^2(\Omega))} \le c,$$
(5.1)

where the exponent $\frac{12p}{12-5p}$ is replaced by 4 for $p \ge 3/2$, cf. (4.30)). Indeed, the additional information resulting from more refined versions of the entropy estimate, and in particular the argument performed in Subsec. 4.3, will produce additional regularity of weak solution, but is not essential for taking the limit. Note also that, from now on, we will stress the dependence on n in the estimates. Moreover, in order to prepare the asymptotic limit $\varepsilon \searrow 0$, we will note by c (respectively c_{ε}) the positive constants occurring in the estimates which are independent of ε (respectively, depending on ε).

That said, we start testing (1.1) by $\xi \in V$ and integrating by parts so to obtain the weak formulation (2.25). Next, by (3.28) and the uniform boundedness of φ_n there follows

$$\|\varphi_n u_n\|_{L^2(0,T;H)} \le c. \tag{5.2}$$

Hence, comparing terms in (1.1), and using (3.39), (2.3) and (5.2), we obtain

$$\|\varphi_{n,t}\|_{L^2(0,T;V')} \le c. \tag{5.3}$$

Let us now test (1.3) by $\eta \in W^{1,\infty}(\Omega)$ (actually, in the procedure it will be clear that the conditions on η can be weakened, as specified in the statements) and integrate by parts so to deduce the weak formulation (2.27). Recall in particular that we have set

$$\boldsymbol{H}(\sigma_n, \varphi_n) := \alpha^{1/2}(\sigma_n) \nabla(\gamma(\sigma_n) - \chi \varphi_n), \tag{5.4}$$

and, with this notation, (3.30) corresponds to

$$\|\boldsymbol{H}(\sigma_n,\varphi_n)\|_{L^2(0,T;H)} \le c.$$
(5.5)

Then, observing that (we assume with no loss of generality $p \in (1, 2)$, the analogue being trivial in the case p = 2)

$$\|\alpha^{1/2}(\sigma_n)\|_{L^{\infty}(0,T;L^{\frac{2p}{2-p}}(\Omega))} \le \|\sigma_n^{\frac{2-p}{2}}\|_{L^{\infty}(0,T;L^{\frac{2p}{2-p}}(\Omega))} \le \|\sigma_n\|_{L^{\infty}(0,T;L^p(\Omega))}^{\frac{2-p}{2}} \le c,$$
(5.6)

combining (5.5) and (5.6), it is easy to deduce

$$\|\alpha^{1/2}(\sigma_n)H(\sigma_n,\varphi_n)\|_{L^2(0,T;L^p(\Omega))} \le c.$$
(5.7)

Note that the above only relies on the "energy" estimate: using the "entropy" regularity, we could actually get some better information.

Next, we observe that (3.29) and the first (5.1), thanks to interpolation and Sobolev's embeddings, imply

$$\|\sigma_n\|_{L^{\frac{5p}{3}}(Q)} + \|\sigma_n\|_{L^{\frac{4p}{6-3p}}(0,T;H)} + \|\sigma_n\|_{L^2(0,T;L^{\frac{3p}{3-p}}(\Omega))} \le c.$$
(5.8)

In particular, we may notice that the last two exponents depending on p are strictly larger than 2 if and only if p is strictly larger than 6/5.

Next, we observe that, from (3.25), the second (5.1) (note that $\frac{12p}{12-5p} > 2$ for p > 12/11) and interpolation, there follows at least

$$\|\nabla\varphi_n\|_{L^{10/3}(Q)} \le c \big(\|\varphi_n\|_{L^{\infty}(0,T;V)} + \|\varphi_n\|_{L^2(0,T;H^2(\Omega))}\big) \le c.$$
(5.9)

Moreover, using again (5.1) with the first (5.6), we get

$$\|\alpha(\sigma_n)\nabla\varphi_n\|_{L^2(0,T;L^{\frac{6p}{12-5p}}(\Omega))} \le c\|\alpha(\sigma_n)\|_{L^{\infty}(0,T;L^{\frac{p}{2-p}}(\Omega))}\|\nabla\varphi_n\|_{L^2(0,T;L^6(\Omega))} \le c.$$
 (5.10)

Since $\frac{6p}{12-5p} > 1$ for p > 12/11, the above, combined with (5.7), implies

$$\begin{aligned} \|\nabla\sigma_n\|_{L^2(0,T;L^S(\Omega))} &\leq c \left(\|\alpha^{1/2}(\sigma_n) \boldsymbol{H}(\sigma_n,\varphi_n)\|_{L^2(0,T;L^p(\Omega))} + \|\alpha(\sigma_n)\nabla\varphi_n\|_{L^2(0,T;L^{\frac{6p}{12-5p}}(\Omega))} \right) \\ &\leq c, \qquad \text{where} \ S = \min\left\{ \frac{6p}{12-5p}, p \right\} > 1. \end{aligned}$$
(5.11)

We now consider the convection term and note that its control is the only step where the condition p > 12/11 is not sufficient to obtain an estimate which also allows for the Brinkman-to-Darcy limit $\varepsilon \searrow 0$. Actually, as $\varepsilon > 0$ is fixed, we may use (3.33), which, combined with the *third* (5.8), thanks to Sobolev's embeddings, implies

$$\left\|\sigma_n \boldsymbol{u}_n\right\|_{L^{Z}(0,T;L^{\frac{4}{3}}(\Omega))} \le c_{\varepsilon},\tag{5.12}$$

where Z is some exponent strictly larger than 1. To deduce the above we used that, for p = 12/11, one has $\frac{3p}{3-p} = 12/7$, so that 1/6 + 7/12 = 3/4. Since in fact p > 12/11, changing a bit the interpolation exponents in (5.8), we see that Z can be taken strictly larger than 1 in (5.12).

