The Gravitational Lens Candidate FBQ 1633+3134[1]

Nicholas D. Morgan[2], Robert H. Becker[3,4], Michael D. Gregg[3,4], Paul L. Schechter[2],
Richard L. White[5]

[2]Department of Physics, Massachusetts Institute of Technology, Cambridge MA 02139;
ndmorgan@mit.edu, schech@achernar.mit.edu

[3]Physics Department, University of California, Davis, CA 95616

[4]Institute of Geophysics and Planetary Physics, Lawrence Livermore National
Laboratory, PO Box 808, L-413, Livermore, CA 94551-9900; bob@igpp.ucllnl.org,
gregg@igpp.ucllnl.org

[5]Space Telescope Science Institute, 3700 San Martin Drive, Baltimore, MD 21218;
rlw@stsci.edu

ABSTRACT

We present our ground-based optical imaging, spectral analysis, and high
resolution radio mapping of the gravitational lens candidate FBQ 1633+3134.
This $z = 1.52$, $B = 17.7$ quasar appears double on CCD images with an image
separation of $0\rlap{.}''66$ and a flux ratio of $\sim 3 : 1$ across $BVRI$ filters. A single 0.27
mJy radio source is detected at 8.46 GHz, coincident to within an arcsecond of
both optical components, but no companion at radio wavelengths is detected
down to a flux level of 0.1 mJy (3σ). Spectral observations reveal a rich
metal-line absorption system consisting of a strong Mg II doublet and associated
Fe I and Fe II absorption features, all at an intervening redshift of $z = 0.684$,
suggestive of a lensing galaxy. Point spread function subtraction however shows
no obvious signs of a third object between the two quasar images, and places
a detection limit of $I \gtrsim 23.0$ if such an object exists. Although the possibility
that FBQ 1633+3134 is a binary quasar cannot be ruled out, the evidence is
consistent with it being a single quasar lensed by a faint, metal-rich galaxy.

Subject headings: gravitational lensing — quasars: individual (FBQ 1633+3134)

---

1. INTRODUCTION

The number of quasars multiply imaged by gravitational lenses has steadily increased in the last decade. There are now over 40 well-established cases of galaxy-size gravitational lenses (Kochanek et al. 1998). The usefulness of gravitational lenses are well known (Kochanek & Hewitt 1996). For example, measurements of a time-delay between the light curves of individual images can lead to a determination of the Hubble constant (Refsdal 1964), and the observed frequency of strong lensing events can provide statistical constraints on the cosmological constant (Kochanek 1996). Also, since the observed image configuration is sensitive to the gravitational potential of the lens, gravitational lensing offers a means to probe the matter distribution of the lensing object (e.g., Maller et al. 2000). This provides a method to sample the dark matter halos of intermediate redshift galaxies which is independent of the galaxy's luminosity.

The identification of a close separation (i.e., $< 10''$) double quasar, however, is not confirmation of a gravitational lens. An alternative explanation is a physical binary quasar (e.g., Kochanek et al. 1999, Mortlock et al. 1999). In principle, the distinction between these two possibilities is clear: either there is direct evidence for an intervening galaxy between the quasar images, confirming the lensing scenario, or the spectral properties of the pair differ to the extent that the most likely explanation is that of two separate quasars. In the latter case, it is usually the failure to detect one of the quasars at radio wavelengths that eliminates the lensing hypothesis (e.g., Djorgovski et al. 1987; Muñoz et al. 1998). In practice, there are a number of quasar pairs that do not fall neatly into either category (Kochanek et al. 1999). The astrometric and spectral characteristics of these pairs are typical of gravitational lenses, yet no direct evidence for a lensing galaxy is found. The double quasar FIRST J163349.0+313412 (hereafter, FBQ 1633+3134), which we analyze here, is best classified as one of these "border-line" systems. In this paper, we present our ground-based optical imaging, spectral observations, and radio mapping of FBQ 1633+3134, and discuss the likelihood that the system is a gravitational lens.

FBQ 1633+3134 ($16^{\rm h}$ $33^{\rm m}$ $48\overset{\rm s}{.}97$, $+31°$ $34'$ $11\overset{''}{.}8$; J2000.0) was originally identified as a $z = 1.52$, $B = 17.3$ quasar from the FIRST Bright Quasar Survey (FBQS; Gregg et al. 1996). The target was first detected as a 2 mJy radio source at 20 cm as part of the FIRST radio survey (Becker, White, & Helfand 1995), and follow-up spectral observations taken with the Lick 3 m telescope confirmed the radio source as a high-redshift quasar (White et al. 2000). The FBQS offers a major advantage over fainter quasar surveys for the purpose of discovering lensed quasars. Since the majority of the FBQS targets are bright ($B < 18$), the probability of finding a quasar multiply imaged by gravitational lensing is enhanced over a sampling of fainter quasars due to the magnification bias effect (Turner, Ostriker, &

Gott 1984). Using the 2.4 m Hiltner Telescope at the MDM Observatory, we have exploited this bias by obtaining optical images of the brightest ($B < 18$) and more distant ($z > 1$) of the FBQS quasars to search for arcsecond scale multiple imaging indicative of strong gravitational lensing. One such observing run in 1996 December discovered the quasar FBQ 0951+2635 as a sub-arcsecond separation gravitational lens (Schechter et al. 1998). A subsequent observing run was carried out during 1997 June and targeted an additional 100 FBQS quasars in a further search for lenses. During this run, a 3 minute R-band exposure of FBQ 1633+3134 showed two stellar-like point sources separated by $\sim 0\rlap{.}''7$, and immediate follow-up exposures showed the same separation and similar flux ratios in BVI bandpasses as well, providing the initial evidence that FBQ 1633+3134 might be gravitationally lensed.