On the other hand, as we deal with the Darcy limit, in order to use the first (3.28), we need to know that

$$\|\sigma_n\|_{L^{Z_1}(Q)} \le c, \text{ for some } Z_1 > 2.$$
 (5.13)

with c independent of ε . Then, interpolating between the last two conditions in (5.8), it is clear that such an exponent can be found if and only if p > 6/5, as required in the statement of Theorem 2.3.

Next, we notice that, from (2.10) and (3.29), there follows

$$\|b(\sigma_n,\varphi_n)\|_{L^{\infty}(0,T;L^p(\Omega))} \le c.$$
(5.14)

Let us now write (1.3) for the approximate sequence and test it by $\eta \in W^{1,\infty}(\Omega)$ (as already noted, this condition may be weakened). Then, using (5.12), (5.10), (5.11) and (5.14), one deduces

$$\|\sigma_{n,t}\|_{L^{Z}(0,T;W^{1,R}(\Omega)')} \le c_{\varepsilon}, \text{ where } Z > 1,$$
 (5.15)

and c_{ε} can be taken independent of ε when either p > 6/5 or $\sigma_0 \in L^q(\Omega)$ for q > 6/5.

Combining (5.15) with (3.29) and (5.11) allows us to apply the Aubin-Lions lemma, which gives

$$\sigma_n \to \sigma \quad \text{strongly in } L^2(0,T;W^{1-\iota,S}(\Omega)),$$
(5.16)

where S is as in (5.11) and $\iota > 0$ is arbitrarily small. Here and below, all convergence relations are to be intended to hold up to the extraction of non-relabelled subsequences of $n \nearrow \infty$. Note that (5.16) implies in particular the pointwise (a.e.) convergence $\sigma_n \to \sigma$.

Proving strong convergence of φ_n is, of course, much simpler. Using (5.3), (3.25) and (5.1), the Aubin-Lions lemma actually guarantees

$$\varphi_n \to \varphi$$
 strongly in $C^0([0,T]; H^{1-\iota}(\Omega)) \cap L^{\frac{12p}{12-5p}}(0,T; H^{2-\iota}(\Omega))$ for every $\iota > 0.$ (5.17)

Again, the above implies in particular almost everywhere convergence both of $\varphi_n \to \varphi$ and of $\nabla \varphi_n \to \nabla \varphi$. This allows us to manage the source terms in (1.1) and (1.3). Actually, using assumptions (2.3) and (2.10) with the first (5.8) (notice that 5p/3 > 20/11 if p > 12/11), we deduce at least

$$h(\sigma_n, \varphi_n) \to h(\sigma, \varphi)$$
 strongly in $L^P(Q)$ for every $P \in [1, \infty)$, (5.18)

$$b(\sigma_n, \varphi_n) \to b(\sigma, \varphi)$$
 strongly in $L^{20/11}(Q)$. (5.19)

We now move to the Korteweg terms in the Brinkman relation (1.4). First of all, notice that, from (5.17), one deduces in particular

$$\nabla \varphi_n \to \nabla \varphi$$
 strongly in $L^4(0,T;H)$. (5.20)

Combining this fact with the weak convergence resulting from (3.39), we then infer

$$\mu_n \nabla \varphi_n \to \mu \nabla \varphi \quad \text{weakly in } L^{4/3}(0,T;L^{3/2}(\Omega)).$$
(5.21)

Next, let us notice that, combining (3.32) with the pointwise convergence $\varphi_n \to \varphi$, there follows

$$\varphi_n \to \varphi$$
 strongly in $L^P(Q)$ for every $P \in [1, \infty)$. (5.22)

The above relation and the weak convergence resulting from (5.11) imply

$$\varphi_n \nabla \sigma_n \to \varphi \nabla \sigma \quad \text{weakly in } L^2(0,T;L^S(\Omega)).$$
 (5.23)

Concerning the cross-diffusion term, as a consequence of the pointwise convergence of σ_n and $\nabla \varphi_n$ and of (5.10), we deduce (at least)

$$\alpha(\sigma_n)\nabla\varphi_n \to \alpha(\sigma)\nabla\varphi$$
 strongly in $L^1(Q)$. (5.24)

Finally, we consider the transport terms, both of which are more conveniently integrated by parts. As already observed, the transport term in (1.1) can be treated easily: indeed, using (5.2), the weak L^2 -convergence of u_n , and the pointwise convergence of φ_n , there follows

$$\boldsymbol{u}_n \varphi_n \to \boldsymbol{u} \varphi$$
 weakly in $L^2(0,T;H)$. (5.25)

Concerning the transport term in (1.3), a similar argument based on (5.12) (notice that, once more, Z > 1 is essential) permits us to obtain at least

$$\boldsymbol{u}_n \sigma_n \to \boldsymbol{u}\sigma$$
 weakly in $L^Z(Q)$ for some $Z > 1.$ (5.26)

As already observed, for $p \leq 6/5$, the argument works only as far as $\varepsilon > 0$ is fixed.

The last term we need to manage is the logarithmic one occurring in (1.2). As is customary, to deal with it we will use some tool from the theory of maximal monotone operators. To this aim, we recall that $\beta(r) := f(r) + \lambda r$ denotes the "monotone part" of f. Hence, of course, it is enough to prove that $\beta(\varphi_n) \to \beta(\varphi)$ in a suitable sense. It is worth observing that, at least formally, this identification can be obtained just using the information resulting from the energy estimate. Indeed, recalling (3.40) and using that p is strictly larger than 1, one deduces

$$\beta(\varphi_n) \to \xi$$
 weakly in $L^p(Q)$. (5.27)

Then, to identify ξ as $\beta(\varphi)$, one may notice that (5.22) holds in particular for P = p' (the conjugate exponent to p). Hence, using the strong-weak closedness of the maximal monotone operator induced by β and acting from $L^{p'}(Q)$ to $L^{p}(Q)$ (which are reflexive Banach spaces in duality), one readily has

$$\beta(\varphi_n) \to \beta(\varphi)$$
 weakly in $L^p(Q)$. (5.28)

We notice, however, that adapting this part of the argument to the proposed approximation scheme will require some amount of additional work (see the next section for details).

As a consequence of the above procedure, all nonlinear terms in the weak formulation (2.25)-(2.28) pass to the expected limits. To conclude, we observe that, from (3.29), (5.15) and a generalized version of the Aubin-Lions lemma, there follows

$$\sigma_n \to \sigma$$
 strongly in $C^0([0,T];X),$ (5.29)

where X is a negative order Sobolev space such that

$$L^{p}(\Omega) \subset X \subset W^{1,\infty}(\Omega)', \tag{5.30}$$

the first inclusion being compact and the second being continuous. The existence of such a space is clear as one considers that $W^{1,\infty}(\Omega) \subset L^{p'}(\Omega)$. Then, using (5.29) and the first (5.17) it is immediate to recover the initial data in the sense of (2.11).

End of proof of Theorem 2.3. To conclude the proof, we first have to check the required regularity properties (2.37)-(2.44). Actually, most of them are direct consequences of the "energy" bounds (3.25)-(3.31), (3.33), (3.39) and of the "entropy" bounds (4.26), (4.28) (or (4.30)), and (in the case when we

assume q > p) (4.53), (4.54) (or (4.55)). In particular, the latter conditions imply the third of (2.37), with the exponent P_0 specified in (2.45).

Next, concerning the last regularity condition in (2.37) and the last in (2.38), we mimick the interpolation argument leading to (5.8), where we can take q in place of p. This gives

$$\|\sigma_n\|_{L^2(0,T;L^{\frac{3q}{3-q}}(\Omega))} \le c.$$
(5.31)

Applying to (1.2) the usual elliptic regularity argument, we then complete (2.37) and (2.38). Concerning the regularity of σ , (2.42) follows directly from (3.29) and (4.26) if q = p (i.e., we do not need to perform the regularity argument in Subsec. 4.3) and from (4.53) if q > p. Regarding (2.41), the second condition follows directly from (5.11), while the first property is a consequence of (5.7) and (5.12); in particular, this regularity conditions allows the use of $\eta \in W^{1,R}(\Omega)$, with R specified by (2.45), as a test function in (2.27).

Next, let us move to the Brinkman-to-Darcy limit. To this aim, let $(\varphi_{\varepsilon}, \mu_{\varepsilon}, \sigma_{\varepsilon}, \boldsymbol{u}_{\varepsilon})$ be a family of weak solutions solving the problem at the level ε and let us let $\varepsilon \searrow 0$. Then, we observe that most of the estimates are uniform with respect to ε and, consequently, the argument to pass to the limit mostly follows the procedure used for letting $n \nearrow \infty$. The only relevant difference regarding the fact that, from (3.28) we can only deduce that (for a subsequence of $\varepsilon \searrow 0$)

$$\boldsymbol{u}_{\varepsilon} \to \boldsymbol{u} \quad \text{weakly in } L^2(0,T;H).$$
 (5.32)

On the other hand, using (5.13), which holds thanks to the assumption p > 6/5, we can now deduce

$$\sigma_{\varepsilon} \to \sigma \quad \text{strongly in } L^2(0,T;H).$$
 (5.33)

Using (5.32), (5.33) and (5.13) again, we then arrive at

$$\sigma_{\varepsilon} \boldsymbol{u}_{\varepsilon} \to \sigma \boldsymbol{u} \quad \text{weakly in } L^{Z}(Q),$$

$$(5.34)$$

for some Z > 1. This allows, from one side, to take the limit of the transport term in (2.27) and, from the other side, to achieve the first convergence relation in (2.50), which basically concludes the proof.

It is finally worth observing that condition (2.43) follows from (3.31), which also implies a uniform in time control of $\ln \sigma$ in the space of signed measures on Ω .

End of proof of Theorem 2.1. In the two-dimensional case, the procedure is simpler and we can rely on the stronger estimates obtained in Subsection 4.1, which imply the desired regularity conditions (2.18)-(2.24), as a direct check shows. In particular, we may notice that the last (2.18) and the last (2.19) are direct consequences of (4.12). Moreover, for fixed $\varepsilon > 0$, combining (2.21) with the second (2.22) and using Sobolev's embeddings there follows

$$\|\sigma \boldsymbol{u}\|_{L^2(0,T;L^{R'}(\Omega))} \le c_{R',\varepsilon}, \quad \text{for every } R' \in [1,2).$$

$$(5.35)$$

Using this piece of information and comparing terms in (1.3), it is then easy to get the first (2.22), with $R \in (2, \infty]$ being the conjugate exponent to R'.