We describe our initial optical observations and analysis of FBQ 1633+3134 in §2, as well as higher resolution follow-up optical imaging taken immediately afterwards using the 3.5 m WIYN telescope. In §3, we describe the original FBQS spectrum of FBQ 1633+3134, as well as a follow-up spectrum taken with the Keck 10 m telescope. While radio observations were originally obtained as part of the FIRST survey, we have since obtained higher resolution VLA radio imaging of FBQ 1633+3134 following the optical detection of the object's double nature. We present these radio observations and their implications in §4. Finally, we discuss lens modeling and interpretation of the system in §5, and §6 summarizes our findings and conclusions for FBQ 1633+3134.

2. OPTICAL OBSERVATIONS AND REDUCTION

2.1. Initial Optical Imaging

Initial optical images of FBQ 1633+3134 were taken by one of us (P.L.S.) with the MDM 2.4 m Hiltner Telescope as part of the optical follow-up program for FBQS quasars described above. Imaging was conducted on the nights of 1997 June 1, 2, and 4. The Tek 1024 × 1024 CCD detector ("Charlotte") was used, with a gain of 3.16 e^- ADU^{-1} and a readnoise of 5.45 e^-. The telescope was operated at an f ratio of $f/7.5$, yielding a plate scale of $0\rlap{.}''2750$ per pixel. Sky conditions were marred by the presence of intermittent clouds on the nights of the 2nd and 4th, although observations reported here were taken during photometric periods. Multiple 3 minute exposures of FBQ 1633+3134 were taken using the Schombert $BVRI$ filter set, with seeing conditions that ranged from $0\rlap{.}''8$ to $1\rlap{.}''0$ FWHM. Multiple $BVRI$ observations of the Landolt standard field PG 1633 (Landolt 1992) were also taken on the night of the 2nd. A summary of the MDM observations for FBQ 1633+3134 are presented in Table 1. Figure 1 shows a $5'$ square field of the target quasar and nearby stars from one of the I-band frames.

All CCD frames were bias-subtracted, flattened, and trimmed using the VISTA reduction program. The flatfield frames consisted of twilight sky exposures taken on June 1, and were removed of cosmic rays using AUTOCLEAN, a program written and kindly provided by J. Tonry. The initial 3 minute R-band exposure of FBQ 1633+3134 showed it to be noticeably misshapen as compared to other stars in the field, with an ellipticity of ~ 0.6 (Fig. 2). This large ellipticity is consistent with a comparatively bright companion at less than an arcsecond away from the brighter optical component. There is also a faint point source $\sim 3''$ to the SW of the elongated image, with a peak flux of $\sim 1\%$ of the peak flux of the target quasar.

Preliminary photometry for the MDM data was performed using DoPHOT (Schechter, Mateo, & Saha 1993), and was able to split FBQ 1633+3134 as a double image in all bandpasses. We fit the elongated image with two empirical point spread functions (PSFs), using a variant of the DoPHOT program designed to deal with close, point-like and extended objects (Schechter & Moore 1993). Hereafter, we will refer to the brighter and fainter components of the two PSF fit as components A and B, respectively. We simultaneously fit for the presence of the faint SW object as well (hereafter, component C), to ensure that this object did not affect the photometric solutions for the main components. Star #5 in Fig. 1 served as the empirical PSF.

For our photometric fits, we have simultaneously solved for the relative fluxes of the system components and the overall position of the system, but have held the relative separations of the components fixed at values determined from the Hubble Space Telescope (HST) NICMOS imaging of E. Falco (PI), obtained as part of the CASTLES program (Kochanek et al. 1998) in 1998 August[2]. Using the archival HST/NICMOS images of FBQ 1633+3134, centroid positions for components A, B, and C were determined by the authors using gaussian fits to the peaks of the flux profiles; this gave a separation between the brighter two components of 0″663. Appropriate steps were then taken to account for the relative orientations and plate scales of the MDM and HST/NICMOS detectors. We report our photometric solutions for the MDM data using fixed relative positions for the three components in Table 2. Error bars in Table 2 are statistical (1σ) errors obtained from the frame-to-frame scatter for each filter.

The broadband magnitude differences between components A and B show little variation across wavelength. The photometric solutions reported in Table 2 correspond to A:B flux ratios of 2.45 in B, and 3.25, 3.24, and 3.36 in V, R, and I, respectively. This is

---

[2]The HST/NICMOS images of FBQ 1633+3134 may be viewed online at the CASTLES homepage at http://cfa-www.harvard.edu/castles

consistent with components A and B having similar spectral energy distributions (SEDs) at optical wavelengths, as would be expected if the system consisted of two quasars (Mortlock et al. 1999), or a single quasar multiply imaged by gravitational lensing. Component B has a $B - V$ color index of 0.114, and $V - R$ and $V - I$ color indices of 0.356 and 0.691, with uncertainties of ~ 0.04 for each index. The $B - V$ index is consistent with an A star, while the $V - R$ and $V - I$ indices are consistent with a late F to early G star. The color of component B is therefore inconsistent with it being a foreground Galactic star. The $B - V$, $V - R$, and $V - I$ color indices for component A are 0.422, 0.353, and 0.725, with uncertainties of ~ 0.02 for each index. Although these colors are consistent with a late F star, spectral observations presented in §3 make it highly unlikely that the brighter component is a Galactic star.

The photometric solutions reported in Table 2 also shed light on the nature of component C. Component C has $B - V$, $V - R$, and $V - I$ color indices of 0.95, 0.50, and 0.95 respectively, with (1σ) statistical errors of ~ 0.2 for $B - V$, and ~ 0.1 for $V - R$ and $V - I$. These colors are consistent with a late G to early K star. In addition, if we suppose components A and B do indeed arise from gravitational lensing, then the relative position of component C with respect to components A and B is inconsistent with it being either a third image of the lensed quasar or the lensing galaxy. Because it also it appears as a point source in our optical imaging, we conclude that component C is a foreground Galactic star and is not physically associated with the two brighter components of the system.

2.2. Astrometry and Photometric Calibration

The apparent magnitudes for FBQ 1633+3134 reported in Table 2 were calibrated using offsets and color terms obtained from observations of the Landolt standard field PG 1633 (Landolt 1992), where we have used extinction coefficients taken from the Kitt Peak direct imaging manual ($k_B = 0.25, k_V = 0.15, k_R = 0.10, k_I = 0.07$; Massey et al. 1997). DoPHOT yields flux ratios for all objects in the field with respect to the PSF template star. Once the PSF star has been calibrated onto the standard system, apparent magnitudes for the remainder of objects in the field are straightforward to obtain. We have also obtained astrometric solutions for reference objects in the field relative to the Automatic Plate Measuring (APM) astrometry of McMahon & Irwin (1992) using one of the I-band frames. To aid in the calibration of future observations, Table 3 presents our photometric and astrometric solutions for 8 reference stars within a 5 arcminute radius of the target quasar. The positional uncertainties in the astrometric solutions are $0\overset{''}{.}4$ (1σ).

2.3. Follow-Up Optical Imaging

To probe components A and B for possible signs of a lensing galaxy, higher resolution follow-up imaging of FBQ 1633+3134 was taken by C. Bailyn immediately following the MDM run. On 1997 June 5, a total of 6 CCD exposures of FBQ 1633+3134 were taken (2 each in Harris B and R, and 1 each in V and I) using the WIYN 3.5 m telescope at the Kitt Peak National Observatory. The 2SKB 2048 × 2048 detector was used, with a gain setting of 2.8 e^- ADU^{-1}, a readnoise of 8 e^-, and a plate scale of 0$.\!\!''$1971 pixel^{-1}. Seeing conditions for the series of exposures were 0$.\!\!''$7 FWHM, slightly better than for the MDM run. Each exposure lasted 120 s. These images were bias-subtracted, flattened, and trimmed using standard IRAF procedures, and the B- and R-band images were stacked using integer pixel offsets before obtaining photometric solutions. PSF fitting was then carried out using the identical procedure as for the MDM data, with star # 5 again providing the empirical PSF.

Results from PSF analysis, again using fixed relative positions taken from the HST/NICMOS data, yield A:B flux ratios of 2.67 in B, and 3.20, 3.21, 3.28 in V, R, and I. These results agree well with the corresponding MDM flux ratios for the two components. In Figure 2, we show an excised portion of the I-band frame, along with residuals after empirical PSF subtraction for the V and I filters. Tick marks for the residual panels indicate the centroid locations of components A (top-right) and B (bottom-left); the orientation of the system is identical in all four panels. We do not show the R-band residuals, as the R-band PSF was discovered to vary slightly over the face of the detector, which masked any residual signature from FBQ 1633+3134 that may have been present due to additional sources of flux in the system. The ellipticity of the R-band PSF increased from ~ 0.01 to ~ 0.06 for stars sampled from the center to the periphery of the detector, leaving an FBQ 1633+3134 residual pattern $\sim 1\%$ of component A's brightness. The R-band PSF-induced residual pattern seen for components A and B of FBQ 1633+3134 was similar to the patterns observed for isolated stars in the field.