Next, for what concerns the limit $\varepsilon \searrow 0$, we let as before $(\varphi_{\varepsilon}, \mu_{\varepsilon}, \sigma_{\varepsilon}, u_{\varepsilon})$ be a family of weak solutions depending on ε . Then, we can mostly proceed as in the three-dimensional case. The only thing we have to notice regards the first (2.32). Indeed, by interpolation, (4.11) guarantees

$$\|\sigma_{\varepsilon}\|_{L^4(Q)} \le c \big(\|\sigma_{\varepsilon}\|_{L^{\infty}(0,T;H)} + \|\sigma_{\varepsilon}\|_{L^2(0,T;V)}\big) \le c.$$
(5.36)

Combining this with (2.33) (which is a consequence of the first (3.28)), we then have

$$\|\sigma_{\varepsilon} \boldsymbol{u}_{\varepsilon}\|_{L^{4/3}(Q)} \le c, \tag{5.37}$$

whence follows the first (2.32). Finally, we observe that, relation (2.27) makes sense for every $\eta \in W^{1,R}(\Omega)$, R > 2. This is actually a consequence of (5.35) combined with the fact that (5.36) and (2.18) imply at least

$$\|\nabla\sigma\|_{L^{2}(0,T;H)} + \|\alpha(\sigma)\nabla\varphi\|_{L^{2}(0,T;H)} \le c.$$
(5.38)

Hence the cross-diffusion term in (2.27) admits $\eta \in W^{1,R}(\Omega)$ as a test function (actually it would also admit $\eta \in V$) and it could be written splitting it into its two summands.

Remark 5.1. Adapting the proof of Theorem 2.3, one may easily check that, for d = 2, existence of solutions for d = 2 can be achieved for every under the sole condition (2.13) with p > 1. Moreover, in this case also the limit Brinkman-to-Darcy (i.e., $\varepsilon \searrow 0$) can be performed for any σ_0 complying with (2.13) and for any p > 1 (compare with assumption (2.46) in the three-dimensional case). Actually, for d = 2 in place of the first (5.8) we get

$$\|\sigma_{\varepsilon}^{p/2}\|_{L^{4}(Q)} \le c \left(\|\sigma_{\varepsilon}^{p/2}\|_{L^{\infty}(0,T;H)} + \|\sigma_{\varepsilon}^{p/2}\|_{L^{2}(0,T;V)}\right) \le c,$$
(5.39)

where we emphasized the dependence on ε in the notation. Hence, being p is strictly larger than 1, the above implies

$$\sigma_{\varepsilon} \to \sigma \quad \text{strongly in} \ L^2(Q), \tag{5.40}$$

so that, using the weak convergence resulting from the first (3.28) (which is independent of ε), one deduces (at least)

$$\sigma_{\varepsilon} \boldsymbol{u}_{\varepsilon} \to \sigma \boldsymbol{u} \quad \text{weakly in } L^1(Q),$$

$$(5.41)$$

which is the key property to achieve the Darcy limit. As noted in the introduction, we are planning to analyze in detail the two-dimensional case in a forthcoming work.

6 Approximation

In this part we sketch a regularization of system (1.1)-(1.5) which may be exploited in order to justify the a priori estimates - compactness argument used in the existence proof. In view of the many nonlinearities involved, we will not present the approximation scheme in full detail, as that would be a rather lengthy and technical procedure. Indeed, we point out that, both for the Cahn-Hilliard model and for the chemotaxis models related to the Keller-Segel system, the literature dealing with (analytical and numerical) approximations is very vast: with no claim of exhaustivity, we refer the reader to [6, 16, 30, 33] for the Cahn-Hilliard model and to [28, 32, 34] for chemotaxis models. Referring to the quoted papers for more details, here, we will just focus on the "new" aspects of the regularization arising in connection with our specific problem. We also point out that some parts of our argument are somehow reminiscent of the procedure devised in [27], to which we refer the reader for additional considerations.

That said, we indicate by $n \in \mathbb{N}$ the regularization parameter, intended to go to infinity in the limit, and propose a regularization of system (1.1)-(1.5) depending on n, for which existence may be proved by standard methods (e.g., a Faedo-Galerkin scheme possibly complemented by a fixed point argument). The approximation will be designed so to be compatible with the a-priori estimates obtained before. To begin, we need to smooth out the singular (logarithmic) term $f(\varphi)$ so to ensure the applicability of some local existence theorem. Actually, in a Faedo-Galerkin scheme one needs that all the elements of the Galerkin base belong to the domain of the potential, which is not true if F is given by (1.6) (and so it takes finite values only on [-1,1]). To regularize f, we observe that it is sufficient to act on its "monotone part" β . Using some tools from the theory of maximal monotone operators (cf. the monographs [2, 4] for some background), in Lemma 6.1 below we will explicitly describe a concrete example of a "smooth" approximating family β_n , with controlled growth at infinity, suitably converging to β as $n \searrow \infty$. Once β_n is given, we define $f_n := \beta_n - \lambda \operatorname{Id}$, Id denoting the identity function. Then, noting as F_n a suitable primitive of f_n (we may assume F_n to be normalized so that $F_n(0) = F(0) = 0$), one also obtains an approximation of the potential.

Once F_n is given, we can introduce a regularized version of the energy functional (3.9), namely

$$\mathcal{E}_n(\varphi,\sigma) := \frac{1}{2} \|\nabla\varphi\|^2 + \frac{1}{2n} \|\Delta\varphi\|^2 + \int_{\Omega} F_n(\varphi) + \int_{\Omega} \left(\widehat{\gamma}(\sigma) - \chi\sigma\varphi\right).$$
(6.1)

Here we have replaced F with F_n and we have added a further regularizing term depending on the Laplacian of φ . The reason for such a choice stands in the fact that, as f is smoothed out, the uniform boundedness of φ , which is extensively used in the estimates (cf. (3.32)) may be lost in the approximation. Adding the second-order term, we somehow compensate this loss of coercivity.