Visual inspection of our ground-based B-, V- and I-band residuals show no obvious signs of unaccounted flux in the system. Although we did not perform a rigorous PSF subtraction of the HST/NICMOS data, a cursory inspection of the CASTLES image also shows no obvious indication of a lensing galaxy. If a lensing galaxy were located between components A and B, its presence would be most evident in our data in I-band, where the quasar-galaxy flux contrast would be more equal. To quantify the I-band magnitude limit for our non-detection, we inserted a series of gaussian flux profiles between components A and B, representative of an intervening galaxy, and investigated the residual patterns that emerged after fitting and subtracting two empirical PSFs as described above. The location of the gaussian profile was dictated by the singular isothermal sphere model (see §5), with

a FWHM of $0.\!\!''71$ (the average seeing for the I-band exposure). Again holding the relative separations of components A and B fixed at the HST/NICMOS values, we were able to detect 5σ residuals for a third object down to $I = 23.0$, where σ is the rms sky noise of the exposure. We therefore rule out the presence of an $I < 23$ magnitude point source at the suspected lensing galaxy position. The significance of this non-detection is discussed in §5.

3.   SPECTROSCOPY

A spectrum of FBQ 1633+3134 had been originally obtained by one of us (R.B.) using the Kast Double Spectrograph at the Lick Observatory on 1996 June 10 as part of the FIRST Bright QSO Survey (see White et al. 2000). The spectrum left no doubt that FBQ 1633+3134 was a quasar, clearly showing redshifted C IV, C III], and Mg II emission features. Several metal-line absorption features were also detected from an intervening line-of-sight absorber. However, the signal to noise (S/N) of the Lick spectrum was only ~ 10 at blue wavelengths, which hampered the absorption analysis.

We have since obtained a higher S/N spectrum using the Echellette Spectrograph and Imager (ESI; Epps & Miller 1998) at the Keck Observatory on 2000 March 9. The wavelength range of the observations was from 3700 Å to 8700 Å, with a dispersion that ranged from 0.7 to 11 Å from blue to red wavelengths. The slit orientation was held at a PA of 40° E of N, roughly perpendicular to the image separation. The seeing FWHM was $0.\!\!''8$ – insufficient to resolve the two components. The 600 s spectrum is shown in Figure 3.

Prominent absorption and emission features for the FBQ 1633+3134 spectrum are presented in Table 4. The source quasar has a redshift of 1.513(2) based on gaussian fits to peaks of the C III] and C IV emission features. The S/N of the Keck spectrum ranges from ~ 50 to ~ 90 from red to blue wavelengths, allowing the detection of an intervening metal-line absorption system at $z_{abs} = 0.684$. The Mg II $\lambda\lambda 2796, 2803$ Å doublet is clearly resolved, as is a series of Fe II $\lambda\lambda 2383, 2586, 2600$ Å and Fe I $\lambda 2474$ Å absorption lines. The Mg II doublet has a total rest-frame equivalent width of 2.4 Å. A strong (1.7 Å equivalent width) absorption feature at an observed wavelength of $\lambda 3708.9$ Å was also detected in the Lick spectrum, which may correspond to Ca I $\lambda 2201$ Å at the slightly higher redshift of $z_{abs} = 0.685$. The detection of a strong Mg II and Fe absorption system is evidence for a damped Lyman α system (DLAS) at the absorber redshift. A DLAS at intermediate redshift is likely to be the dusty, metal-rich disk of a late-type galaxy (Boisse et al. 1998).

The high S/N of the Keck spectrum provides a means to rule out the possibility that either component is a Galactic star. In general, the equivalent width of a spectral

absorption feature is reduced by a factor of $(1 + f)^{-1}$ when flooded with continuum flux from a companion object, where f represents the ratio of continuum intensities of the brighter object to the fainter object exhibiting the absorption feature. Prominent stellar absorption features typically have equivalent widths that are of order 1 Å (Jacoby et al. 1984). Thus, if the fainter optical component was indeed a Galactic star, the 3:1 optical flux ratio would predict stellar absorption features of typically 0.25 Å equivalent width. The equivalent width associated with the rms noise in the Keck spectrum ranges from 0.01 Å in the blue to 0.1 Å in the far red. The Keck spectrum therefore has a high enough S/N to confidently detect stellar absorption features if they were present. None of the common stellar absorption features are detected at the indicated positions (dashed lines) of Figure 3, which argues against either component being a foreground Galactic star. The simplest explanation for the combined spectrum is that we are seeing the light from two identical redshift quasars.

4.   RADIO OBSERVATIONS

FBQ 1633+3134 is a milli-Jansky radio source at 20 cm in the FIRST survey. The survey was carried out using the Very Large Array (VLA) in the B configuration, providing a resolution of $\sim 5''$ (Becker et al. 1995). The original radio observations lack the required resolution to probe for subarcsecond companions to the target quasar. Once optical observations had revealed the quasar's double nature, higher resolution VLA imaging was obtained.

On 1998 March 14, follow-up VLA observations were carried out by one of us (R.B.) at 8.46 GHz while the VLA was in the A configuration, providing an angular resolution of $0''.2$. Integration time was 60 minutes. Only one source was securely detected, with a peak intensity of 0.27 mJy and an rms noise level in the field of 0.03 mJy.

If we assume the radio emission comes from the brighter component of the optical pair, then the optical flux ratio of $\sim 3 : 1$ would predict a peak radio intensity for the fainter component of $\sim .1$ mJy. The failure to detect a radio component is therefore significant at the 3σ level.

The position of the radio source from the high resolution VLA imaging was determined to be 16^h 33^m $48^s.943$, $+31°$ $34'$ $11''.19$ (J2000.0), with a positional uncertainty of $\sim 0.2''$. This differs from our ground-based optical positions of component A (16^h 33^m $48^s.974$, $+31°$ $34'$ $11''.85$) by $\sim 0''.8$ and for component B (16^h 33^m $49^s.020$, $+31°$ $34'$ $11''.44$) by $\sim 1''.0$. The APM catalog position for the composite AB source shows the same difference. We also

queried the U.S. Naval Observatory A2.0 (USNO-A2.0) catalog, which uses the Precision Measuring Machine (PPM) astrometry of Monet et al. (1996), and found a similar $1\overset{''}{.}0$ offset with respect to the radio position. We are unable to identify with confidence which of the two optical components corresponds to the radio detection. A similar $1''$ discrepancy between VLA and optical astrometric solutions was also identified for the FIRST lensed quasar FBQ 0951+2635 (Schechter et al. 1998).

5.   MODEL AND INTERPRETATION

The identification of an intervening metal-line absorption system in the FBQ 1633+3134 spectrum raises the obvious question of whether this material could serve as the lens. Although this possibility is attractive from the lensing point of view, the mere presence of an intervening absorption system should not be taken as direct evidence of a lensing galaxy. The observed incidence of Mg II absorption systems in quasar spectra is ~ 1 such system per unit redshift range down to a minimum equivalent width of 0.3 Å (Steidel & Sargent 1992), so it is common to find an intervening metal-line absorber in the spectra of high-redshift quasars. In this section, we explore the likelihood that the $z_{abs} = 0.684$ absorber could indeed indicate the presence of the lensing galaxy, and investigate the properties of the putative lensing galaxy as constrained by other observables of the system. Throughout this section, we have adopted an $\Omega_m = 0.3$, $\Omega_\Lambda = 0.7$ cosmological model and have parameterized the Hubble constant as $H_o = 100h$ km s^{-1} Mpc^{-1}.

We will adopt the singular isothermal sphere (SIS) model to describe the gravitational potential of the hypothesized lensing galaxy. The SIS model is described by the line-of-sight velocity dispersion σ of the lensing galaxy. The corresponding Einstein radius is $\Theta_E = 4\pi \left(\frac{\sigma}{c}\right)^2 \frac{D_{LS}}{D_{OS}}$, where D_{LS} and D_{OS} are the angular diameter distances from the lens to the source and from the observer to the source, respectively (e.g., Narayan & Bartleman 1998). The SIS produces a two image lens configuration, with an angular separation between the two images of $2\Theta_E$. If the angular distance between the brighter (fainter) image and the core of the potential is denoted by AG (BG), then the ratio of image distances AG/BG is the same as the lensing-induced flux ratio of the two images. We further take the B-band luminosity of galaxies to be related to their line-of-sight velocity dispersions according to a Faber-Jackson (1976) relationship of the form $L/L_\star = (\sigma/\sigma_\star)^\gamma$. Following Keeton et al. (1998), we adopt parameters of $\sigma_\star = 220$ (144) km/s and $\gamma = 4$ (2.6) for early-type ellipticals (late-type spirals), where L_* corresponds to a B-band magnitude of $M_B^* = -19.7 + 5\log h$.

For a given lensing redshift z_l, one can use the observed image separation to estimate

the velocity dispersion and luminosity of the lensing galaxy. Figure 4 shows the predicted luminosity of the lensing galaxy as a function of redshift for both an early-type elliptical and a late-type spiral. At the redshift of the metal-line absorber (indicated by the dashed vertical line in Figure 4), the predicted luminosities for both galaxy types are comparable to an L_* galaxy; $\sim 0.3L_*$ for the elliptical, and $\sim 1.3L_*$ for the spiral, with a corresponding dark matter velocity dispersion of ~ 160 km/s for both galaxy types. The observed image separation can therefore be reproduced by a slightly overluminous spiral or underluminous elliptical (with respect to an L_* galaxy) at the absorber redshift. We also use the techniques of Kochanek (1992) to estimate the lens redshift probability distribution for the system. Using a critical lens radius of $r = 0\rlap{.}''33$ for FBQ 1633+3134, we compute a median redshift for an elliptical (spiral) lens galaxy of $z = 0.72$ (0.59), with a 2σ redshift interval of $0.21 \lesssim z \lesssim 1.15$ $(0.17 \lesssim z \lesssim 1.04)$. The detected metal-line absorption at $z_{abs} = 0.684$ lies within 1σ of the median lensing redshift for either type of galaxy.

Although the redshift and predicted velocity dispersion of any intervening galaxy are consistent with gravitational lensing, the predicted luminosity of the lens is difficult to reconcile with the detection limits presented in §2. Using spectral energy distributions for early- and late-type galaxies, we predict the apparent magnitude of the lensing galaxy as a function of the lensing redshift z_l. To calculate broadband apparent magnitudes, we adopt SEDs from Lilly (1997, private communication), which consist of interpolated and extrapolated curves from Coleman, Wu, and Weedman (1980). No evolution correction was applied to the energy distributions. The SEDs were normalized to their Faber-Jackson B-band luminosities at $4400(1 + z_l)$ Å, which yields the predicted absolute magnitudes on the AB magnitude system. The corresponding apparent magnitudes are then computed using the cosmological distance modulus relation

$$m_{AB}(\lambda_{obs}) - M_{AB}(\lambda_{rest}) = 5\log(\frac{D_L}{10\mathrm{pc}}) + 7.5\log(1 + z_l), \tag{1}$$

where D_L is the angular diameter distance to the lensing galaxy. The apparent AB magnitudes are finally transformed onto the standard $BVRI$ system by subtracting $-0.110, 0.011, 0.199$, and 0.456, respectively (Fukugita, Shimasaku, & Ichikawa 1995). The resulting magnitudes are accurate to within ± 1 magnitude.