Then, the regularization of the energy, in turn, generates a regularized version of the system, which may be stated as

$$\varphi_{n,t} + \boldsymbol{u}_n \cdot \nabla \varphi_n - \Delta \mu_n = h(\sigma_n, \varphi_n) - \ell \varphi_n, \qquad (6.2)$$

$$\mu_n = \frac{1}{n} \Delta^2 \varphi_n - \Delta \varphi_n + f_n(\varphi_n) - \chi \sigma_n, \tag{6.3}$$

$$\sigma_{n,t} + \boldsymbol{u}_n \cdot \nabla \sigma_n - \Delta \sigma_n + \chi \operatorname{div}(\alpha(\sigma_n) \nabla \varphi_n) = b(\sigma_n, \varphi_n), \tag{6.4}$$

$$-\varepsilon \operatorname{div}(D\boldsymbol{u}_n) + \boldsymbol{u}_n = \nabla \pi_n + \mu_n \nabla \varphi_n - \chi \varphi_n \nabla \sigma_n, \qquad (6.5)$$

$$\operatorname{div} \boldsymbol{u}_n = 0. \tag{6.6}$$

Observe that, due to the occurrence of the fourth order term in (6.3), an additional boundary condition is needed. We shall assume

$$\partial_{\boldsymbol{n}} \Delta \varphi_{\boldsymbol{n}} = 0 \quad \text{on } \Gamma \times (0, T).$$
 (6.7)

Actually, (6.2)-(6.3) is a particular version of the so-called "sixth-order Cahn-Hilliard system", which has been extensively studied in the literature (we quote, among the other contributions [21, 29, 30]) and for which the boundary condition (6.7) is a rather standard choice.

Notice also that there is no need of regularizing the functions h and b, since assumptions (2.3) and (2.10) already guarantee sufficient smoothness properties as well as a controlled growth for large values of σ and φ . We also remark that the expression of equation (6.4) has not been modified compared to (1.3); in particular, the equivalent of the "minimum principle" (3.31) still holds in the approximation.

We now see how the above regularization impacts on the various parts of the existence proof presented in the previous sections.

Energy estimate. We first observe that the balance of mass (3.2) as well as the "weak" minimum principle (3.3) still hold in the approximation. Concerning the energy bound, the argument outlined in Section 3 is fully compatible with the regularization; however, some differences arise in the outcome of the procedure. To explain this point, we first observe that, as we reproduce the estimate by working on (6.2)-(6.6), we still obtain the analogue of (3.24). Applying Grönwall's lemma, and neglecting for simplicity the contribution of the dissipative terms on the left-hand side, this leads to the bound

$$\left\| \mathcal{E}_{n}(\varphi_{n},\sigma_{n}) + \|\ln\sigma_{n}\|_{L^{1}(\Omega)} \right\|_{L^{\infty}(0,T)} \leq c_{T} \left| \mathcal{E}_{n}(\varphi_{0,n},\sigma_{0,n}) + \|\ln\sigma_{0,n}\|_{L^{1}(\Omega)} \right|.$$
(6.8)

Here, $(\varphi_{0,n}, \sigma_{0,n})$ are suitable regularizations of the initial data designed in such a way that the righthand side above remains bounded uniformly with respect to n and $(\varphi_{0,n}, \sigma_{0,n}) \rightarrow (\varphi_0, \sigma_0)$ in a suitable sense. Actually, one may directly take $\sigma_{0,n} \equiv \sigma_0$, while smoothing out φ_0 is a bit more involved and can be performed by following the lines of the procedure devised in [29] or referring to other papers dealing with the Cahn-Hilliard system with singular potential. In particular, one may ensure that the analogue of (2.12) holds uniformly with respect to n and that

$$\|\varphi_{0,n}\|_{H^2(\Omega)} \le cn^{1/2},$$
(6.9)

so that the regularized initial energy on the right-hand side of (6.8) is controlled uniformly in n.

The effects of the regularization also have some impact on (3.25)-(3.32). In particular, as F_n is defined over the whole real line, there is no reason for (3.32) to hold. Nevertheless, if we choose β_n as in Lemma 6.1 below, i.e., we ask that

$$\int_{\Omega} F_n(\varphi) \ge \kappa \|\varphi\|_{L^{q_0}(\Omega)}^{q_0} - c, \qquad (6.10)$$

for some constants $\kappa > 0$ and $c \ge 0$ independent of n and for $q_0 > p'$ (p' being the conjugate exponent to p), then the energy functional keeps some coercivity. Indeed, recalling (3.6), one has

$$\int_{\Omega} F_{n}(\varphi) + \int_{\Omega} \left(\widehat{\gamma}(\sigma) - \chi \sigma \varphi \right) \\
\geq \kappa \|\varphi\|_{L^{q_{0}}(\Omega)}^{q_{0}} + \frac{1}{p(p-1)} \|\sigma\|_{L^{p}(\Omega)}^{p} - c \left(1 + \|\sigma\|_{L^{1}(\Omega)} + \int_{\Omega} |\sigma \ln \sigma| + \|\sigma\|_{L^{p}(\Omega)} \|\varphi\|_{L^{p'}(\Omega)} \right) \\
\geq \frac{\kappa}{2} \|\varphi\|_{L^{q_{0}}(\Omega)}^{q_{0}} + \frac{1}{2p(p-1)} \|\sigma\|_{L^{p}(\Omega)}^{p} - c,$$
(6.11)

the last inequality following by properly applying Young's inequality and using that $q_0 > p'$. The above relation entails that the approximate energy is *uniformly coercive* with respect to n. As a consequence, when performing the approximate energy estimate, at least a uniform control of φ in $L^{\infty}(0,T; L^{q_0}(\Omega))$ is achieved. Moreover, due to the additional Laplacian in (6.1), (6.8) entails the additional estimate

$$\|\varphi_n\|_{L^{\infty}(0,T;H^2(\Omega))} \le cn^{1/2},\tag{6.12}$$

with the constant c independent of n.