Figure 5 presents our I-band magnitude predictions for both an early-type elliptical and late-type spiral galaxy as a function of lensing redshift. The predicted apparent magnitudes initially become fainter with increasing redshift, but eventually increase to brighter magnitudes as the lens redshift approaches the source redshift. This increase is because the velocity dispersion required to produce a fixed image separation increases without bound as the lens approaches the source, leading to a corresponding increase in the lens luminosity. For FBQ 1633+3134, the I-band detection limit of $I = 23$ from §2

roughly corresponds to the faintest predicted magnitude for an early-type elliptical, but ~ 1.5 magnitudes fainter than the faintest prediction for a late-type spiral. At the absorber redshift of $z_{abs} = 0.684$, the predicted magnitudes lie above the detection limit by ~ 0.5 (~ 2.0) magnitudes for an elliptical (spiral) galaxy. While an early-type elliptical lensing galaxy could conceivable escape our I-band detection limit, a late-type spiral should have been bright enough to detect in our optical imaging.

6. SUMMARY AND CONCLUSIONS

From an observational standpoint, the requirements for a double quasar to be confidently classified as a gravitational lens can be stated as follows: First, since gravitational lensing is achromatic, the optical flux ratio of the components should (ideally) be self-consistent across broadband filters. In practice, the ratios will not agree exactly due to differential reddening through the lens galaxy (McLeod et al. 1998) and microlensing by stars in the lens galaxy (Wisotzki et al. 1993). More importantly, the radio flux ratio, if available, must also agree with the optical ratio, again because of the achromatic nature of lensing. Second, spectral observations of the quasar pair must support the conclusion that both objects are identical redshift quasars with similar (although not necessarily identical; Wisotzki et al. 1993) emission and continuum properties. Third, the lensing galaxy itself must be detected, ideally by direct imaging, although weaker arguments based on the presence of intervening absorption features in the QSO spectrum have been used in the past (e.g., Hewett et al. 1994). In practice, the first two requirements are necessary for a lens classification, while the third (direct imaging of the lensing galaxy) is usually sufficient.

FBQ 1633+3134 meets the necessary requirements for a gravitational lens. The broadband flux ratio ($\sim 2.5 : 1$ in B, $\sim 3.3 : 1$ in VR, & I) is self-consistent across broadband wavelengths. Although only one component was detected in radio, the lower limit to the radio flux ratio of $\gtrsim 3 : 1$ is not in severe conflict with the corresponding optical flux ratio. Also, the spectral analysis argues against either component of the pair being a Galactic star, and the unresolved emission profiles from the pair are consistent with two identical-redshift quasars.

The sufficient condition that could elevate FBQ 1633+3134 to the status of a lensed quasar, direct detection of the lensing galaxy, is not realized with our optical imaging. Subtraction of two PSFs for the WIYN data presented in §2 show no significant indication of unaccounted flux in the system, and place a magnitude limit of $I > 23.0$ using a simple SIS model for the lensing potential. Although the hypothesized lens galaxy is not detected directly, a weaker argument for its presence can be made from the spectral identification

of a rich $z_{abs} = 0.684$ metal-line absorption system, which is within 1σ of the median lensing redshift. Also, the estimated luminosity of the suspected galaxy does not differ strongly from an L_* galaxy for either of the galaxy types considered here. On the other hand, the predicted apparent magnitude of a late-type spiral at the absorber redshift is ~ 2 magnitudes above the I-band detection threshold, while an early-type galaxy is only marginally consistent with the detection limit.

We conclude that FBQ 1633+3134 is a strong candidate for a close separation ($0\rlap{.}''66$) gravitational lens. The evidence presented here suggests lensing by a relatively faint ($I > 23.0$), $z = 0.684$ metal-rich galaxy, although the binary quasar scenario cannot be positively ruled out. Confirmation of the lensing hypothesis can be provided by a deeper radio probe for the fainter component of the system. We have obtained VLA A array time for late 2000 with this aim in mind, and intend to present these results, along with a systematic analysis of the HST data for this system, in a future paper.