Entropy estimates. We start observing that the first, and simpler, version of the argument, as presented in Subsec. 4.1 for the two-dimensional case, basically requires no significant variation. On the other hand, for the argument presented in Subsec. 4.2, some discussion is in order. Indeed, the procedure uses in an essential way the outcome of the Gagliardo-Nirenberg inequality (4.13), which, in turn, subsumes the uniform boundedness of φ stated in (3.32) as a consequence of the choice of the "singular" potential (1.6), which is however lost in the approximation as F is smoothed out. Actually, at the regularized level, the best information we have is given by the combination of (6.11) and (6.12). In particular, the latter relation implies, via Sobolev's embeddings, that, at fixed n, φ_n is bounded in the uniform norm; however, at least in principle, this bound needs not be uniform with respect to n.

Fortunately, in the next lemma we may prove that β_n can be constructed so to ensure the boundedness of φ_n independently of the approximation parameter. We believe this property to have an independent interest; hence we present it as a lemma so to facilitate future reference. We refer again to the monographs [2, 4] for the terminology as well as to [1, Chap. 3] for what concerns the concepts of graph-convergence and G-convergence.

Lemma 6.1. Let Ω be a smooth, bounded domain of \mathbb{R}^d , $d \in \{2,3\}$, and $\beta : \mathbb{R} \to 2^{\mathbb{R}}$ be a maximal monotone graph, with the domain $D(\beta)$ satisfying $\overline{D(\beta)} = [-1,1]$. Let β be normalized in such a way that $\beta(0) \ni 0$. Let also $q_0 \in (2,\infty)$ be a given exponent. Then, there exists an approximating family of monotone and locally Lipschitz continuous functions $\beta_n : \mathbb{R} \to \mathbb{R}$, with $\beta_n(0) = 0$, such that, defining

$$\widehat{\beta}_n(s) := \int_0^s \beta_n(r) \, \mathrm{d}r, \tag{6.13}$$

then it follows that

$$K_n(1+|s|^{q_0}) \ge \widehat{\beta}_n(s) \ge \kappa |s|^{q_0} - c \quad \text{for every } n \in \mathbb{N}, \ s \in \mathbb{R},$$
(6.14)

for suitable constants $\kappa > 0$ and $c \ge 0$ (independent of n) and $K_n > 0$ (depending on n and diverging as $n \nearrow \infty$). Moreover, for every constant $c_0 \ge 0$ there exists a constant $\overline{C} > 0$ independent of n and such that, if $\{v_n\}$ is a sequence of functions in $H^2_n(\Omega)$ satisfying

$$\|v_n\|_V + n^{-\frac{1}{2}} \|v_n\|_{H^2(\Omega)} + \int_{\Omega} \widehat{\beta}_n(v_n) \le c_0,$$
(6.15)

then $\{v_n\}$ also satisfies

$$\|v_n\|_{L^{\infty}(\Omega)} \le \overline{C}.\tag{6.16}$$

Finally, the family β_n converges to β in the sense of graphs (or G-convergence) in $\mathbb{R} \times \mathbb{R}$.

PROOF. We denote as $\beta_{0,n}$ the Yosida approximation of β of index n^{-1} . Moreover, we define the following sequence of functions:

$$j_n(s) := \begin{cases} 0 & \text{if } s \in [-1,1], \\ q_0 n^{8q_0} (s-1)^{q_0-1} & \text{if } s > 1, \\ -q_0 n^{8q_0} |s+1|^{q_0-1} & \text{if } s < -1. \end{cases}$$
(6.17)

Using that, for every s > 1 (respectively, for every s < -1) $j_n(s)$ is monotone increasing (respectively, is monotone decreasing) with respect to n and diverges to $+\infty$ (respectively, $-\infty$), it is easy to check that the sequence $\{j_n\}$ converges, in the sense of graphs, to the graph $\partial I_{[-1,1]}$ (the subdifferential of the indicator function of [-1,1]). To be precise, this property may be shown by proving that the antiderivatives of j_n vanishing at 0 monotonically converge to $I_{[-1,1]}$ and applying [1, Thm. 3.20 and Thm. 3.66]. Then, setting $\beta_n(s) := \beta_{0,n}(s) + j_n(s)$, we can also see that β_n converges to β in the sense of graphs. In addition, defining $\hat{\beta}_n$ as in (6.13), it is clear that properties (6.14) hold. Moreover, one has $\beta_n(0) = 0$ for every $n \in \mathbb{N}$ thanks to the properties of Yosida approximations.