REFERENCES

Becker, R. H., Helfand, R. L., & Helfand, D. J., 1995, ApJ, 450, 559

Boissé, P., Le Brun,, V., Bergeron, J., & Deharveng, J.-M., 1998, A&A, 333, 841

Coleman, G. D., Wu, C. -C., & Weedman, D. W. 1980, ApJS, 43, 393

Djorgovski, S., Perley, R., Meylan, G., & McCarthy, P., 1987, ApJ, 321, L17

Epps, H. W. & Miller, J. S., 1998, Proc. SPIE, 3355, 48

Faber, S. M., & Jackson, R. E., 1976, ApJ, 204, 668

Fukugita, M., Shimasaku, K., & Ichikawa, T. 1995, PASP, 107, 945

Gregg, M. D., Becker, R. H., White, R. L., Helfand, D. J., McMahon, R. G., & Hook, I. M., 1996, AJ, 112, 407

Hewett, P. C., Irwin, M. J., Foltz, C. B., Harding, M. E., Corrigan, R. T., Webster, R. L., & Dinshaw, N., 1993, AJ, 108, 1534

Jacoby, G. H., & Hunter, D. A., 1984, ApJ, 56, 257

Keeton, C. R., Kochanek, C. S., Falco, E. E., 1997, ApJ, 509, 561

Kochanek, C. S. 1992, ApJ, 384, 1

Kochanek, C. S. 1996, ApJ, 466, 638

Kochanek, C. S., & Hewitt, J. N. 1996 *Astrophysical Applications of Gravitational Lensing* (Dordrecht: Kluwer)

Kochanek, C. S., Falco, E. E., Impey, C. D., Lehár, J., McLeod, B. A., & Rix, H., -W. 1998, in AIP Conf. Proc. 470, After the Dark Ages: When galaxies Were Young, ed. S. S. Holt & E. P. Smith (New York: AIP), 163

Kochanek, C. S., Falco, E. E., & Muñoz, J. A. 1999, ApJ, 510, 590

Landolt, A. U. 1992, AJ, 104, 340

Maller, A. H, Simard, L., Guhathakurta, P., Hjorth, J., Jaunsen, A. O., Flores, R. A., & Primack, J. R., 2000, ApJ, 533, 194

Massey, P., Armandroff, T., De Veny, J., Claver, C., Harmer, C., Jacoby, G., Schoening, B., & Silva, D., 1997, Direct Imaging Manual for Kitt Peak (Tucson: NOAO)

McLeod, B. A., Bernstein, G. M., Rieke, M. J., Weedman, D. W., 1998, AJ, 115, 1377

McMahon, R. G., & Irwin, M. J., 1992, in Digitised Optical Sky Surveys, ed. H. T. MacGillivray & E. B. Thompson (Dordrecht: Kluwer), 417

Monet, D., Bird, A., Canzian, B., Harris, H., Reid, N., Rhodes, A., Sell, S., Ables, H., Dahn, C., Guetter, H., Henden, A., Leggett, S., Levison, H., Luginbuhl, C., Martini, J., Monet, A., Pier, J., Riepe, B., Stone, R., Vrba, F., Walker, R. 1996, USNO-SA1.0, (U.S. Naval Observatory, Washington DC).

Mortlock, D. J., Webster, R. L., & Francis, P. J., 1999, MNRAS, 309, 836

Muñoz, J. A., Falco, E. E., Kochanek, C. S., & Lehár, J., 1998, ApJ, 492, L9

Narayan, R., & Bartlemann, M. 1998, in Formation of Structure in the Universe, ed. A. Dekel & J. Ostriker (Cambridge: Cambridge Univ. Press)

Refsdal, S., 1964, MNRAS, 128, 307

Schechter, P. L., Mateo, M., & Saha, A., 1993, PASP, 105, 1342

Schechter, P. L. & Moore, C. B., 1993, AJ, 105, 1

Schechter, P. L., Gregg, M. D., Becker, R. H., Helfand, D. J., & White, R. L., 1998, AJ, 115, 1371

Steidel, C. C., & Sargent, W. L. W., 1991, AJ, 102, 1610

Turner, E. L, Ostriker, J. P., & Gott, J. R. III, 1984, ApJ, 284, 1

White, R. L., Becker, R. H., Gregg, M. D., Laurent-Muehleisen, S. A., Brotherton, M. S., Impey, C. D., Petry, C. E., Foltz, C. B., Chaffee, F. H., Richards, G. T., Oegerle, W. R., Helfand, D. J., McMahon, R. G., & Cabanela, J. E., 2000, ApJ, 126, 133

Wisotzki, L., Köhler, T., Kayser, R., & Reimers, D., 1993, A&A, 278, L15

Fig. 1.— A 3 minute I-band exposure of FBQ 1633+3134 (center crosshairs) and surrounding field. North is up and east is to the left. The scale of the image is shown in the bottom left of the figure.

Fig. 2.— Extracted subrasters of FBQ 1633+3134 taken with the WIYN 3.5 m telescope. North is up and East is to the left. Each panel is $\sim 4''$ wide. Top left: A 120 s I-band exposure, showing component A (NW) and component B (SE). The small tick mark SW of the pair indicates the centroid location of component C (which is too faint to be seen at the given contrast level). Top right: I-band residuals from two empirical PSF subtraction of components A and B. Tick marks indicate the centroid location of the two components. Bottom left: Stacked B-band residuals from two component PSF subtraction. Bottom right: V-band residuals from two component PSF subtraction. The saturation levels for the residual panels are at $\pm 10\sigma$, where σ is the respective rms sky noise for each image. Component A has a peak intensity of $\sim 580\sigma$, $\sim 482\sigma$, and $\sim 420\sigma$ for the B-(stacked), V-, and I- band exposures.