Let now (6.15) hold. Then, computing explicitly the primitive of j_n , one can easily realize that

$$c_{0} \ge \int_{\Omega} \widehat{\beta}_{n}(v_{n}) \ge n^{8q_{0}} \bigg(\int_{\Omega} \big((v_{n} - 1)_{+} \big)^{q_{0}} + \int_{\Omega} \big((v_{n} + 1)_{-} \big)^{q_{0}} \bigg).$$
(6.18)

Let now $T \in C^2(\mathbb{R};\mathbb{R})$ a truncation operator satisfying the following properties:

$$T(s) \equiv 0$$
 for all $s \in [-1,1]$, $T(s) = s - 2$ for all $s \ge 3$, $T(s) = s + 2$ for all $s \le -3$, (6.19)

$$|T(s)| \le 1 \text{ for all } s \in [-3,3], \quad 0 \le T'(s) \le 1 \text{ for all } s \in \mathbb{R}, \quad |T''(s)| \le c \text{ for all } s \in \mathbb{R}, \quad (6.20)$$

and for some c > 0. It is not difficult to construct explicitly a function T satisfying all the above properties. Noting that

$$\Delta T(v_n) = T'(v_n)\Delta v_n + T''(v_n)|\nabla v_n|^2, \qquad (6.21)$$

using Sobolev's embedding and elliptic regularity (recall that $v_n \in H^2_n(\Omega)$) one deduces

$$\|\Delta T(v_n)\| \le c \|\Delta v_n\| + c \|\nabla v_n\|_{L^4(\Omega)}^2 \le c (\|v_n\|_V^2 + \|\Delta v_n\|^2).$$
(6.22)

Hence, recalling (6.15) and using again $v_n \in H^2_n(\Omega)$, there follows

$$\|v_n\|_{H^2(\Omega)} \le c(\|v_n\| + \|\Delta T(v_n)\|) \le c_1 n,$$
(6.23)

where $c_1 > 0$ depends on c_0 and embedding constants, but is independent of n.

Next, let us observe that

$$|T(v_n)| \le (v_n - 1)_+ + (v_n + 1)_- \quad \text{a.e. in } \Omega.$$
(6.24)

Hence, from (6.18) we deduce

$$\|T(v_n)\|_{L^{q_0}(\Omega)}^{q_0} \le c \big(\|(v_n-1)_+\|_{L^{q_0}(\Omega)}^{q_0} + \|(v_n+1)_-\|_{L^{q_0}(\Omega)}^{q_0}\big) \le cn^{-8q_0}.$$
(6.25)

As a consequence,

$$|T(v_n)|| \le c ||T(v_n)||_{L^{q_0}(\Omega)} \le cn^{-8}.$$
(6.26)

Thus, by interpolation (we consider the more difficult case d = 3, the two-dimensional setting working with better exponents)

$$||T(v_n)||_{L^{\infty}(\Omega)} \leq c ||T(v_n)||_{H^{7/4}(\Omega)} \leq c ||T(v_n)||_{H^2(\Omega)}^{7/8} ||T(v_n)||^{1/8}$$
$$\leq c n^{7/8} n^{-1} = c n^{-1/8}.$$
(6.27)

Recalling (6.19), the thesis (6.16) follows.

With the lemma at disposal, we may notice that, from the approximate energy estimate (6.8), there follows in particular that

$$n^{-\frac{1}{2}} \|\Delta \varphi_n\|_{L^{\infty}(0,T;H)} + \|F_n(\varphi_n)\|_{L^{\infty}(0,T;L^1(\Omega))} \le c.$$
(6.28)

Hence, applying the lemma, one obtains

$$\|\varphi_n\|_{L^{\infty}(0,T;L^{\infty}(\Omega))} \le c, \tag{6.29}$$

with c independent of n, which allows for the use of the uniform boundness of φ even when the estimates are performed at the regularized level. The remainder of the entropy argument seems to require no further variations.

Passage to the limit. Here, the main point is concerned with the logarithmic term. Actually, in the previous part (cf., in particular, (5.28)), the identification of the limit was obtained by relying on the strong-weak closedness of the maximal monotone operator induced by β (the monotone part of f) from $L^{p'}(Q)$ to $L^{p}(Q)$ (which are reflexive Banach spaces in duality). This property may be stated as

$$\varphi_n \to \varphi \text{ strongly in } L^{p'}(Q) \text{ and } \beta(\varphi_n) \to \xi \text{ weakly in } L^p(Q) \implies \xi \equiv \beta(\varphi),$$
 (6.30)

the last equality holding in $L^p(Q)$, hence almost everywhere.

However, as the system is approximated by (6.2)-(6.6), some differences arise. First of all, β is now replaced by a varying family $\{\beta_n\}$. In itself, this would not be a problem because the fact that β_n converges to β in the sense of graphs (in \mathbb{R}) implies, in turn, that the maximal monotone operators induced by β_n G-converge (we refer, again, to [1, Chap. 3] for further details) to the the maximal monotone operator induced by β . Such a property would actually imply

$$\varphi_n \to \varphi$$
 strongly in $L^{p'}(Q)$ and $\beta_n(\varphi_n) \to \xi$ weakly in $L^p(Q) \implies \xi \equiv \beta(\varphi)$. (6.31)

There is, however, a further issue which prevents the application of the above argument, related to the presence of the fourth-order elliptic term in (6.3). Indeed, it turns out that the operator

$$B_n(v) := \frac{1}{n} \Delta^2 v + \beta_n(v), \qquad (6.32)$$

complemented with the boundary conditions (6.7), is not maximal monotone in the duality between $L^{p'}$ and L^{p} .

To explain this fact, we go back to the Hilbert setting (indeed, in our modified version of the argument we will be able to use Hilbert spaces) and neglect, for simplicity, the effects of the time variable. What happens is that, while the operators $v \mapsto n^{-1}\Delta^2 v$ and $v \mapsto \beta_n(v)$ (with proper domains) are separately maximal monotone in H, their sum is not. The reason stants in the lack of the so-called "angle condition". Roughly speaking, this corresponds to the fact that the L^2 -scalar product of $\Delta^2 v$ and $\beta_n(v)$ has no sign properties. In other words, an L^2 -estimate of the form

$$\left\|n^{-1}\Delta^2 v_n + \beta_n(v_n)\right\| \le c,\tag{6.33}$$

due to the lack of the angle condition, does not imply

$$\left\| n^{-1} \Delta^2 v_n \right\| + \left\| \beta_n(v_n) \right\| \le c.$$
(6.34)

For this reason, the argument in (6.31) cannot be reproduced. To bypass this issue, we go back to the (approximate) entropy bound and we first recall that (4.28), as p > 12/11, gives at least

$$\left\|\Delta\varphi_n\right\|_{L^2(0,T;H)} \le c. \tag{6.35}$$

On the other hand, even in the worst case scenario for p, the second (5.8) guarantees

$$\|\sigma_n\|_{L^{8/5}(0,T;H)} \le c. \tag{6.36}$$

Hence, using (3.39) and comparing terms in (6.3), we deduce the analogue of (6.33), namely

$$\left\| n^{-1} \Delta^2 \varphi_n + \beta_n(\varphi_n) \right\|_{L^{8/5}(0,T;H)} \le c.$$
(6.37)

This estimate, however, as observed above, cannot be decoupled. For this reason, we need to go back again to (5.8) and we observe that it also implies

$$\|\sigma_n\|_{L^2(0,T;V')} \le c \|\sigma_n\|_{L^2(0,T;L^{12/7}(\Omega))} \le c.$$
(6.38)

Then, a further comparison of terms in (6.3) gives the estimate

$$\left\|n^{-1}\Delta^2\varphi_n + \beta_n(\varphi_n)\right\|_{L^2(0,T;V')} \le c,\tag{6.39}$$

which carries a weaker information with respect to space variables, but, on the other hand, can be decoupled. To see this, we recall (4.14), which, due to the presence of the elliptic regularizing term, takes now the form

$$\|\Delta\varphi_n\|^2 + n^{-1} \|\nabla\Delta\varphi_n\|^2 \le c \big(1 + \|\nabla\mu_n\|\big) + \chi \int_{\Omega} \nabla\varphi_n \cdot \nabla\sigma_n.$$
(6.40)

Then, following the lines of the entropy estimate, we may see that, in place of (4.27), one has

$$\|\Delta\varphi_n\|^{\frac{12p}{12-5p}} + n^{-\frac{6p}{12-5p}} \|\nabla\Delta\varphi_n\|^{\frac{12p}{12-5p}} \le c + c\|\nabla\mu_n\|^{\frac{6p}{12-5p}} + c\|\sigma_n\|^{\frac{p(5p-4)}{12-5p}}_{L^p(\Omega)} \left(1 + \|\nabla\sigma_n^{p/2}\|^2\right).$$
(6.41)

Hence, using again $\frac{12p}{12-5p} \ge 2$ we additionally deduce (at least)

$$\|n^{-1}\Delta\varphi_n\|_{L^2(0,T;V)} \le c$$
, or, more precisely, $\|n^{-1}\Delta^2\varphi_n\|_{L^2(0,T;V')} \le c$. (6.42)

Comparing with (6.39), we then obtain the desired "decoupled" information

$$\|\beta_n(\varphi_n)\|_{L^2(0,T;V')} \le c.$$
 (6.43)

Actually, this suffices to pass to the limit provided that we use a proper duality argument (an alternative, but in fact equivalent, approach may rely on the theory of variational inequalities). Namely, one may first observe that the sequence of monotone functions $\{\beta_n\}$ also induce a family of maximal monotone operators from $L^2(0,T;V)$ to $L^2(0,T;V')$. Moreover, such a family converges, in the sense of G-convergence, to the maximal monotone operator induced by β from $L^2(0,T;V)$ to the power set $2^{L^2(0,T;V')}$: let us denote by β_w , i.e., "weak β ", the operator obtained in this way. Actually, as first noticed in [5], β_w , due to the singular character of F, is a multivalued mapping; in particular there may occur concentration phenomena on the set where $|\varphi| = 1$.

Within this perspective, using the second (5.17) and (6.43), (6.30) is replaced by

$$\varphi_n \to \varphi$$
 strongly in $L^2(0,T;V)$ and $\beta_n(\varphi_n) \to \xi$ weakly in $L^2(0,T;V') \implies \xi \in \beta_w(\varphi)$, (6.44)

where we used the inclusion symbol because, as said, $\beta_w(\varphi)$ can contain more than an element. On the other hand, as the identification $\xi \in \beta_w(\varphi)$ is achieved, a posteriori one can prove that, in fact, $\beta_w(\varphi)$ can be replaced by $\beta(\varphi)$. More precisely, the limit ξ of $\beta_n(\varphi_n)$ is a function (and not just a functional), and it coincides almost everywhere with $\beta(\varphi)$. To see this, we first notice that, comparing terms in the *limit* equation (1.2) (where the bi-Laplacian no longer occurs), one deduces $\xi \in L^{8/5}(0,T;H)$ (the exponent 8/5 is due to (6.36)). Then, applying, e.g., [3, Prop. 2.5] (see also [29, Subsec. 2.1] for some additional background) one gets the desired (pointwise) equality $\xi = \beta(\varphi)$.

That said, the remainder of the argument used for passing to the limit with $n \nearrow \infty$ can be adapted to the approximation scheme up to purely tecnical modifications.

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