Fig. 3.— Composite spectrum of FBQ 1633+3134 A,B taken with the Keck 10 m telescope. Integration time was 600 s. Prominent emission features are indicated for this $z = 1.52$ quasar, as well as the identification of a $z_{abs} = 0.684$ metal-rich absorption system. The lack of any stellar absorption features at the indicated wavelengths (dashed lines) argues against either component being a foreground Galactic star.

Fig. 4.— Predicted B-band luminosity (in units of L_*) of the hypothesized lensing galaxy as a function of its lensing redshift z_l. Results for a late-type spiral and early-type elliptical are shown. The redshift of the intervening absorption system is indicated by the vertical dashed line. The corresponding line-of-sight velocity dispersion at the absorber redshift is ~ 160 km s^{-1}.

Fig. 5.— Predicted I-band apparent magnitudes of the hypothesized lensing galaxy as a function of its lensing redshift z_l. Results are shown for a late-type spiral (dot-dash line) and early-type elliptical (solid line). The $I = 23.0$ detection limit for any such galaxy is indicated by the heavy horizontal line, and the redshift of the intervening absorption system is shown by the vertical dashed line.

Table 1.   Summary of MDM Observations (June 1997)

Filter	N_{img}	Average FWHM
B	5	$1\rlap{.}''00$
V	12	$0\rlap{.}''82$
R	4	$0\rlap{.}''87$
I	8	$0\rlap{.}''87$

Table 2.   MDM Photometric Solutions for FBQ 1633+3134

Comp.	Δ R.A.	Δ Dec.	B	V	R	I
A........	—	—	18.046 ± 0.017	17.624 ± 0.006	17.271 ± 0.013	16.899 ± 0.010
B........	$0\rlap{.}''524$	$-0\rlap{.}''406$	19.019 ± 0.033	18.905 ± 0.019	18.549 ± 0.037	18.214 ± 0.034
C........	$-1\rlap{.}''793$	$-2\rlap{.}''282$	23.268 ± 0.167	22.315 ± 0.071	21.814 ± 0.037	21.370 ± 0.063

Note. — Relative astrometric positions were obtained from the archival CASTLES HST/NICMOS imaging. Magnitude solutions for the MDM data were obtained while holding the relative separations fixed at the HST values. Error bars are statistical (1σ) and do not include uncertainties in the PSF star calibration. See note to Table 3.

Table 3.   Relative Astrometry and Apparent Magnitudes for Field Stars

Object	$\Delta\alpha$ (^{s})	$\Delta\delta$ ($''$)	B	V	R	I
1	-3.010	115.11	18.299 ± 0.002	17.128 ± 0.001	16.407 ± 0.001	15.827 ± 0.002
2	-1.206	82.37	19.158 ± 0.002	18.537 ± 0.003	18.215 ± 0.002	17.864 ± 0.003
3	5.580	66.18	20.557 ± 0.006	19.296 ± 0.002	18.489 ± 0.002	17.775 ± 0.003
4	-1.829	63.61	19.853 ± 0.001	18.605 ± 0.002	17.776 ± 0.002	17.058 ± 0.002
5	0.000	0.00	18.670 ± 0.000	17.187 ± 0.000	16.059 ± 0.000	14.673 ± 0.000
6	4.275	-5.35	17.200 ± 0.002	16.225 ± 0.001	15.708 ± 0.001	15.242 ± 0.002
7	4.216	-8.93	17.556 ± 0.003	16.417 ± 0.001	15.729 ± 0.001	15.108 ± 0.002
8	0.885	-42.38	18.702 ± 0.003	17.810 ± 0.001	17.339 ± 0.002	16.904 ± 0.002

Note. — Object numbers correspond to the labels shown in Figure 1. Reported error bars are statistical (1σ) errors from the observed dispersion between frames, and do not include uncertainties in the calibration of the PSF star (object #5). Magnitude uncertainties (1σ) for the PSF star are 0.007, 0.008, 0.005, and 0.004 for $BVRI$ filters, respectively, and must be added in quadrature to the errors reported above.

Table 4.   Spectral Analysis for FBQ 1633+3134

Emission	Absorption	λ_{obs}	λ_{rest}	z
C IV	.....	3889.9	1549.3	1.5107(5)
C III]	.....	4799.7	1908.7	1.5146(11)
Mg II	.....	7064.4	2799.5	1.5235(5)
.....	Ca I	3708.9	2201.4	0.6848(4)
.....	Fe II	4012.0	2382.8	0.6837(2)
.....	Fe I	4165.8	2473.9	0.6839(3)
.....	Fe II	4354.8	2586.7	0.6836(2)
.....	Fe II	4377.9	2600.2	0.6837(1)
.....	Mg II	4707.5	2796.4	0.6834(3)
.....	Mg II	4719.8	2803.5	0.6835(1)

Note. — Numbers in parenthesis are 1σ error bars.

FBQ 1633+3134

elliptical

spiral cd

$z_{abs} = 0.684$

I = 23.0

predicted I band magnitude

lens redshift z_l